

Finite Difference Discretization of Hyperbolic Equations: Linear Problems

Lectures 8, 9 and 10

First Order Wave Equation

INITIAL BOUNDARY VALUE PROBLEM (IBVP)

$$\frac{\partial u}{\partial t} + U \frac{\partial u}{\partial x} = 0, \quad x \in (0, 1)$$

Initial condition: $u(x, 0) = u^0(x)$

Boundary conditions: $\begin{cases} u(0, t) = g_0(t) & \text{if } U > 0 \\ u(1, t) = g_1(t) & \text{if } U < 0 \end{cases}$

First Order Wave Equation

Solution

$$du = \frac{\partial u}{\partial t} dt + \frac{\partial u}{\partial x} dx = \left(\frac{\partial u}{\partial t} + \frac{dx}{dt} \frac{\partial u}{\partial x} \right) dt$$

If $\frac{dx}{dt} = U \Rightarrow [x = Ut + \xi]$ Characteristics



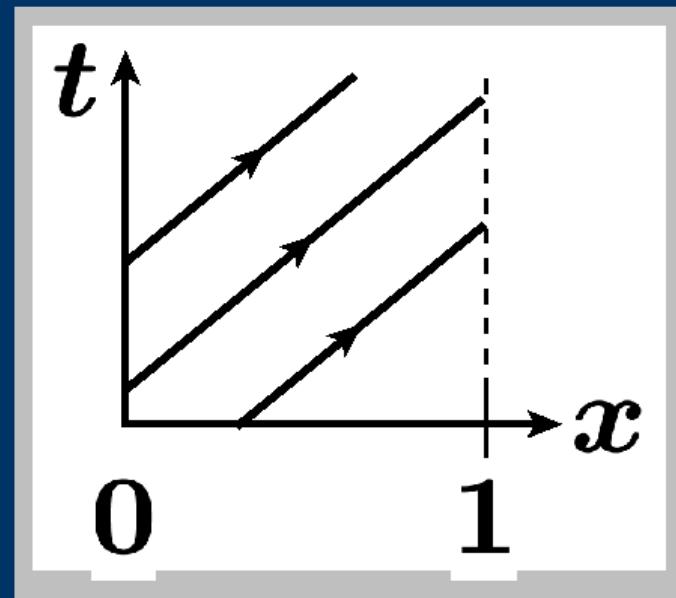
$$du = 0, \Rightarrow [u(x, t) = f(\xi) = f(x - Ut)]$$

General solution

First Order Wave Equation

Solution

$U > 0$

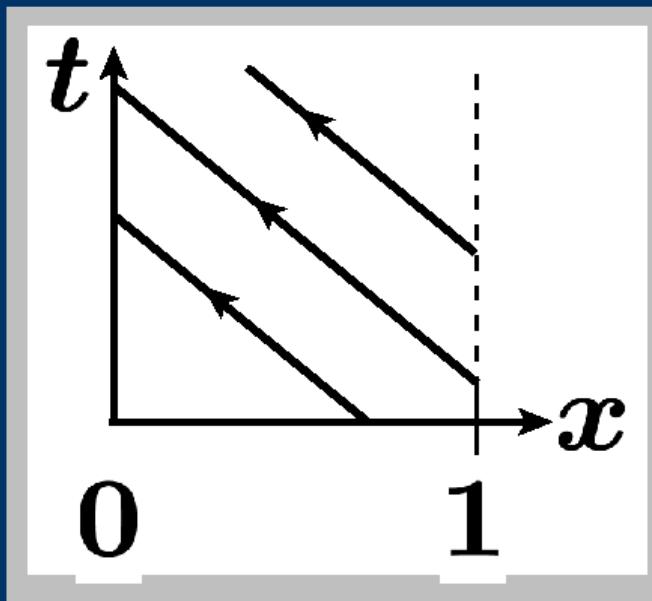


$$u(x, t) = \begin{cases} u^0(x - Ut), & \text{if } x - Ut > 0 \\ g_0(t - x/U), & \text{if } x - Ut < 0 \end{cases}$$

First Order Wave Equation

Solution

$U < 0$



$$u(x, t) = \begin{cases} u^0(x - Ut), & \text{if } x - Ut < 1 \\ g_1(t - x/U), & \text{if } x - Ut > 1 \end{cases}$$

First Order Wave Equation

Stability

$L^2([0, 1])$ -norm

$$\|u\|_2(t) = \left(\int_0^1 u^2(x, t) dx \right)^{\frac{1}{2}}$$

$$\int_0^1 u \left(\frac{\partial u}{\partial t} + U \frac{\partial u}{\partial x} \right) dx = 0$$

$$\frac{d}{dt} \|u\|_2^2 = -U(u^2(1, t) - u^2(0, t))$$

Model Problem

$$\frac{\partial \mathbf{u}}{\partial t} + U \frac{\partial \mathbf{u}}{\partial x} = \mathbf{0}, \quad x \in (0, 1)$$

Initial condition:

$$\mathbf{u}(x, 0) = \mathbf{u}^0(x)$$

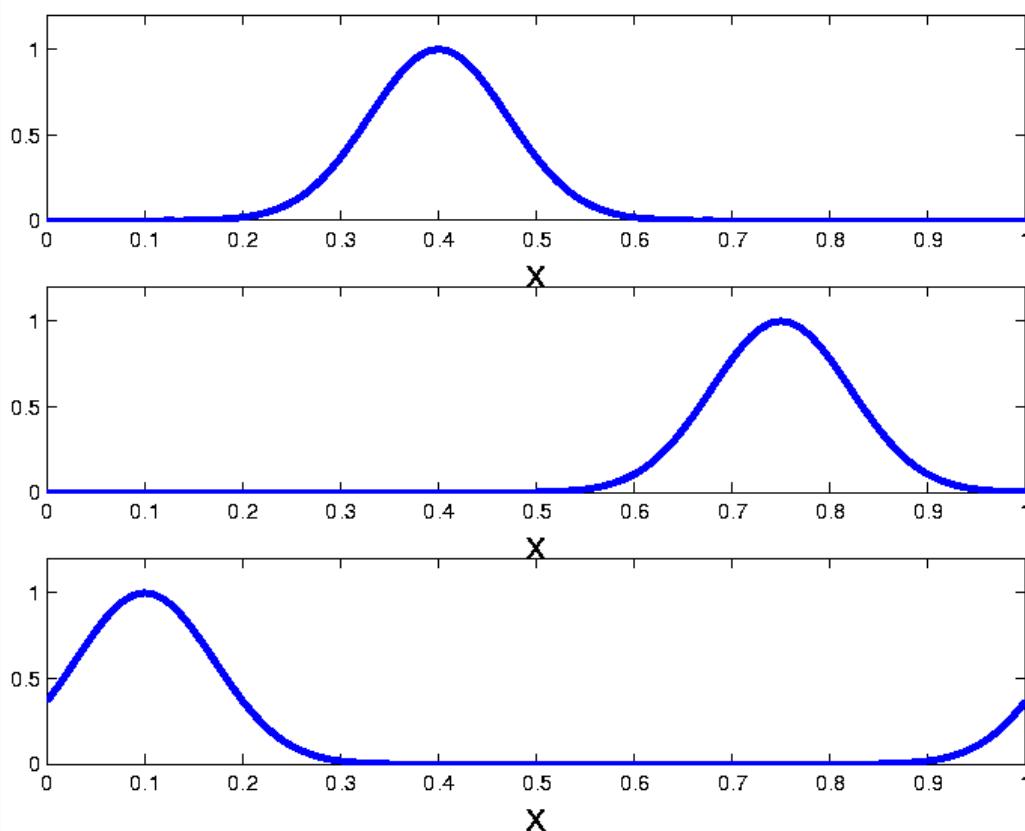
Periodic Boundary conditions: $\mathbf{u}(0, t) = \mathbf{u}(1, t)$

$$\frac{d}{dt} \|\mathbf{u}\|_2^2 = 0 \quad \Rightarrow \quad \|\mathbf{u}\|_2(t) = \|\mathbf{u}^0\|_2 = \text{constant}$$

Model Problem

Example

Periodic Solution ($U > 0$)



$t = 0$

$t = T$

$t = 2T$

Finite Difference Solution

Discretization

Discretize $(0, 1)$ into J equal intervals Δx

$$\Delta x = \frac{1}{J}, \quad x_j = j\Delta x$$

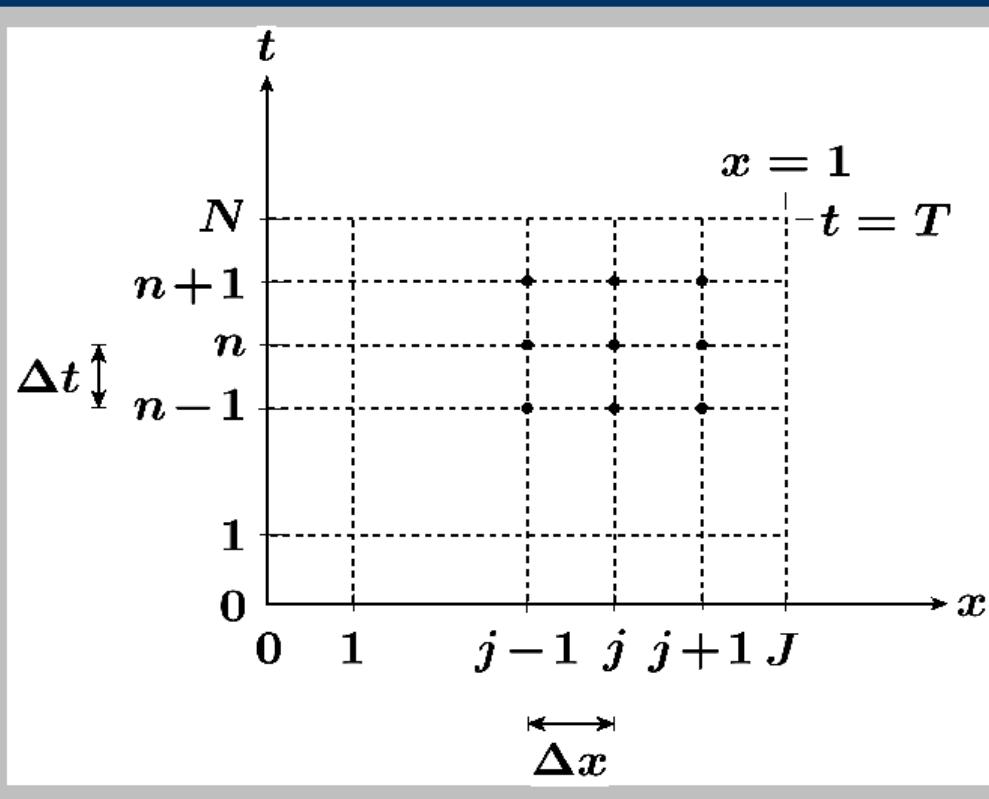
and $(0, T)$ into N equal intervals Δt

$$\Delta t = \frac{T}{N}, \quad t^n = n\Delta t$$

$$\hat{u}_j^n \approx u_j^n \equiv u(x_j, t^n), \quad \text{for } \begin{cases} 0 \leq j \leq J \\ 0 \leq n \leq N \end{cases}$$

Finite Difference Solution

Discretization



Discretization

Finite Difference Solution

NOTATION:

- \hat{v}_j^n approximation to $v(x_j, t^n) \equiv v_j^n$
- $\underline{v}^n \in \mathbb{R}^J$ vector of approximate values at time n ;

$$\underline{v}^n = \{\hat{v}_j^n\}_{j=1}^J$$

- $\underline{v}^n \in \mathbb{R}^J$ vector of exact values at time n ;

$$\underline{v}^n = \{v(x_j, t^n)\}_{j=1}^J$$

Approximation

Finite Difference Solution

For example . . . (for $U > 0$)

$$\left. \frac{\partial v}{\partial x} \right|_j^n \approx \frac{v(x_j, t^n) - v(x_{j-1}, t^n)}{\Delta x} = \frac{v_j^n - v_{j-1}^n}{\Delta x}$$

$$\left. \frac{\partial v}{\partial t} \right|_j^n \approx \frac{v(x_j, t^{n+1}) - v(x_j, t^n)}{\Delta t} = \frac{v_j^{n+1} - v_j^n}{\Delta t}$$

Forward in Time Backward (Upwind) in Space

Finite Difference Solution

First Order Upwind Scheme

$u_t + U u_x = 0$ suggests ...

$$\frac{\hat{u}_j^{n+1} - \hat{u}_j^n}{\Delta t} + U \frac{\hat{u}_j^n - \hat{u}_{j-1}^n}{\Delta x} = 0 \quad \Rightarrow$$

