

Lecture # 22

① ✓

11/4/05

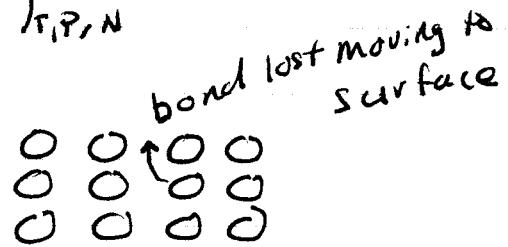
Last time

Surface tension - free energy cost of creating new surface

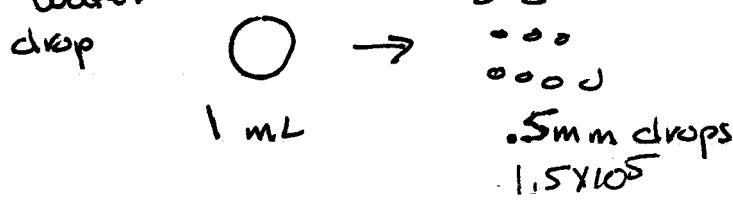
$$\gamma = \frac{\partial A}{\partial A})_{T,V,N}$$

$$\gamma = \frac{\partial G}{\partial A})_{T,P,N}$$

$$\gamma = - \frac{w_{AA}}{2a}$$



How much energy does it take for water drop



$$\Delta G = \gamma \Delta A = \gamma (A\pi) [1.5 \times 10^5 (0.025)^2 - 0.62^2]$$

$$= 734 - 4.8 = 729 \text{ cm}^2 = .073 \text{ m}^2$$

$$\gamma = 72.8 \text{ dyne/cm} = 0.0728 \text{ kg/s}^2$$

$$\Delta G = (.0728)(.073) = 5.34 \times 10^{-3} \text{ J}$$

Getting to be the time of year - what temp does sea water freeze? Why put salt on ice?

⇒ Dirty water freezes at a lower temp than pure water. (boils higher, also)
Let's now predict "colligative property"

Using the lattice model, we derived ^{for} regular solutions, this relationship

$$M_A = kT \left[\frac{\partial}{\partial N_A} \left(\frac{F}{kT} \right) \right]_{N_A, T}$$

$$= kT \left[\ln x_A + \frac{z_{WAA}}{kT} + x_{A_B} (1-x_A)^2 \right]$$

Microscopic $M_A = \frac{z_{WAA}}{2} + kT \left[\ln(x_A) + x_{A_B} (1-x_A)^2 \right]$

We could rewrite as

Microscopic $M_A = \frac{z_{WAA}}{2} + kT \ln \left(e^{x_{A_B}(1-x_A)^2} x_A \right)$

leads to an obvious definition of macroscopic M_A

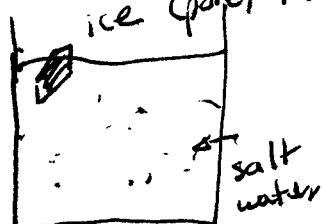
Macroscopic $M_A = M_A^o + kT \ln \gamma_A x_A$ Eqⁿ 15.17

↑ ↴
pure A activity coeff captures nonidealities
 $\gamma_A \rightarrow 1$ as $x_A \rightarrow \phi$

Book analyzes boiling \Rightarrow that is for summer version of course.

Now lets analyze freezing salt water \Rightarrow

similar approach to other phase transitions.



water = B

salt = A

$$\begin{array}{ccc} \text{pure} & & \text{dirty} \\ M_B(\text{ice}, T) = M_B(\text{l. q}, T, x_B) & & \\ \downarrow & & \downarrow \\ M_B^\circ(\text{ice}) = M_B^\circ(\text{l. q}) + KT \ln \gamma_B x_B & & \\ & \text{pure}_B & \end{array}$$

Eq ⁿ for boiling	16.16
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$$\ln \gamma_B x_B = \frac{1}{KT} [M_B^\circ(\text{ice}, T) - M_B^\circ(\text{l. q}, T)]$$

Recall $M_j = \frac{\partial G}{\partial N_j} \Big|_{T, P, N_{i \neq j}}$

$$M_j = \frac{\partial H}{\partial N_j} - T \frac{\partial S}{\partial N_j} = h_j - TS_j \quad \text{eqn 9.25}$$

Sub in

$$\ln \gamma_B x_B = \frac{1}{KT} [(h^\circ_{\text{ice}} - h^\circ_{\text{l. q}}) - \cancel{T(S^\circ_{\text{ice}} - S^\circ_{\text{l. q}})}]$$

Now for pure H₂O @ T = T_f ← freezing temp

$$\ln(1) = \frac{1}{KT_f} (h^\circ_{\text{ice}} - h^\circ_{\text{l. q}}) - \frac{1}{K} (S^\circ_{\text{ice}} - S^\circ_{\text{l. q}})$$

Subtract $\ln \gamma_B x_B - 0 = \left(\frac{1}{KT} - \frac{1}{KT_f} \right) \underbrace{(h^\circ_{\text{ice}} - h^\circ_{\text{l. q}})}_{\Delta h_f} - \phi$

$$\ln \gamma_B x_B = \frac{\Delta h_f (T_f - T)}{KT_f}$$

$$\boxed{\begin{aligned} T &\sim T_f \\ \frac{1}{T} - \frac{1}{T_f} &= \frac{T_f - T}{T_f T} \sim \frac{T_f - T}{T_f^2} \end{aligned}}$$

Simpl. for

For small solute concentrations

$$\begin{aligned}\ln \gamma_B x_B &= \ln [\gamma_B (1-x_A)] = \ln \gamma_B + \ln (1-x_A) \\ &\quad \text{from micro def} \\ &= \chi_{AB} x_A^2 + (-x_A - \frac{x_A^2}{2} - \frac{x_A^3}{3} \dots) \\ &= -[x_A + (\frac{1}{2} - \chi_{AB}) x_A^2 + \frac{x_A^3}{3} \dots]\end{aligned}$$

2nd order approx

$$\Delta T = T - T_f = \frac{k T_f^2}{\Delta h_f^\circ} [x_A + (\gamma_2 - \chi_{AB}) x_A^2]$$

and for $x_A \ll 1$

$$\Delta T = \frac{RT_f^2 x_A}{\Delta h_f^\circ} \text{ per mol}$$

Does not depend
on nature of
the solute?
only x_A

For Seawater

$$x_A \sim 0.035$$

$$T_f = 273 K$$

$$\Delta h_f^\circ = 6001 J/mol$$

$$R = 8.315 J/mol \cdot K$$

$$\Delta T_f = \frac{8.315 (273)^2}{6001} (0.035) = 3.6 K$$

Can use FPD to get M_w of an unknown dilute:

$$x_A = m_A / M_B \leftarrow \text{molar mass of solvent}$$

\uparrow
moles A / 1000 kg solvent

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