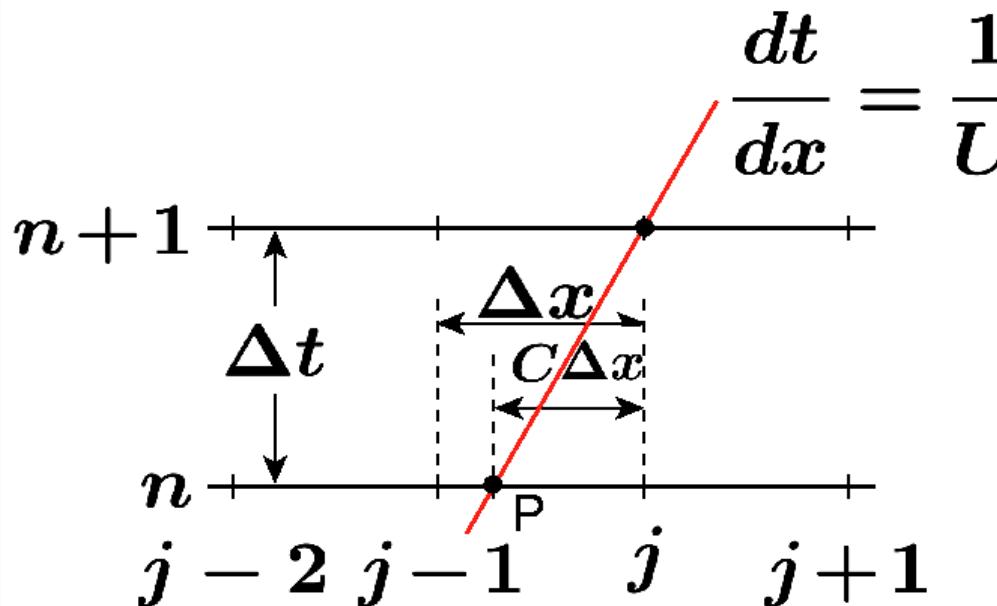
$$\begin{aligned}\hat{u}_j^{n+1} &= \hat{u}_j^n - C(\hat{u}_j^n - \hat{u}_{j-1}^n) & \begin{cases} 1 \leq j \leq J \\ 0 \leq n \leq N \end{cases} \\ \hat{u}_0^n &= \hat{u}_J^n & 0 \leq n \leq N\end{aligned}$$

Courant number $C = U \Delta t / \Delta x$

Finite Difference Solution

First Order Upwind Scheme

Interpretation



$$u_j^{n+1} = u_P$$

Use Linear
Interpolation

$$j-1, j$$

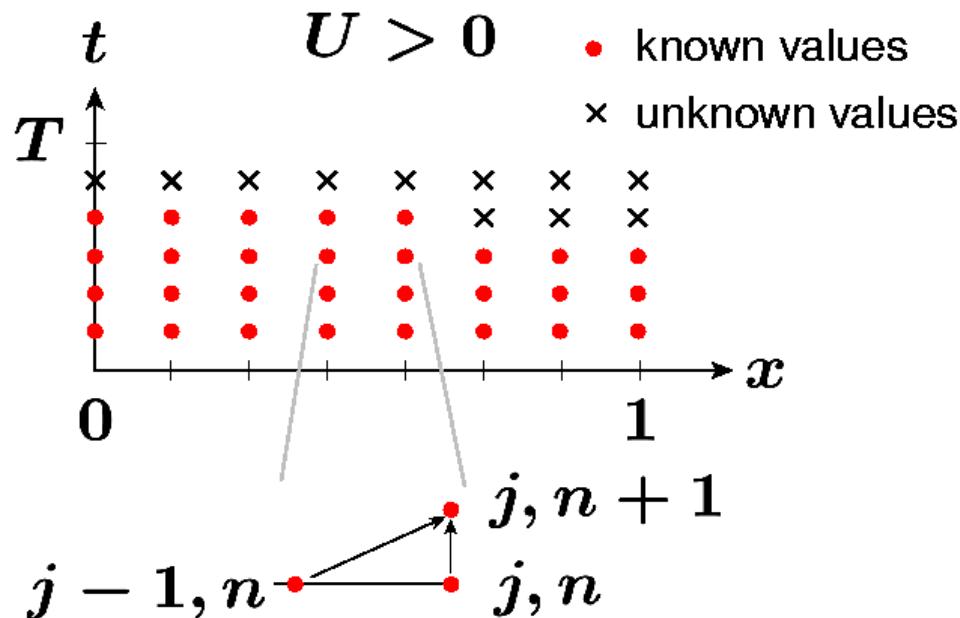
N1

$$u_P \approx C\hat{u}_{j-1}^n + (1 - C)\hat{u}_j^n$$

Finite Difference Solution

First Order Upwind Scheme

Explicit Solution



no matrix inversion

\hat{u}^n exists and is unique

$$\hat{u}_j^{n+1} = \hat{u}_j^n - C(\hat{u}_j^n - \hat{u}_{j-1}^n)$$

Finite Difference Solution

First Order Upwind Scheme

Matrix Form

We can write

$$\hat{\underline{u}}^n = \hat{\mathcal{S}} \hat{\underline{u}}^{n-1} = \hat{\mathcal{S}}^n \hat{\underline{u}}^0$$

$$\hat{\underline{u}}^0 \equiv \underline{u}^0$$

$$\underbrace{\begin{pmatrix} (1-C) & 0 & 0 & \cdots & C \\ C & (1-C) & 0 & \cdots & 0 \\ 0 & \cdots & \cdots & \ddots & 0 \\ \vdots & \cdots & C & (1-C) & 0 \\ 0 & \cdots & 0 & C & (1-C) \end{pmatrix}}_{\hat{\mathcal{S}}}$$

Finite Difference Solution

First Order Upwind Scheme

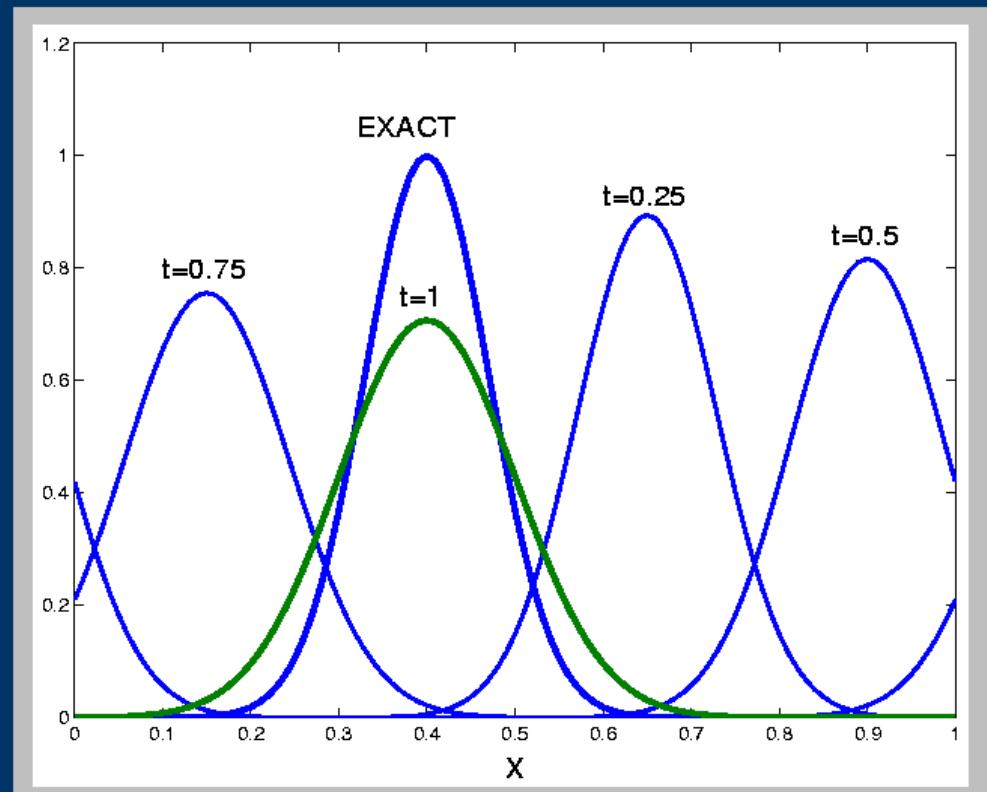
Example

$$u_t + u_x = 0$$

$$\Delta x = \frac{1}{100}$$

$$C = \frac{\Delta t}{\Delta x} = 0.5$$

$$T = 1 \Rightarrow N = 200$$



Definition

Convergence

The finite difference algorithm **converges** if

$$\lim_{\Delta x, \Delta t \rightarrow 0} \|\hat{\underline{u}}^n - \underline{u}^n\| = 0, \quad 1 \leq n \leq N$$
$$N\Delta t = T$$
$$J\Delta x = 1$$

for **any** initial condition $\underline{u}^0(x)$.

$$\|\underline{v}\| = \left(\Delta x \sum_{j=1}^J v_j^2 \right)^{1/2} = \sqrt{\Delta x} \|\underline{v}\|_2$$

N2

Definition

Consistency

The difference scheme $\hat{\mathcal{L}}\hat{\mathbf{u}}^n = \mathbf{0}$,
is **consistent** with the differential equation $\mathcal{L}\mathbf{u} = \mathbf{0}$

If:

For **all** smooth functions \mathbf{v}

$$(\hat{\mathcal{L}}\underline{\mathbf{v}}^n)_j - (\mathcal{L}\mathbf{v})_j^n \rightarrow \mathbf{0}, \quad \text{for } \begin{cases} 1 \leq j \leq J \\ 1 \leq n \leq N \end{cases}$$

when $\Delta x, \Delta t \rightarrow 0$.

Consistency

Difference operator

$$\hat{\mathcal{L}}\underline{v}^n = \frac{1}{\Delta t}\{\underline{v}^{n+1} - \hat{\mathcal{S}}\underline{v}^n\}$$

Differential operator

$$\mathcal{L}v \equiv \frac{\partial v}{\partial t} + U \frac{\partial v}{\partial x}$$

Consistency

$$\begin{aligned}
 (\hat{\mathcal{L}}\underline{v}^n)_j &\equiv \frac{\underline{v}_j^{n+1} - \underline{v}_j^n}{\Delta t} + \mathbf{U} \frac{\underline{v}_j^n - \underline{v}_{j-1}^n}{\Delta x} \\
 &= (\underline{v}_t + \mathbf{U}\underline{v}_x)_j^n + \frac{\Delta t}{2}(\underline{v}_{tt})_j^n + \mathbf{U} \frac{\Delta x}{2}(\underline{v}_{xx})_j^n + \dots
 \end{aligned}$$

$$(\mathcal{L}\underline{v})_j^n \equiv (\underline{v}_t + \mathbf{U}\underline{v}_x)_j^n$$

$$(\hat{\mathcal{L}}\underline{v}^n)_j - (\mathcal{L}\underline{v})_j^n = \mathcal{O}(\Delta x, \Delta t)$$

⇒ First order accurate in space and time.

Truncation Error

Insert exact solution \underline{u} into difference scheme

$$(\hat{\mathcal{L}} \underline{u})_j^n - \underbrace{(\mathcal{L} \underline{u})_j^n}_{=0} = \tau_j^n, \quad \text{for } \begin{cases} 1 \leq j \leq J \\ 1 \leq n \leq N \end{cases}$$

$$\underline{u}^{n+1} = \hat{\mathcal{S}} \underline{u}^n + \Delta t \underline{\tau}^n$$

Consistency $\Rightarrow \|\underline{\tau}^n\| = \mathcal{O}(\Delta x, \Delta t), \quad 1 \leq n \leq N$

Definition

Stability

The difference scheme $\hat{\underline{u}}^{n+1} = \hat{\mathcal{S}}\hat{\underline{u}}^n$ is **stable** if:
there exists C_T such that

$$\|\underline{v}^n\| = \|\hat{\mathcal{S}}^n \underline{v}^0\| \leq C_T \|\underline{v}^0\|$$

for all \underline{v}^0 ; and $n, \Delta t$ such that $0 \leq n\Delta t \leq T$

Above condition can be written as

$$\|\hat{\mathcal{S}} \underline{v}\| \leq (1 + \mathcal{O}(\Delta t)) \|\underline{v}\|$$

First Order Upwind Scheme

Stability

$$\hat{u}_j^{n+1} = \hat{u}_j^n - C(\hat{u}_j^n - \hat{u}_{j-1}^n)$$

$$= (1 - C) \hat{u}_j^n + C \hat{u}_{j-1}^n$$

$$= \alpha \hat{u}_j^n + \beta \hat{u}_{j-1}^n$$

Stability

$$\begin{aligned} \sum_{j=1}^J |\hat{u}_j^{n+1}|^2 &= \sum_{j=1}^J |\alpha \hat{u}_j^n + \beta \hat{u}_{j-1}^n|^2 \\ &\leq \sum_{j=1}^J |\alpha|^2 |\hat{u}_j^n|^2 + 2|\alpha||\beta| |\hat{u}_j^n| |\hat{u}_{j-1}^n| + |\beta|^2 |\hat{u}_{j-1}^n|^2 \\ &\leq \sum_{j=1}^J |\alpha|^2 |\hat{u}_j^n|^2 + |\alpha||\beta| (|\hat{u}_j^n|^2 + |\hat{u}_{j-1}^n|^2) + |\beta|^2 |\hat{u}_{j-1}^n|^2 \\ &= \sum_{j=1}^J (|\alpha|^2 + 2|\alpha||\beta| + |\beta|^2) |\hat{u}_j^n|^2 = (|\alpha| + |\beta|)^2 \sum_{j=1}^J |\hat{u}_j^n|^2 \end{aligned}$$

First Order Upwind Scheme

Stability

$$\|\underline{u}^{n+1}\|_2^2 \leq (|\alpha| + |\beta|)^2 \|\underline{u}^n\|_2^2$$

Stability if

$$|\alpha| + |\beta| \leq 1, \quad \Rightarrow$$

$$|1 - C| + |C| \leq 1, \quad 0 \leq C \leq 1$$

Upwind scheme is **stable** provided

$$U > 0, \quad \Delta t \leq \frac{\Delta x}{U}$$

Lax Equivalence Theorem

A **consistent** finite difference scheme for a partial differential equation for which the initial value problem is well-posed is **convergent** if and only if it is **stable**.

Lax Equivalence Theorem

Proof

$$\begin{aligned}\|\hat{\underline{u}}^n - \underline{u}^n\| &= \|\hat{\mathcal{S}}\hat{\underline{u}}^{n-1} - \hat{\mathcal{S}}\underline{u}^{n-1} + \Delta t \underline{\tau}^{n-1}\| \\ &\leq \|\hat{\mathcal{S}}(\hat{\underline{u}}^{n-1} - \underline{u}^{n-1})\| + \Delta t \mathcal{O}(\Delta x, \Delta t) \\ &\leq \|\hat{\underline{u}}^{n-1} - \underline{u}^{n-1}\| + \Delta t \mathcal{O}(\Delta x, \Delta t) \\ &\quad \vdots \\ &\leq \underbrace{\|\hat{\underline{u}}^0 - \underline{u}^0\|}_{=0} + \underbrace{n \Delta t}_{\leq T} \mathcal{O}(\Delta x, \Delta t) \\ &\leq \mathcal{O}(\Delta x, \Delta t) \quad (\text{first order in } \Delta x, \Delta t)\end{aligned}$$

Lax Equivalence Theorem

First Order Upwind Scheme

- Consistency: $\|\underline{\tau}\| = \mathcal{O}(\Delta x, \Delta t)$
- Stability: $\|\underline{\hat{u}}^{n+1}\| \leq \|\underline{\hat{u}}^n\|$ for $C \equiv U \Delta t / \Delta x \leq 1$
- \Rightarrow Convergence $\underline{e} = \underline{u} - \underline{\hat{u}}$

$$\|\underline{e}^n\| \leq (C_x \Delta x + C_t \Delta t), \quad 1 \leq n \leq N$$

$$\text{or } |e_j^n| \leq (C_x \Delta x + C_t \Delta t), \quad \begin{cases} 1 \leq j \leq J, \\ 1 \leq n \leq N \end{cases}$$

C_x and C_t are constants independent of $\Delta x, \Delta t$

Lax Equivalence Theorem

First Order Upwind Scheme

Example

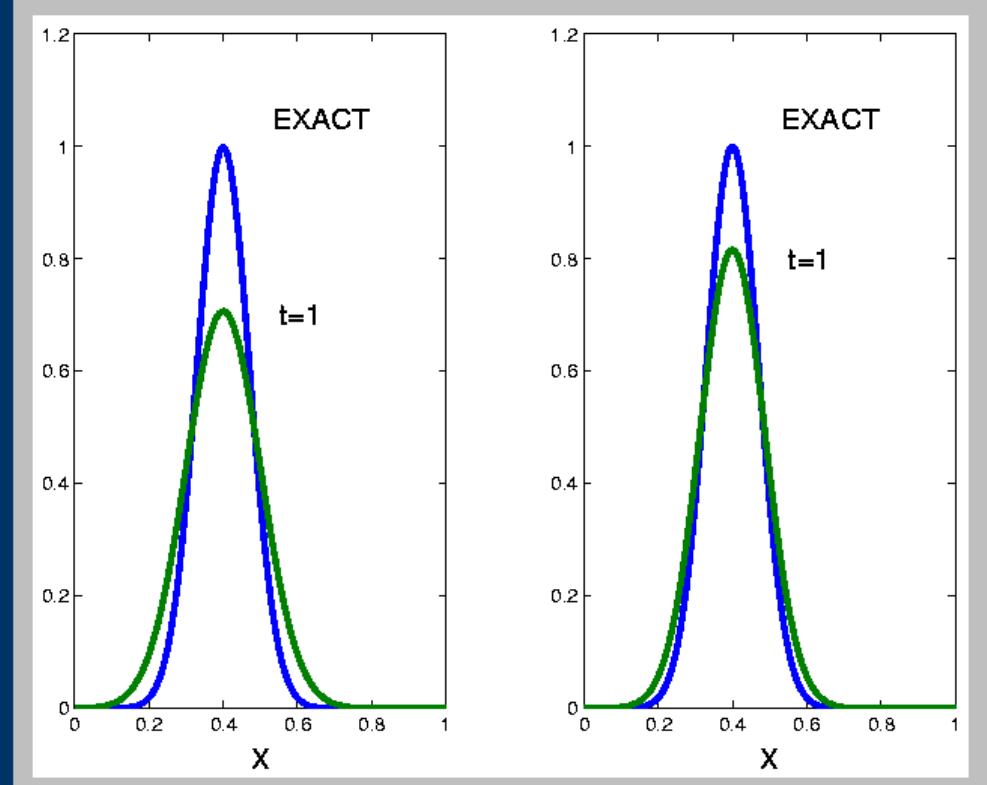
Solutions for:

$$C = 0.5$$

$$\Delta x = 1/100 \text{ (left)}$$

$$\Delta x = 1/200 \text{ (right)}$$

Convergence is slow !!



CFL Condition

Mathematical Domain of Dependence of $u(x_j, t^N)$

Set of points in (x, t) where the initial or boundary data may have some effect on $u(x_j, t^N)$.

Numerical Domain of Dependence of \hat{u}_j^N

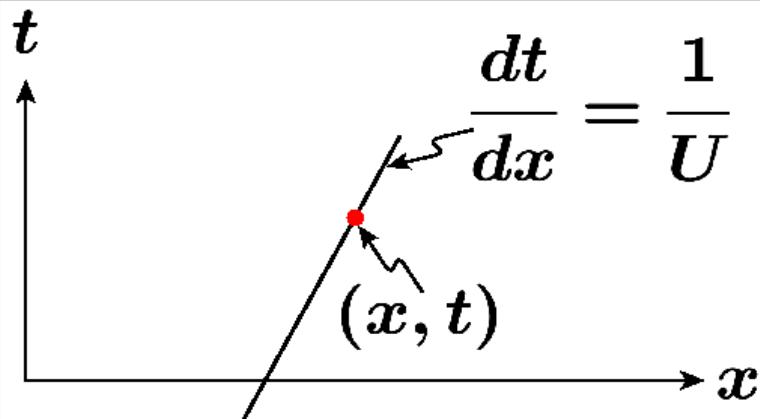
Set of points x_k, t^n where the initial or boundary data may have some effect on \hat{u}_j^N .

N3

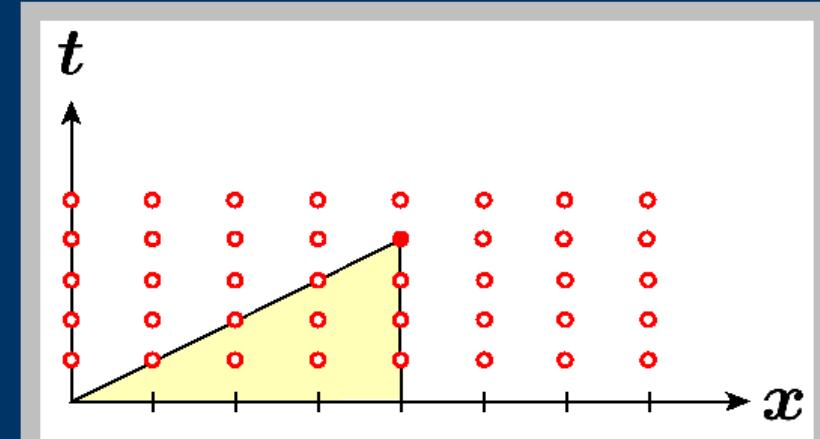
Domains of Dependence

CFL Condition

First Order Upwind Scheme



Analytical



Numerical ($U > 0$)

CFL Theorem

CFL Condition

CFL Condition

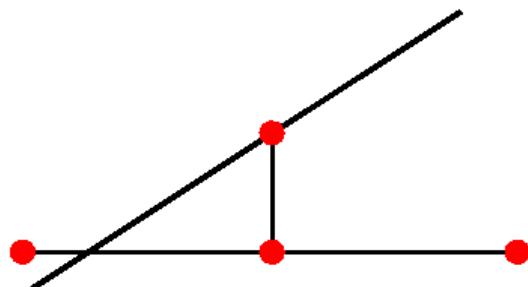
For each (x_j, t^N) the **mathematical domain of dependence** is **contained** in the **numerical domain of dependence**.

CFL Theorem

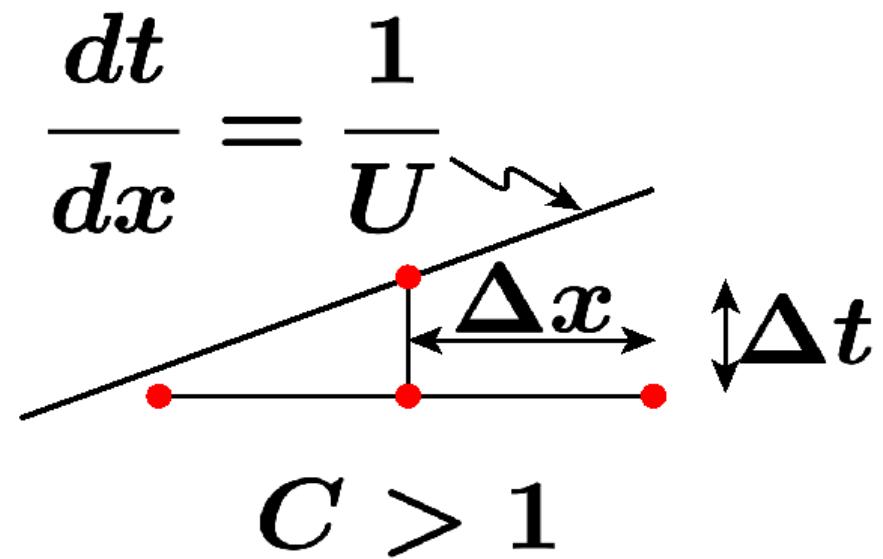
The **CFL condition** is a **necessary** condition for the **convergence** of a numerical approximation of a partial differential equation, linear or nonlinear.

CFL Theorem

CFL Condition



$$C < 1$$



$$C > 1$$

Stable

Unstable

Fourier Analysis

- Provides a systematic method for determining stability → von Neumann Stability Analysis
- Provides insight into discretization errors

Fourier Analysis

Continuous Problem

Fourier Modes and Properties...

Fourier mode: $\Phi_k(x) = e^{i2\pi kx}$, $k \in \mathbb{Z}$ (integer)

- Periodic (period = 1)
- Orthogonality

$$\int_0^1 \Phi_k(x) \Phi_{-k'}(x) dx = \delta_{kk'}$$

- Eigenfunction of $\frac{\partial^m}{\partial x^m} \Phi_k(x) = (i2\pi k)^m \Phi_k(x)$

Fourier Analysis

Continuous Problem

...Fourier Modes and Properties

- Form a basis for periodic functions in $L^2([0, 1])$

$$v(x) = \sum_{k=-\infty}^{\infty} \mathbb{V}_k \Phi_k(x) = \sum_{k=-\infty}^{\infty} \mathbb{V}_k e^{i2\pi kx}$$

- Parseval's theorem

$$\|v\|_2^2 = \sum_{k=-\infty}^{\infty} |\mathbb{V}_k|^2$$

Fourier Analysis

Continuous Problem

Wave Equation

$$u(x, t) = \sum_{k=-\infty}^{\infty} \mathbb{U}_k(t) \Phi_k(x) = \sum_{k=-\infty}^{\infty} \mathbb{U}_k(t) e^{i2\pi kx}$$

$$u_t + U u_x = 0 \Rightarrow \sum_{k=-\infty}^{\infty} \left(\frac{d\mathbb{U}_k}{dt} + i2\pi k U \mathbb{U}_k \right) e^{i2\pi kx} = 0$$

$$u^0(x) = \sum_{k=-\infty}^{\infty} \mathbb{U}_k^0 e^{i2\pi kx} \Rightarrow \mathbb{U}_k(t) = \mathbb{U}_k^0 e^{-i2\pi k U t}$$

Fourier Analysis

Discrete Problem

Fourier Modes and Properties...

Fourier mode: $\underline{\Phi}_k = \{\Phi_k(x_j)\}_{j=0}^{J-1},$

k (integer) $\in (-J/2 + 1, J/2)$

$$\Phi_k(x_j) = e^{i2\pi k j \Delta x} \equiv e^{ij\theta} = \Phi_{\theta j}, \quad \theta = 2\pi k \Delta x$$

$$k \in (-J/2 + 1, J/2) \Rightarrow \theta \in (-\pi + 2\pi \Delta x, \pi)$$

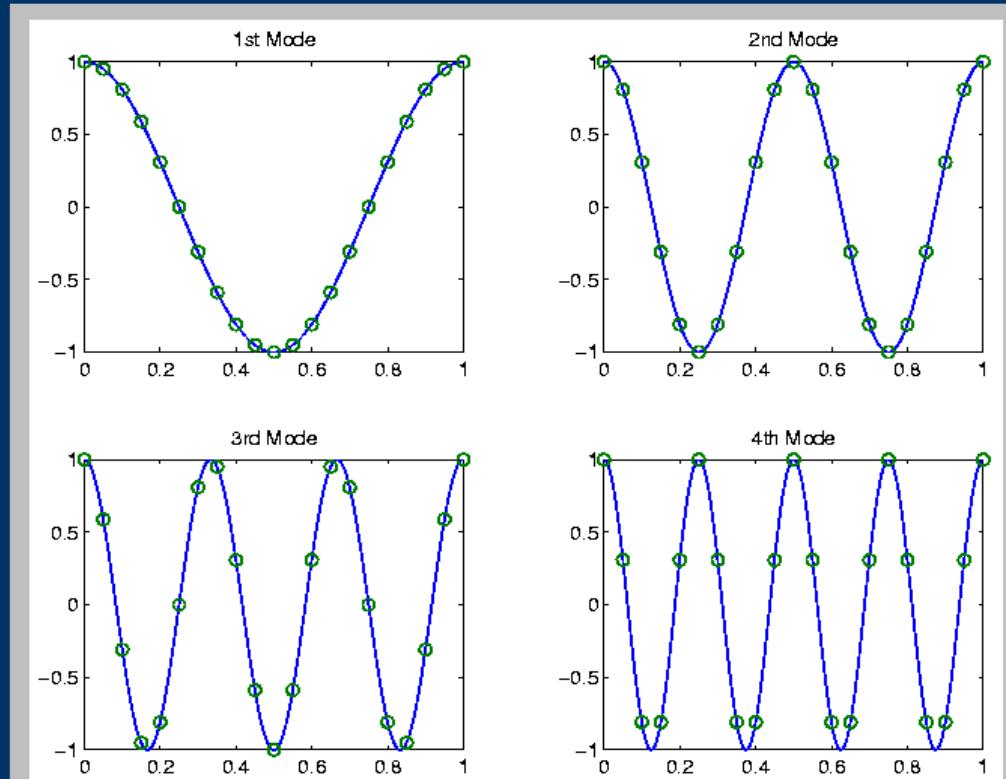
Fourier Analysis

Discrete Problem

...Fourier Modes and Properties...

Real part of first 4 Fourier modes

$$\Delta x = 1/20$$



Fourier Analysis

Discrete Problem

...Fourier Modes and Properties...

- Periodic (period = J)
- Orthogonality

N4

$$\frac{1}{J} \underline{\Phi}_{\theta}^T \underline{\Phi}_{-\theta'} = \frac{1}{J} \sum_{j=0}^{J-1} e^{i2\pi kj\Delta x} e^{-i2\pi k' j\Delta x} = \delta_{kk'}$$

$$= \frac{1}{J} \sum_{j=0}^{J-1} e^{ij\theta} e^{-ij\theta'} = \begin{cases} 1 & \text{if } \theta = \theta' \\ 0 & \text{if } \theta \neq \theta' \end{cases}$$

Fourier Analysis

Discrete Problem

...Fourier Modes and Properties...

- Eigenfunctions of difference operators e.g.,

N5

$$- \delta_{2x} \underline{v}|_j = v_{j+1} - v_{j-1}$$

$$\delta_{2x} \underline{\Phi}_\theta = i 2 \sin(\theta) \underline{\Phi}_\theta$$

$$- \delta_x^2 \underline{v}|_j = v_{j+1} - 2v_j + v_{j-1}$$

$$\delta_x^2 \underline{\Phi}_\theta = -4 \sin^2(\theta/2) \underline{\Phi}_\theta$$

$$- \Delta_x^- \underline{v}|_j = v_j - v_{j-1}$$

$$\Delta_x^- \underline{\Phi}_\theta = (1 - e^{-i\theta}) \underline{\Phi}_\theta$$

Fourier Analysis

Discrete Problem

...Fourier Modes and Properties

- Basis for periodic (discrete) functions $\underline{v} = \{v_j\}_{j=1}^J$

$$\underline{v} = \sum_{\theta = -\pi \\ +2\pi\Delta x}^{\pi} \mathbb{V}_\theta \underline{\Phi}_\theta \quad \rightarrow \quad v_j = \sum_{\theta = -\pi \\ +2\pi\Delta x}^{\pi} \mathbb{V}_\theta e^{ij\theta}$$

- Parseval's theorem

$$\|\underline{v}\|^2 \equiv \underbrace{\Delta x}_{1/J} \|\underline{v}\|_2^2 = \sum_{\theta = -\pi \\ +2\pi\Delta x}^{\pi} |\mathbb{V}_\theta|^2$$

Fourier Analysis

Write $\hat{\underline{u}}^{n+1} = \sum_{\theta} \hat{\mathbb{U}}_{\theta}^{n+1} \underline{\Phi}_{\theta}$, $\hat{\underline{u}}^n = \sum_{\theta} \hat{\mathbb{U}}_{\theta}^n \underline{\Phi}_{\theta}$

Stability $\|\hat{\underline{u}}^{n+1}\| \leq (1 + \mathcal{O}(\Delta t)) \|\hat{\underline{u}}^n\|$

$$\Rightarrow \sum_{\theta} |\hat{\mathbb{U}}_{\theta}^{n+1}|^2 \leq (1 + \mathcal{O}(\Delta t)) \sum_{\theta} |\hat{\mathbb{U}}_{\theta}^n|^2$$

Stability for all data \Rightarrow

$$|\hat{\mathbb{U}}_{\theta}^{n+1}| \leq (1 + \mathcal{O}(\Delta t)) |\hat{\mathbb{U}}_{\theta}^n|, \quad \forall \theta$$

Fourier Analysis

von Neumann Stability Criterion

First Order Upwind Scheme...

$$\hat{u}_j^n = \sum_{\theta} \hat{\mathbb{U}}_{\theta}^n \Phi_{\theta j} = \sum_{\theta} \hat{\mathbb{U}}_{\theta}^n e^{ij\theta}$$

$$\hat{u}_j^{n+1} - \hat{u}_j^n + C(\hat{u}_j^n - \hat{u}_{j-1}^n) = 0, \quad \forall j \Rightarrow$$

$$\sum_{\theta} (\hat{\mathbb{U}}_{\theta}^{n+1} - \hat{\mathbb{U}}_{\theta}^n + C(1 - e^{-i\theta}) \hat{\mathbb{U}}_{\theta}^n) e^{ij\theta} = 0, \quad \forall j \Rightarrow$$

Fourier Analysis

von Neumann Stability Criterion

...First Order Upwind Scheme...

$$\hat{U}_\theta^{n+1} = \frac{((1 - C) + Ce^{-i\theta})}{g(C, \theta)} \hat{U}_\theta^n = g(C, \theta) \hat{U}_\theta^n$$

amplification factor

Stability if $|\hat{U}_\theta^{n+1}| \leq |\hat{U}_\theta^n|, \quad \forall \theta$ which implies

$$|g(C, \theta)| \leq 1, \quad \forall \theta$$

Fourier Analysis

von Neumann Stability Criterion

...First Order Upwind Scheme

$$\begin{aligned}|g(C, \theta)|^2 &= |(1 - C) + Ce^{-i\theta}|^2 \\&= (1 - C + C \cos(\theta))^2 + C^2 \sin^2(\theta) \\&= (1 - 2C \sin^2(\theta/2))^2 + 4C^2 \sin^2(\theta/2) \cos^2(\theta/2) \\&= 1 - 4C(1 - C) \sin^2(\theta/2)\end{aligned}$$

Stability if:

$$|g(C, \theta)| \leq 1 \Rightarrow 0 \leq C \equiv \frac{U \Delta t}{\Delta x} \leq 1$$

Fourier Analysis

von Neumann Stability Criterion

FTCS Scheme...

$$\frac{\hat{u}_j^{n+1} - \hat{u}_j^n}{\Delta t} + U \frac{\hat{u}_{j+1}^n - \hat{u}_{j-1}^n}{2\Delta x} = 0$$

$$\Rightarrow \underline{\hat{u}}^{n+1} = \underline{\hat{u}}^n - \frac{C}{2} \delta_{2x} \underline{\hat{u}}^n$$

Fourier Decomposition: $\underline{u}_j^n = \sum_{\theta} \hat{\mathbb{U}}_{\theta}^n e^{ij\theta}$

$$\Rightarrow \sum_{\theta} (\hat{\mathbb{U}}_{\theta}^{n+1} - \hat{\mathbb{U}}_{\theta}^n + iC \sin(\theta) \hat{\mathbb{U}}_{\theta}^n) e^{ij\theta} = 0$$

Fourier Analysis

von Neumann Stability Criterion

...FTCS Scheme

$$\hat{U}_\theta^{n+1} = \underbrace{(1 - iC \sin(\theta))}_{g(C, \theta)} \hat{U}_\theta^n = g(C, \theta) \hat{U}_\theta^n$$

amplification factor

$$|g(C, \theta)|^2 = 1 + C^2 \sin^2(\theta) \geq 1, \quad \text{for } C \neq 0$$

⇒ Unconditionally Unstable ⇒ Not Convergent

Lax-Wendroff Scheme

Time Discretization

Write a Taylor series expansion in time about t^n

$$u(x, t^{n+1}) = u(x, t^n) + \Delta t \frac{\partial u}{\partial t} \Big|_{t^n} + \frac{\Delta t^2}{2} \frac{\partial^2 u}{\partial t^2} \Big|_{t^n} + \dots$$

But ...

$$\frac{\partial u}{\partial t} = -U \frac{\partial u}{\partial x} \quad (\text{from } u_t + U u_x = 0)$$

$$\frac{\partial^2 u}{\partial t^2} = \frac{\partial}{\partial t} \left(-U \frac{\partial u}{\partial x} \right) = -U \frac{\partial}{\partial x} \left(\frac{\partial u}{\partial t} \right) = U^2 \frac{\partial u^2}{\partial x^2}$$

Lax-Wendroff Scheme

Spatial Approximation

$$u(x, t^{n+1}) = u(x, t^n) - U \Delta t \left. \frac{\partial u}{\partial x} \right|_j + \frac{U^2 \Delta t^2}{2} \left. \frac{\partial^2 u}{\partial x^2} \right|_j + \dots$$

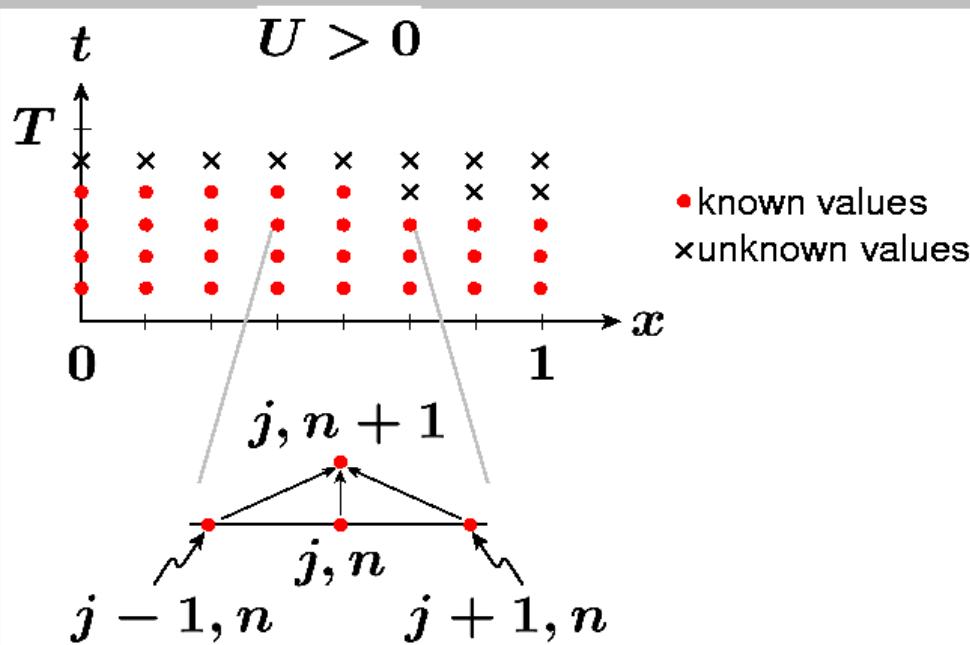
Approximate spatial derivatives

$$\left. \frac{\partial v}{\partial x} \right|_j \approx \frac{1}{2 \Delta x} \delta_{2x} v|_j = \frac{v_{j+1} - v_{j-1}}{2 \Delta x}$$

$$\left. \frac{\partial^2 v}{\partial x^2} \right|_j \approx \frac{1}{\Delta x^2} \delta_x^2 v|_j = \frac{v_{j+1} - 2v_j + v_{j-1}}{\Delta x^2}$$

Equations

Lax-Wendroff Scheme



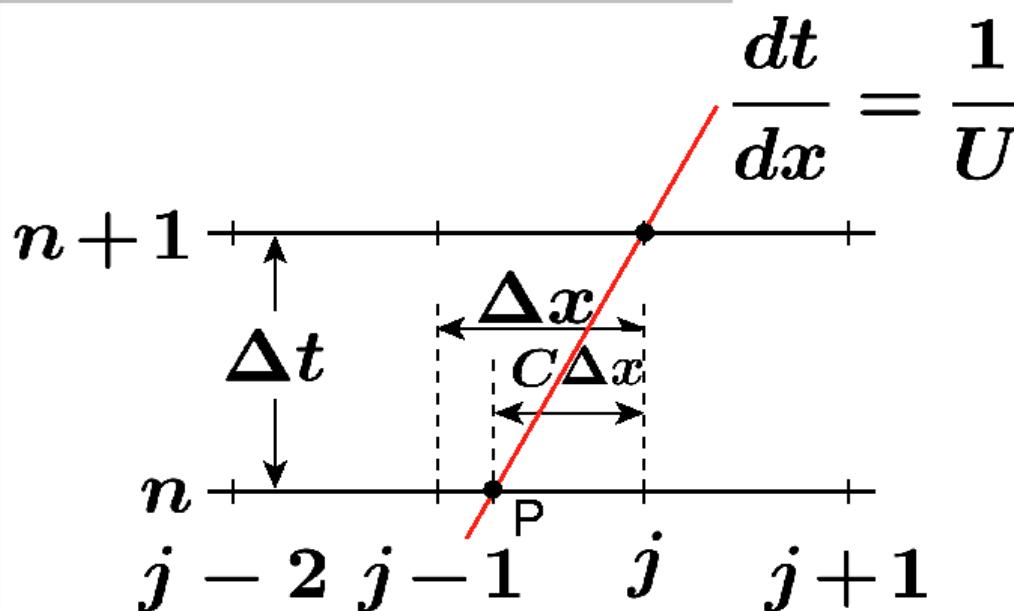
no matrix inversion

\hat{u}^n exists and is unique

$$\hat{u}_j^{n+1} = \hat{u}_j^n - \frac{C}{2}(\hat{u}_{j+1}^n - \hat{u}_{j-1}^n) + \frac{C^2}{2}(\hat{u}_{j+1}^n - 2\hat{u}_j^n + \hat{u}_{j-1}^n)$$

Lax-Wendroff Scheme

Interpretation



$$u_j^{n+1} = u_P$$

Use Quadratic
Interpolation
 $j-1, j, j+1$

$$u_P \approx \frac{C}{2}(1+C)\hat{u}_{j-1}^n + (1+C)(1-C)\hat{u}_j^n - \frac{C}{2}(1-C)\hat{u}_{j+1}^n$$

Lax-Wendroff Scheme

Analysis

Consistency

$$\begin{aligned}(\hat{\mathcal{L}}\underline{\boldsymbol{v}}^n)_j &\equiv \frac{\boldsymbol{v}_j^{n+1}-\boldsymbol{v}_j^n}{\Delta t} + \boldsymbol{U} \frac{\boldsymbol{v}_{j+1}^n-\boldsymbol{v}_{j-1}^n}{2\Delta x} - \frac{\boldsymbol{U}^2 \Delta t}{2} \frac{\boldsymbol{v}_{j+1}^n-2\boldsymbol{v}_j^n+\boldsymbol{v}_{j-1}^n}{\Delta x^2} \\&= (\boldsymbol{v}_t + \boldsymbol{U}\boldsymbol{v}_x)_j^n + \underbrace{\frac{\Delta t}{2} \left(\boldsymbol{v}_{tt}|_j^n - \boldsymbol{U}^2 \boldsymbol{v}_{xx}|_j^n \right)}_{=0 \text{ (for } v=u\text{)}} + \dots \\(\mathcal{L}\boldsymbol{v})_j^n &\equiv (\boldsymbol{v}_t + \boldsymbol{U}\boldsymbol{v}_x)_j^n\end{aligned}$$

$$(\hat{\mathcal{L}}\underline{\boldsymbol{v}}^n)_j - (\mathcal{L}\boldsymbol{v})_j^n = \mathcal{O}(\Delta x^2, \Delta t^2)$$

⇒ Second order accurate in space and time.

Lax-Wendroff Scheme

Analysis

Truncation Error

Insert exact solution \underline{u} into difference scheme

$$(\hat{\mathcal{L}} \underline{u})_j^n - \underbrace{(\mathcal{L} \underline{u})_j^n}_{=0} = \tau_j^n, \quad \text{for } \begin{cases} 1 \leq j \leq J \\ 1 \leq n \leq N \end{cases}$$

$$\underline{u}^{n+1} = \hat{\mathcal{S}} \underline{u}^n + \Delta t \tau^n$$

Consistency $\Rightarrow \|\tau^n\| = \mathcal{O}(\Delta x^2, \Delta t^2), \quad 1 \leq n \leq N$

Lax-Wendroff Scheme

Analysis

Stability

$$\begin{aligned}\hat{\underline{u}}^{n+1} &= \hat{\underline{u}}^n - \frac{C}{2} \delta_{2x} \hat{\underline{u}}^n + \frac{C^2}{2} \delta_x^2 \hat{\underline{u}}^n \\ \Rightarrow \hat{\mathbb{U}}_\theta^{n+1} &= \hat{\mathbb{U}}_\theta^n - iC \sin(\theta) \hat{\mathbb{U}}_\theta^n - C^2(1 - \cos(\theta)) \hat{\mathbb{U}}_\theta^n \\ &= \underbrace{\left(1 - 2C^2 \sin^2(\theta/2) - iC \sin(\theta)\right)}_{g(C,\theta)} \hat{\mathbb{U}}_\theta^n \\ |g(C,\theta)|^2 &= 1 - 4C^2(1 - C^2) \sin^4(\theta/2)\end{aligned}$$

Stability if: $|g(C, \theta)| \leq 1 \Rightarrow |C| \equiv |U|\Delta t/\Delta x \leq 1$

Lax-Wendroff Scheme

Analysis

Convergence

- Consistency: $\|\underline{\tau}\| = \mathcal{O}(\Delta x^2, \Delta t^2)$
- Stability: $\|\hat{\underline{u}}^{n+1}\| \leq \|\hat{\underline{u}}^n\|$ for $C \equiv U\Delta t/\Delta x \leq 1$
- \Rightarrow Convergence

$$\underline{e} = \underline{u} - \hat{\underline{u}}$$

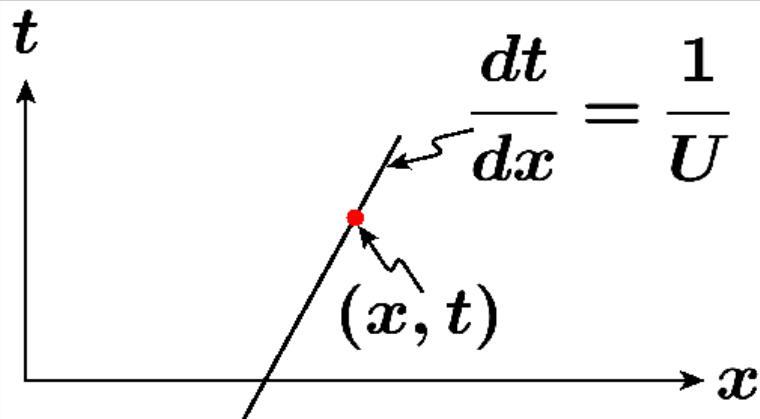
$$\|\underline{e}^n\| \leq (C_x \Delta x^2 + C_t \Delta t^2), \quad 1 \leq n \leq N$$

$$\text{or } |e_j^n| \leq (C_x \Delta x^2 + C_t \Delta t^2), \quad \begin{cases} 1 \leq j \leq J, \\ 1 \leq n \leq N \end{cases}$$

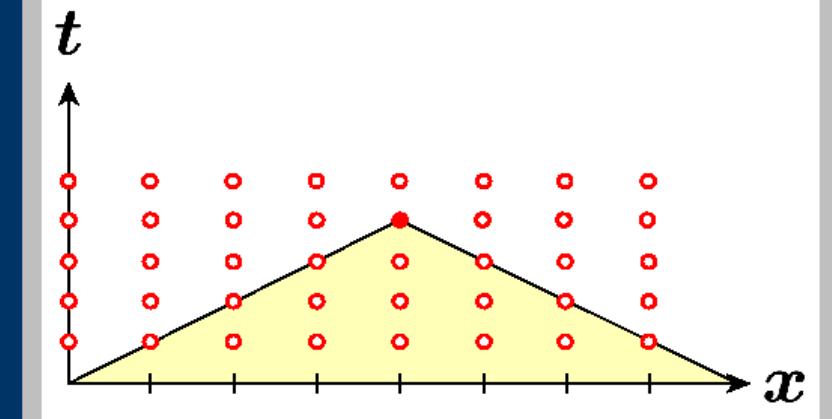
C_x and C_t are constants independent of Δx , Δt

Lax-Wendroff Scheme

Domains of Dependence



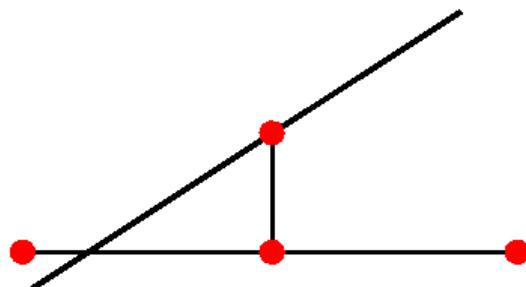
Analytical



Numerical

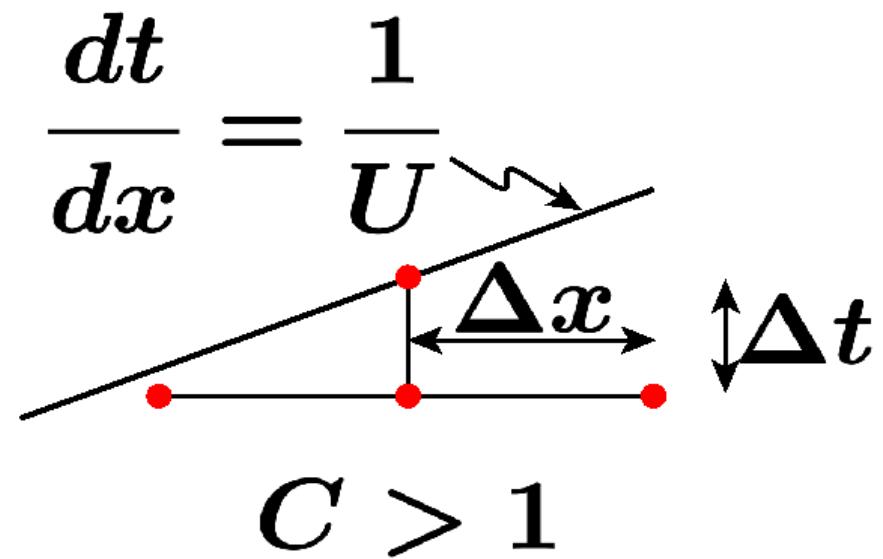
Lax-Wendroff Scheme

CFL Condition



$$C < 1$$

Stable



$$C > 1$$

Unstable

Lax-Wendroff Scheme

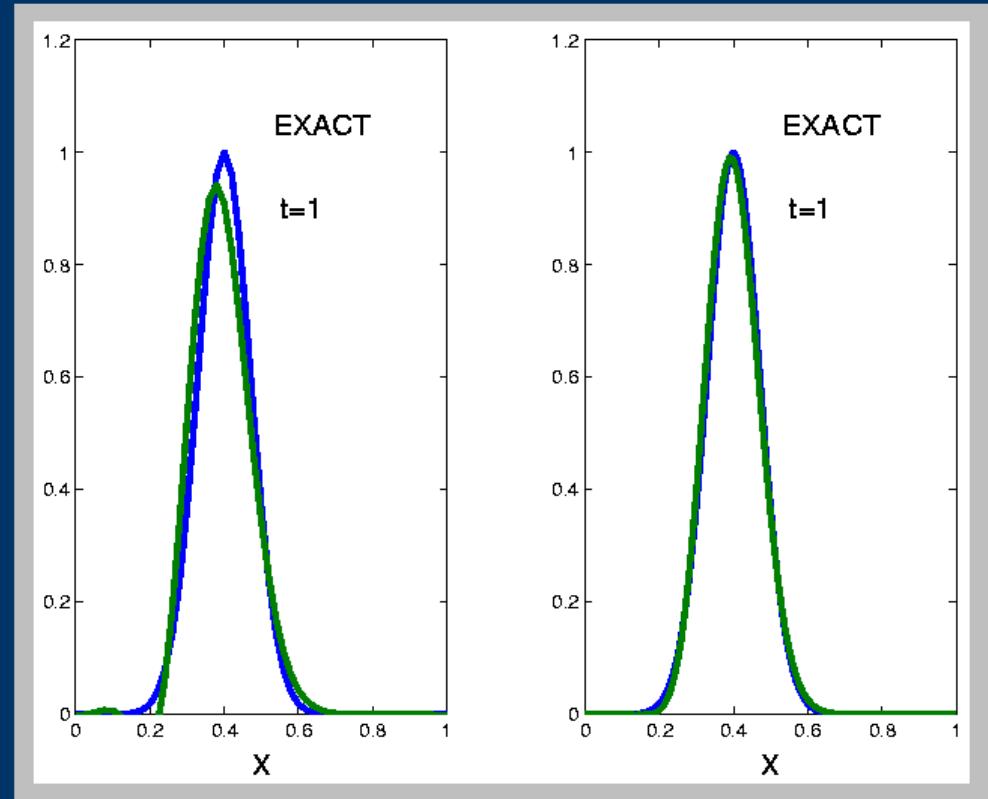
Example

Solutions for:

$$C = 0.5$$

$$\Delta x = 1/50 \text{ (left)}$$

$$\Delta x = 1/100 \text{ (right)}$$



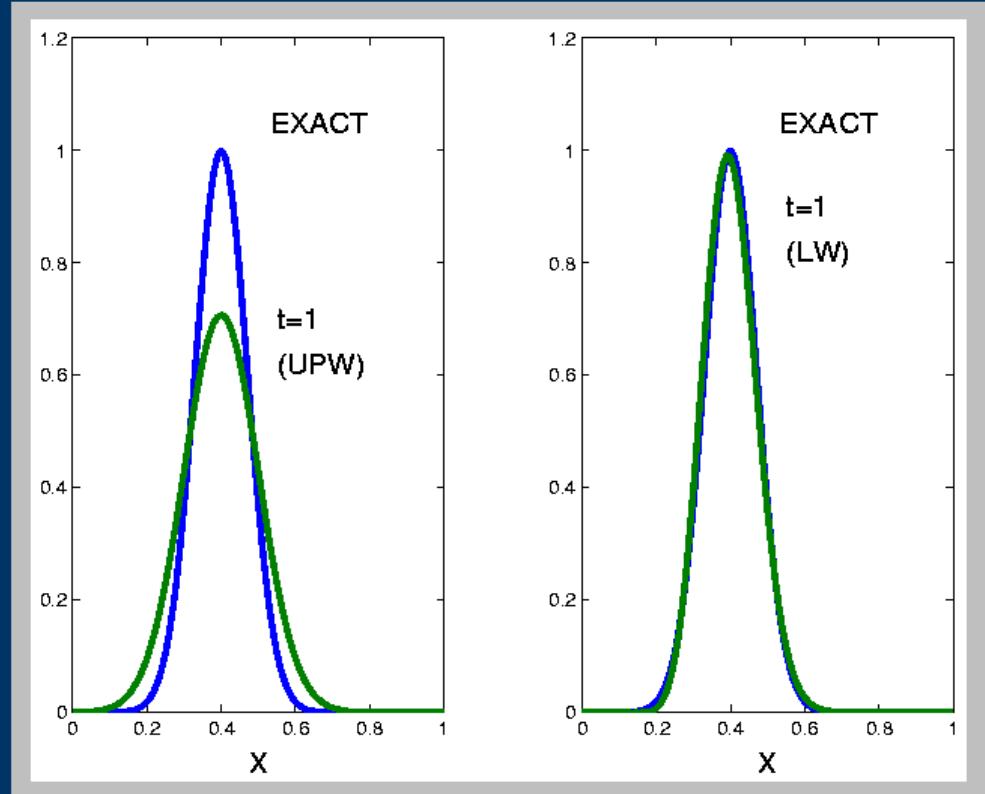
Lax-Wendroff Scheme

$\Delta x = 1/100$

$C = 0.5$

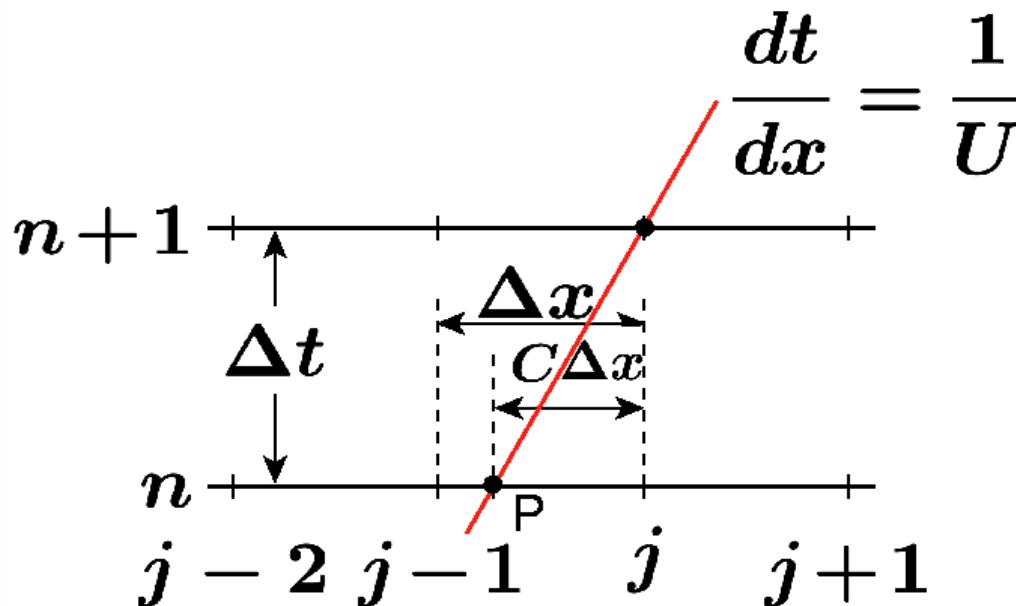
Upwind (left)
vs.
Lax-Wendroff (right)

Example



Beam-Warming Scheme

Derivation



$$u_j^{n+1} = u_P$$

Use Quadratic Interpolation
 $j-2, j-1, j$

$$u_P \approx -\frac{C}{2}(1-C)\hat{u}_{j-2}^n + C(2-C)\hat{u}_{j-1}^n + \frac{1}{2}(1-C)(2-C)\hat{u}_j^n$$

Beam-Warming Scheme

Consistency and Stability

$$\hat{u}_j^{n+1} = \hat{u}_j^n - \frac{C}{2}(3\hat{u}_j^n - 4\hat{u}_{j-1}^n + \hat{u}_{j-2}^n) + \frac{C^2}{2}(\hat{u}_j^n - 2\hat{u}_{j-1}^n + \hat{u}_{j-2}^n)$$

- Consistency, $\|\boldsymbol{\tau}\| \sim \mathcal{O}(\Delta x^2, \Delta t^2)$
- Stability

$$|g(C, \theta)|^2 = 1 - 4C(1-C)^2(2-C) \sin^4(\theta/2)$$

$$|g(C, \theta)| < 1 \quad \Rightarrow \quad \boxed{0 \leq C \leq 2}$$

Method of Lines

Generally applicable to time evolution PDE's

- **Spatial discretization**

- ⇒ Semi-discrete scheme (system of coupled ODE's)

- **Time discretization** (using ODE techniques)

- ⇒ Discrete scheme

By studying the semi-discrete scheme we can better understand spatial and temporal discretization errors

Method of Lines

NOTATION:

- $\bar{v}_j(t)$ approximation to $v(x_j, t) \equiv v_j(t)$
- $\bar{v}(t)$ vector of semi-discrete approximations;

$$\bar{v}(t) = \{\bar{v}_j(t)\}_{j=1}^J$$

Spatial Discretization

Method of Lines

$$\frac{\partial u}{\partial t} + U \frac{\partial u}{\partial x} = 0$$

Central differences... (for example)

$$\frac{d\bar{u}_j}{dt} + \frac{U}{2\Delta x} (\bar{u}_{j+1} - \bar{u}_{j-1}) = 0, \quad 1 \leq j \leq J$$

or, in vector form,

$$\frac{d\bar{u}}{dt} + \frac{U}{2\Delta x} \delta_{2x}\bar{u} = 0$$

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Method of Lines

Spatial Discretization

Fourier Analysis...

Write semi-discrete approximation as

$$\bar{u}_j(t) = \sum_{\theta = -\pi + 2\pi\Delta x}^{\pi} \bar{U}_{\theta}(t) e^{ij\theta}$$

Inserting into semi-discrete equation

$$\sum_{\theta} \left(\frac{d \bar{U}_{\theta}}{dt} + i \frac{U}{\Delta x} \sin(\theta) \bar{U}_{\theta} \right) e^{ij\theta} = 0, \quad 1 \leq j \leq J$$

Method of Lines

Spatial Discretization

...Fourier Analysis...

For each θ , we have a **scalar** ODE

$$\frac{d \bar{U}_\theta}{dt} + i \frac{U}{\Delta x} \sin(\theta) \bar{U}_\theta = 0$$

$$\Rightarrow \bar{U}_\theta(t) = \bar{U}_\theta^0 e^{-i \frac{U}{\Delta x} \sin(\theta) t}$$

$$|\bar{U}_\theta(t)| = |\bar{U}_\theta^0| \quad \text{Neutrally stable}$$

Method of Lines

Spatial Discretization

...Fourier Analysis...

Exact solution

$$u_j(t) = \sum_k \bar{U}_k^0 e^{i2\pi(kx_j - kU)t}$$

$$\omega_{EX} = kU$$

Semi-discrete solution

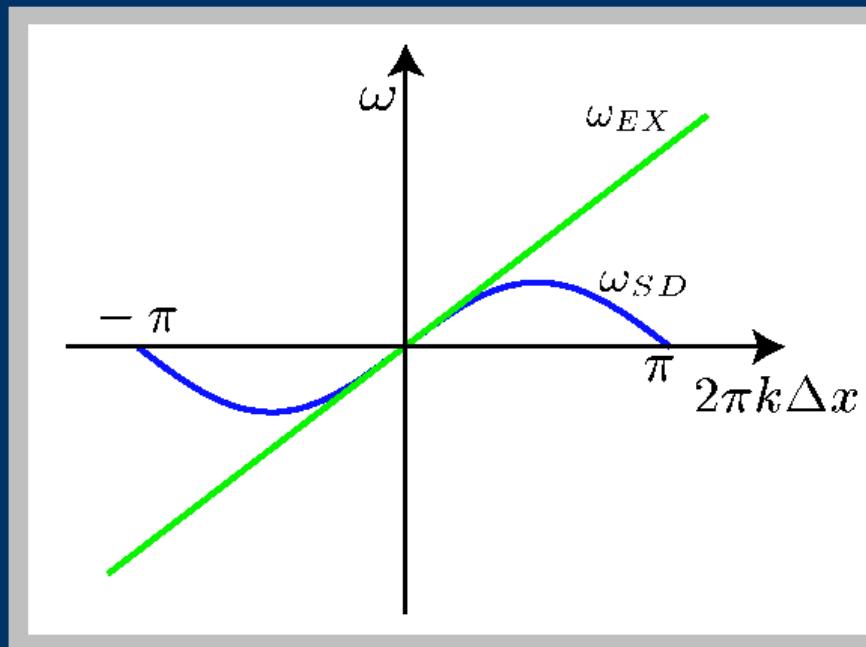
$$\begin{aligned}\bar{u}_j(t) &= \sum_\theta \bar{U}_\theta^0 e^{ij\theta} e^{-i\frac{U}{\Delta x} \sin(\theta)t} \\ &= \sum_k \bar{U}_k^0 e^{i2\pi(kx_j - \frac{U}{2\pi\Delta x} \sin(2\pi k \Delta x) t)}\end{aligned}$$

$$\omega_{SD} = \frac{U}{2\pi\Delta x} \sin(2\pi k \Delta x)$$

Method of Lines

Spatial Discretization

...Fourier Analysis



$$\omega_{EX} = kU$$

vs.

$$\omega_{SD} = \frac{U}{2\pi\Delta x} \sin(2\pi k\Delta x)$$

Method of Lines

Time Discretization

Predictor/Corrector Algorithm...

Model ODE

$$\frac{du}{dt} = \lambda u$$

$$\begin{aligned}\hat{u}^p &= \hat{u}^n + \Delta t \lambda \hat{u}^n \\ \hat{u}^{n+1} &= \hat{u}^n + \Delta t \lambda \hat{u}^p\end{aligned}$$

Predictor
Corrector

Combining the two steps we have

$$\hat{u}^{n+1} = \hat{u}^n + \Delta t \lambda \hat{u}^n + \Delta t^2 \lambda^2 \hat{u}^n = (1 + z + z^2) \hat{u}^n$$

$$z = \Delta t \lambda$$

Method of Lines

Time Discretization

...Predictor/Corrector Algorithm

Semi-discrete equation

$$\frac{d\underline{\mathbf{u}}}{dt} + \frac{U}{2\Delta x} \delta_{2x} \underline{\mathbf{u}} = 0$$

$$\hat{\underline{\mathbf{u}}}^p = \underline{\mathbf{u}}^n + \frac{C}{2} \delta_{2x} \hat{\underline{\mathbf{u}}}^n \quad \text{Predictor}$$

$$\underline{\mathbf{u}}^{n+1} = \underline{\mathbf{u}}^n + \frac{C}{2} \delta_{2x} \hat{\underline{\mathbf{u}}}^p \quad \text{Corrector}$$

Combining the two steps we have

$$\underline{\mathbf{u}}^{n+1} = \underline{\mathbf{u}}^n + \frac{C}{2} \delta_{2x} \hat{\underline{\mathbf{u}}}^n + \frac{C^2}{4} \delta_{2x}^2 \hat{\underline{\mathbf{u}}}^n$$

Method of Lines

$$\hat{\underline{u}}^{n+1} = \hat{\underline{u}}^n + \frac{C}{2} \delta_{2x} \hat{\underline{u}}^n + \frac{C^2}{4} \delta_{2x}^2 \hat{\underline{u}}^n$$

Fourier transform



$$\begin{aligned}\hat{\mathbb{U}}_\theta^{n+1} &= \hat{\mathbb{U}}_\theta^n - iC \sin(\theta) \hat{\mathbb{U}}_\theta^n - C^2 \sin^2(\theta) \hat{\mathbb{U}}_\theta^n \\ &= (1 + z_\theta + z_\theta^2) \hat{\mathbb{U}}_\theta^n, \quad \forall \theta\end{aligned}$$

$$z_\theta = -iC \sin(\theta)$$

Method of Lines

Amplification factor

$$g(C, \theta) = 1 + z_\theta + z_\theta^2$$

$z_\theta = i\alpha_\theta$ with $\alpha_\theta \in \mathbb{R}$

$$|g(C, \theta)|^2 = (1 - \alpha_\theta^2)^2 + \alpha_\theta^2 = 1 - \alpha_\theta^2(1 - \alpha_\theta^2)$$

Stability $\Rightarrow \alpha_\theta^2 \leq 1 \quad \forall \theta \quad \Rightarrow \quad C \leq 1$

Fourier Stability Analysis

Method of Lines

PDE $u(x, t)$



Semi-discrete
 $\bar{u}_j(t)$



A

Discrete
 \hat{u}_j^n



Semi-discrete Fourier
 $\bar{\mathbb{U}}_\theta(t)$

B

→

Discrete Fourier
 $\hat{\mathbb{U}}_\theta^n$

Fourier Stability Analysis

Method of Lines

Path B...

Semi-discrete

$$\frac{d\bar{u}}{dt} + \frac{U}{2\Delta x} \delta_{2x} \bar{u} = 0$$

Fourier semi-discrete

$$\frac{d\bar{\mathbb{U}}_\theta}{dt} + i \frac{U}{\Delta x} \sin(\theta) \bar{\mathbb{U}}_\theta = 0$$

Predictor

$$\hat{\mathbb{U}}^p = \hat{\mathbb{U}}^n - iC \sin(\theta) \hat{\mathbb{U}}^n$$

Corrector

$$\hat{\mathbb{U}}^{n+1} = \hat{\mathbb{U}}^n - iC \sin(\theta) \hat{\mathbb{U}}^p$$

Discrete

$$\hat{\mathbb{U}}_\theta^{n+1} = (1 + z_\theta + z_\theta^2) \hat{\mathbb{U}}_\theta^n$$

Method of Lines

...Path B

- Gives the same discrete Fourier equation
- Simpler
- “Decouples” spatial and temporal discretizations
 - For each θ , the discrete Fourier equation is the result of discretizing the **scalar** semi-discrete **ODE** for the θ Fourier mode

Method of Lines

Model Equation:

$$\frac{du}{dt} = \lambda u$$

u, λ complex-valued

Discretization

$$\frac{\hat{u}^{n+1} - \hat{u}^n}{\Delta t} = \lambda \hat{u}^n \quad \text{EF}$$

$$\frac{\hat{u}^{n+1} - \hat{u}^n}{\Delta t} = \lambda \hat{u}^{n+1} \quad \text{EB}$$

$$\frac{\hat{u}^{n+1} - \hat{u}^n}{\Delta t} = \frac{1}{2} \lambda (\hat{u}^n + \hat{u}^{n+1}) \quad \text{CN}$$

Methods for ODE's

Method of Lines

Absolute Stability Diagrams...

Given $\frac{du}{dt} = \lambda u$ and

u, λ complex-valued

$$\frac{\hat{u}^{n+1} - \hat{u}^n}{\Delta t} = \lambda \hat{u}^n \text{ (EF) or } \lambda \hat{u}^{n+1} \text{ (EB) or ...;}$$

\mathcal{R}_{EF}^{abs} or ... $\in \mathbb{C}$ is defined such that

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$$z \equiv \Delta t \lambda \in \mathcal{R}_{EF}^{abs} \Leftrightarrow |\hat{u}^{n+1}| < |\hat{u}^n|$$

$$\Rightarrow |\hat{u}^n| \rightarrow 0 \text{ as } n \rightarrow \infty$$

Methods for ODE's

Method of Lines

...Absolute Stability Diagrams...

$$\hat{u}^{n+1} - \hat{u}^n = \Delta t \lambda \hat{u}^n \quad \text{EF}$$

$$\Rightarrow \quad \hat{u}^{n+1} = (1 + z) \hat{u}^n$$

$$\hat{u}^{n+1} - \hat{u}^n = \Delta t \lambda \hat{u}^{n+1} \quad \text{EB}$$

$$\Rightarrow \quad \hat{u}^{n+1} = \frac{1}{1-z} \hat{u}^n$$

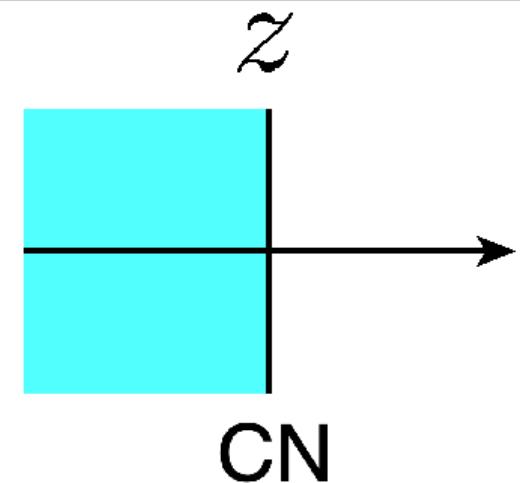
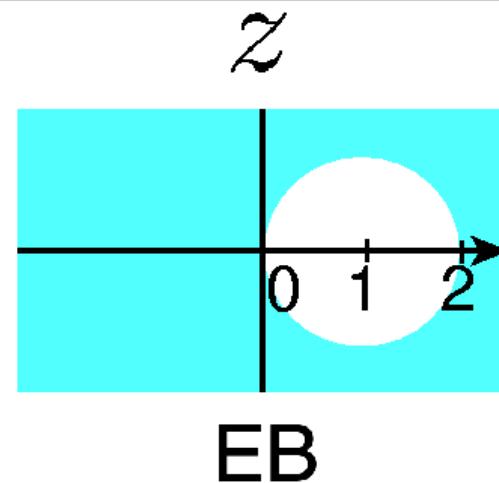
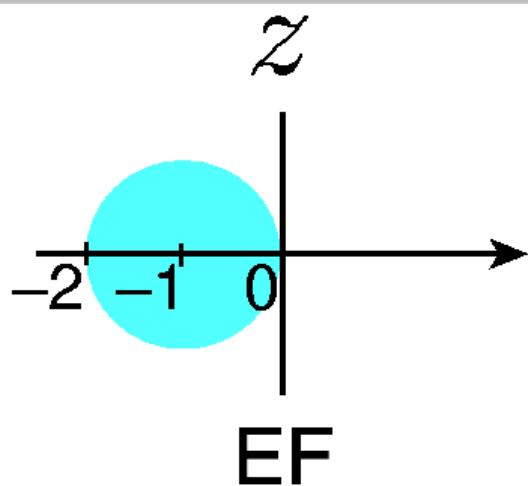
$$\hat{u}^{n+1} - \hat{u}^n = \frac{1}{2} \Delta t \lambda (\hat{u}^n + \hat{u}^{n+1}) \quad \text{CN}$$

$$\Rightarrow \quad \hat{u}^{n+1} = \frac{1+z/2}{1-z/2} \hat{u}^n$$

Method of Lines

Methods for ODE's

...Absolute Stability Diagrams



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Method of Lines

Methods for ODE's

Application to the Wave Equation...

For each θ

$$\frac{d \bar{U}_\theta}{dt} + i \frac{U}{\Delta x} \sin(\theta) \bar{U}_\theta = 0, \quad \text{or} \quad \boxed{\frac{d \bar{U}_\theta}{dt} = \lambda_\theta \bar{U}_\theta}$$

Thus,

$$\lambda_\theta = -i \frac{U}{\Delta x} \sin(\theta)$$

- λ_θ (and $z_\theta = \Delta t \lambda_\theta$) is purely imaginary
- $\lambda_\theta \rightarrow \infty$ for $\Delta x \rightarrow 0$

Method of Lines

Methods for ODE's

...Application to the Wave Equation...

$$\frac{d \bar{U}_\theta}{dt} = \lambda_\theta \bar{U}_\theta$$

- ⇒ **EF** is unconditionally **unstable**
- ⇒ **EB** is unconditionally **stable**
- ⇒ **CN** is unconditionally **stable**

Method of Lines

Methods for ODE's

...Application to the Wave Equation...

Stable schemes can be obtained by:

- 1) Selecting explicit time stepping algorithms which have some stability on the imaginary axis
- 2) Modifying the original equation by adding “artificial viscosity” $\Rightarrow \Re(\lambda_\theta) < 0$

Method of Lines

Methods for ODE's

...Application to the Wave Equation...

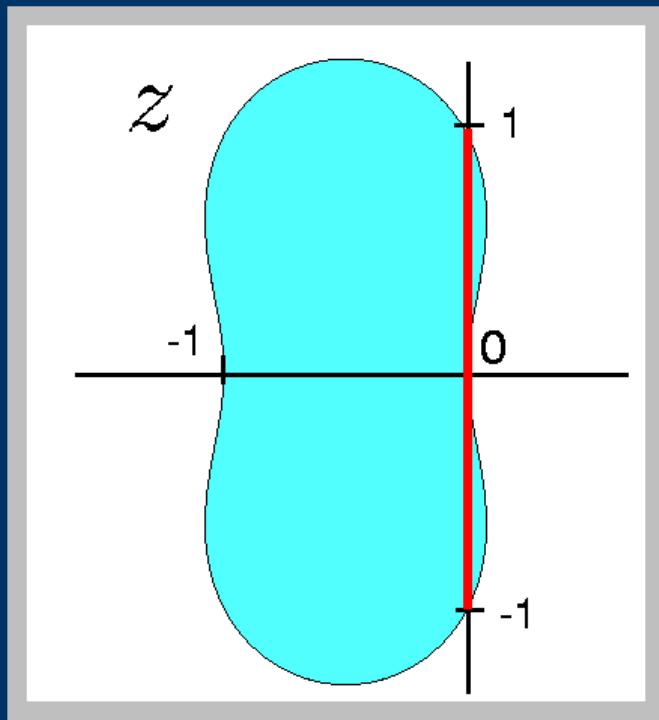
Explicit Time stepping Schemes

Predictor/Corrector

$$\hat{u}^{n+1} = (1 + z + z^2)\hat{u}^n$$

$$z_\theta = iC \sin(\theta)$$

$$\Rightarrow C \leq 1$$



Method of Lines

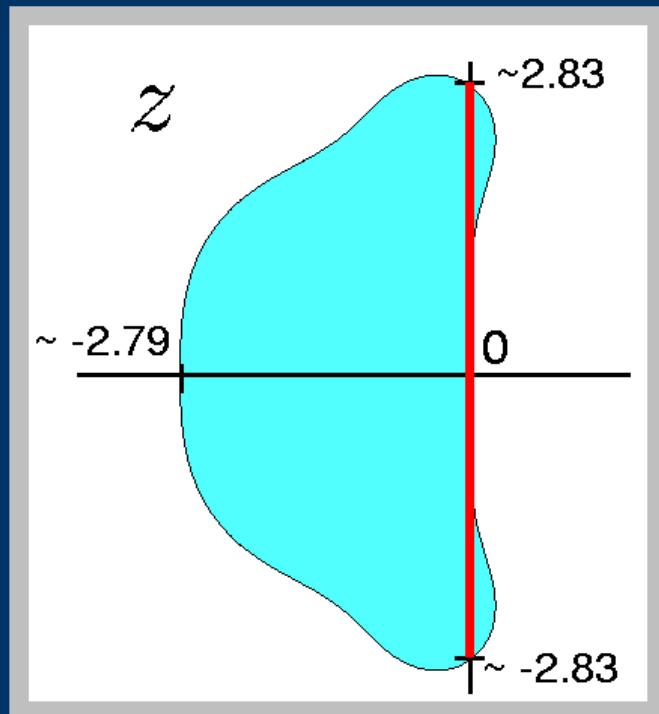
Methods for ODE's

...Application to the Wave Equation...

Explicit Time stepping Schemes

4 Stage Runge-Kutta

$$\hat{u}^{n+1} = \left(1 + z + \frac{z^2}{2} + \frac{z^3}{6} + \frac{z^4}{24}\right) \hat{u}^n$$



$$z_\theta = iC \sin(\theta)$$

$$\Rightarrow C \leq 2\sqrt{2} \sim 2.83$$

Method of Lines

Methods for ODE's

...Application to the Wave Equation...

Adding Artificial Viscosity

$$\frac{d\bar{u}}{dt} + \frac{U}{2\Delta x} \delta_{2x} \bar{u} - \underbrace{\mu \frac{U}{2\Delta x} \delta_x^2 \bar{u}}_{\text{Additional Term}} = 0$$

EF Time + $\mu = 1$ \Rightarrow First Order Upwind

EF Time + $\mu = C$ \Rightarrow Lax-Wendroff

Method of Lines

Methods for ODE's

...Application to the Wave Equation...

Adding Artificial Viscosity

For each Fourier mode θ ,

$$\frac{d \bar{U}_\theta}{dt} + \left\{ i \frac{U}{\Delta x} \sin(\theta) - 2\mu \underbrace{\frac{U}{\Delta x} \sin^2(\theta/2)}_{\text{Additional Term}} \right\} \bar{U}_\theta = 0$$

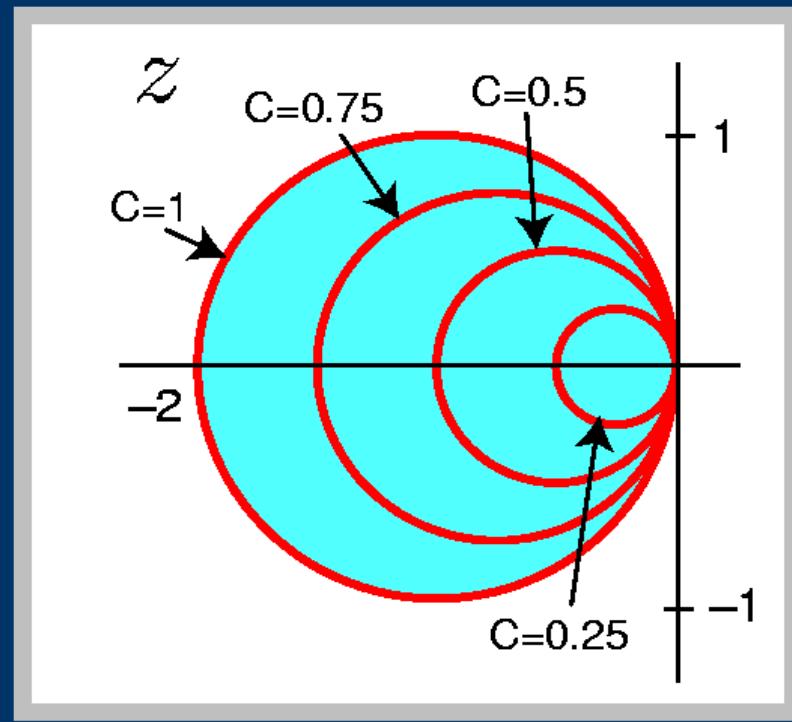
$$z_\theta = -2\mu C \sin^2(\theta/2) - iC \sin(\theta)$$

Method of Lines

Methods for ODE's

...Application to the Wave Equation...

First Order Upwind Scheme $\mu = 1$

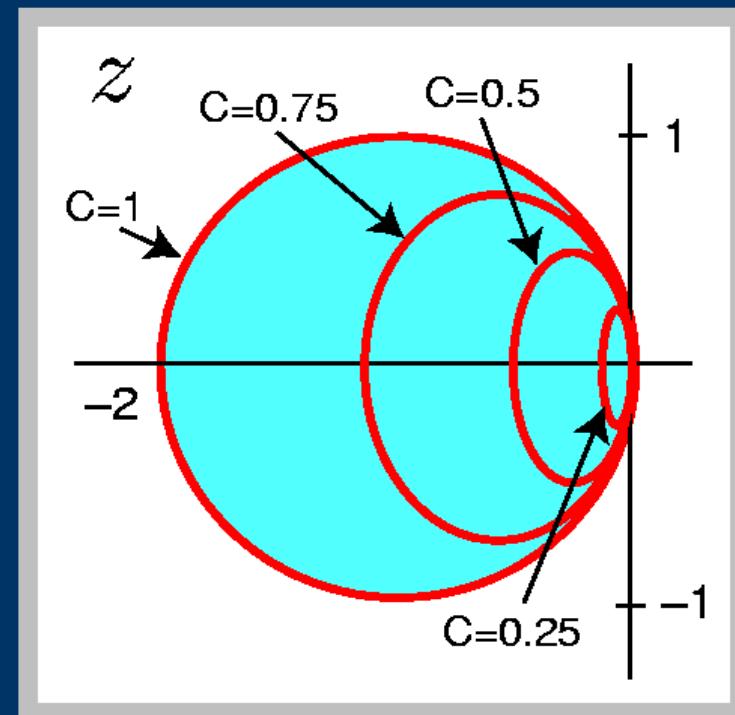


Method of Lines

Methods for ODE's

...Application to the Wave Equation

Lax-Wendroff Scheme $\mu = C$



Dissipation and Dispersion

Model Problem

$$\frac{\partial u}{\partial t} + U \frac{\partial u}{\partial x} = \kappa \frac{\partial^2 u}{\partial x^2} - a \frac{\partial^3 u}{\partial x^3}, \quad x \in (0, 1)$$

with $u(x, 0) = u^0(x)$ and periodic boundary conditions.

Solution

$$u(x, t) = \sum_{k=-\infty}^{k=\infty} \mathbb{U}_k^0 e^{-4\pi^2 \sigma(k)t} e^{i2\pi(kx - \omega(k)t)}$$

$$\sigma(k) = \kappa k^2, \quad \omega(k) = U k - a 4\pi^2 k^3$$

Dissipation and Dispersion

Model Problem

$e^{-4\pi^2\sigma(k)t}$ represents **Decay**
 $\sigma(k)$ dissipation relation

$e^{i2\pi(kx-\omega(k)t)}$ represents **Propagation**
 $\omega(k)$ dispersion relation

For the exact solution of $\mathbf{u}_t + \mathbf{U}\mathbf{u}_x = \mathbf{0}$

$\sigma = 0$ no dissipation

$\omega = kU$, or $\omega/k = U$ (constant) no dispersion

Dissipation and Dispersion

Modified Equation

First Order Upwind

$$u_t + U u_x = \frac{U \Delta x}{2} (1 - C) u_{xx} - \frac{U \Delta x^2}{6} (1 - C^2) u_{xxx}$$

Lax-Wendroff

$$u_t + U u_x = -\frac{U \Delta x^2}{6} (1 - C^2) u_{xxx}$$

Beam-Warming

$$u_t + U u_x = \frac{U \Delta x^2}{6} (2 - C)(1 - C) u_{xxx}$$

- For the upwind scheme **dissipation** dominates over dispersion \Rightarrow **Smooth** solutions
- For Lax-Wendroff and Beam-Warming **dispersion** is the leading error effect \Rightarrow **Oscillatory** solutions (if not well resolved)
- Lax-Wendroff has a **negative** phase error
- Beam-Warming has (for $C < 1$) a **positive** phase error

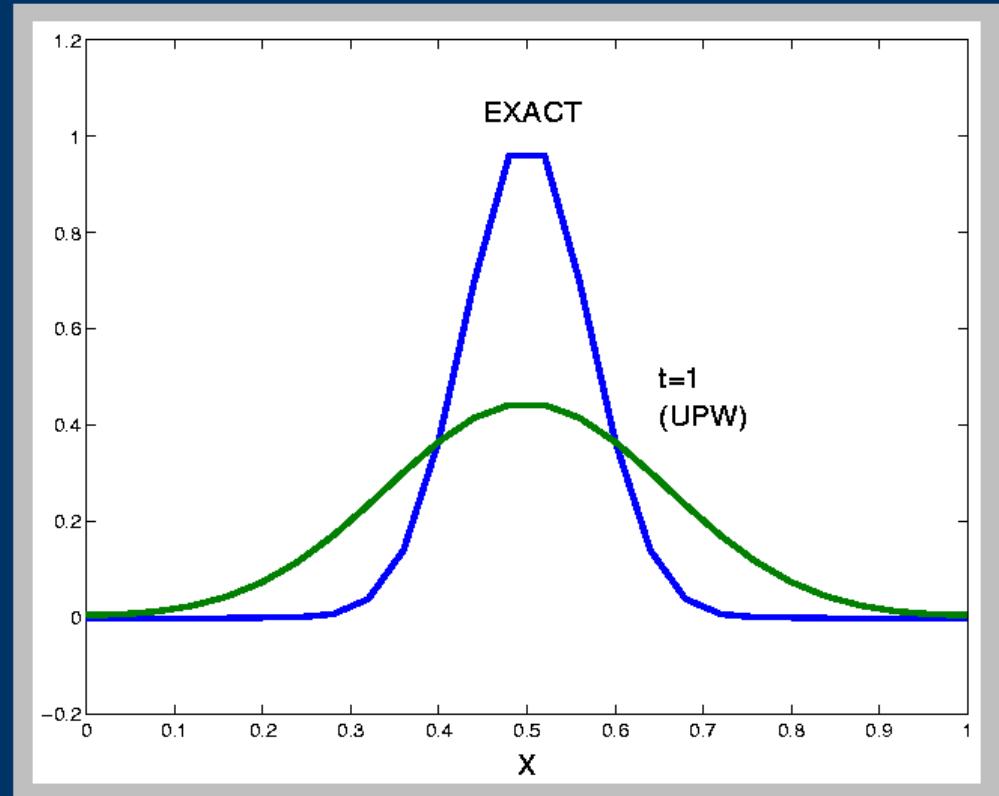
Dissipation and Dispersion

Examples

$$\Delta x = 1/25$$

$$C = 0.5$$

First Order Upwind



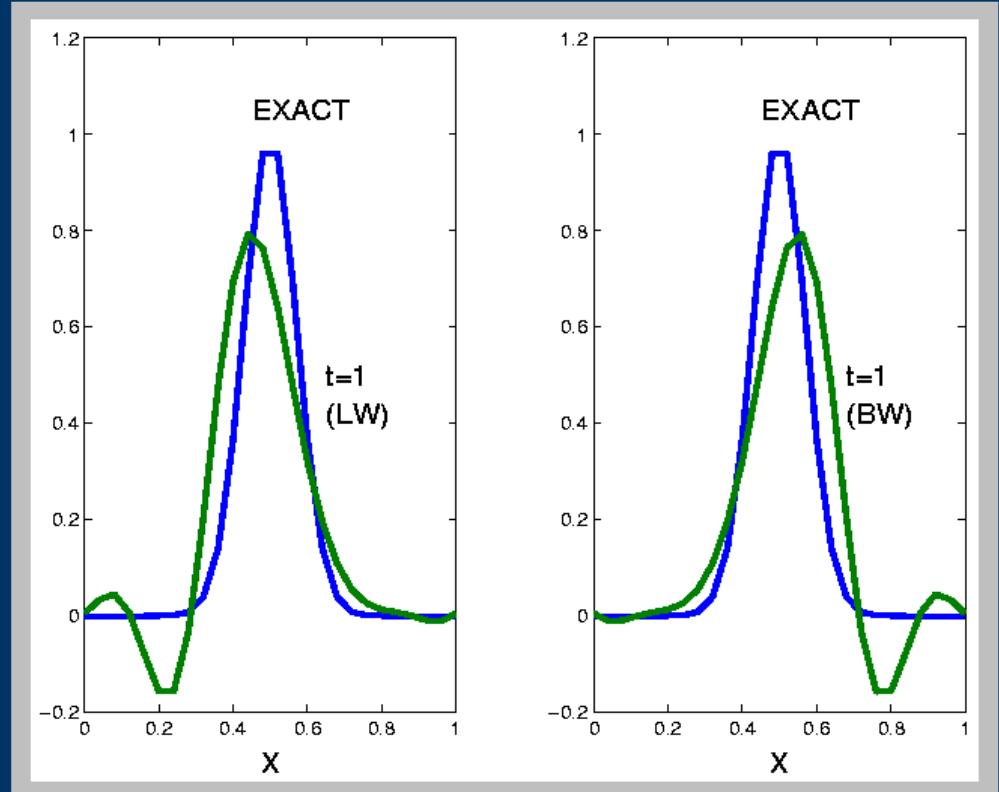
Dissipation and Dispersion

Examples

$$\Delta x = 1/25$$

$$C = 0.5$$

Lax-Wendroff (left)
vs.
Beam-Warming (right)



Dissipation and Dispersion

Exact Discrete Relations

For the exact solution

$$\mathbb{U}_\theta^{n+1} = e^{i2\pi k U \Delta t} \mathbb{U}_\theta^n$$

$$\Rightarrow \omega_{EX} = kU = \theta U / 2\pi \Delta x, \text{ and } \sigma_{EX} = 0$$

For the discrete solution

$$\hat{\mathbb{U}}_\theta^{n+1} = g(C, \theta) \hat{\mathbb{U}}_\theta^n$$

$$g(C, \theta) = e^{-i2\pi\omega(\theta)\Delta t - 4\pi^2\sigma(\theta)\Delta t}$$
$$\Rightarrow \omega(\theta), \text{ and } \sigma(\theta)$$