

MASSACHUSETTS INSTITUTE OF TECHNOLOGY

**5.74 Quantum Mechanics II**  
**Spring, 2004**

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*Problem Set #7*

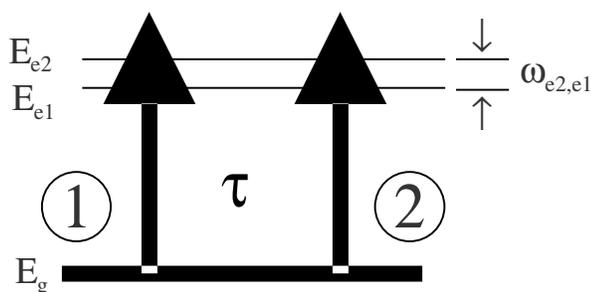
**DUE:** At the start of Lecture on Friday, April 23.

**Reading:** HLB-RWF 9.1.7, 9.1.8, 9.4.9, 9.2.1, 9.2.2, 9.2.3, 9.2.4

**Problems:**

I. One-Color Pump-Probe Experiment.

Consider a three-level system.



Pulses (1) and (2) are obtained by splitting a 100fs laser pulse into two identical pulses and systematically delaying pulse (2) relative to pulse (1) using a delay line.

The intensity of pulse (1) is set so that, immediately after pulse (1) has exited the sample

$$\Psi(0) = [1 - 0.02]^{1/2} \psi_g + 0.1\psi_{e1} + 0.1\psi_{e2}.$$

A. For excitation by pulse (1) only, what does  $P(t) = |\langle \Psi(t) | \Psi(0) \rangle|^2$  look like?

- B. For excitation at  $t = 0$  by both pulses (1) and (2), with zero delay between the two pulses, what is  $\Psi(0)$ ? Be extremely careful to justify your value for the coefficients in front of  $\psi_{e1}$  and  $\psi_{e2}$ .
- C. For excitation at  $t = 0$  by pulse (1) and  $t = \tau$  by pulse (2), what is  $\Psi(\tau)$ ?
- D. Compute  $P(t - \tau) = |\langle \Psi(t - \tau) | \Psi(\tau) \rangle|^2$  for  $t = 0$ ,  $\tau = \frac{2\pi}{\omega_{e2,e1}}(1/4)$ , and  $\tau = \frac{2\pi}{\omega_{e2,e1}}(1/2)$ .
- E. Suppose you were going to monitor  $P(t - \tau)$  in a fluorescence-dip scheme. Describe this one-color pump-probe fluorescence-dip experiment in the  $\rho(0)$ , **E**, **U**, **D** formulation.
- F. Consider an experiment in which you have a cw laser oscillating at  $\left(\frac{E_{e2} + E_{e1}}{2} - E_g\right)/\hbar = \omega_{cw}$ . Your excitation pulse at  $t = 0$  is obtained by pulse amplifying  $\omega_{cw}$  from  $\sim 1\text{mW}$  cw to a 100fs,  $1\mu\text{J}$  pulse (peak power  $10^7\text{W}$ ). The laser beam traverses your sample and impinges on a “square law” detector. The signal is

$$I(t) \propto |\epsilon_{\text{laser}}(t) + \epsilon_{\text{molecule}}(t)|^2.$$

Use the  $\rho$ , **E**, **U**, **D** formulation to describe  $I(t)$  at  $t > 0$  for a single amplified pulse at  $t = 0$ .

## II. Anharmonic Coupling

$$\mathbf{H} = \mathbf{h}_1 + \mathbf{h}_2 + \mathbf{h}_{12}$$

$$\mathbf{h}_1/hc = (\omega_1/2)(\mathbf{a}_1^\dagger \mathbf{a}_1 + \mathbf{a}_1 \mathbf{a}_1^\dagger) + (x_{11}/4)(\mathbf{a}_1^\dagger \mathbf{a}_1 + \mathbf{a}_1 \mathbf{a}_1^\dagger)^2$$

$$\mathbf{h}_2/hc = (\omega_2/2)(\mathbf{a}_2^\dagger \mathbf{a}_2 + \mathbf{a}_2 \mathbf{a}_2^\dagger) + (x_{22}/4)(\mathbf{a}_2^\dagger \mathbf{a}_2 + \mathbf{a}_2 \mathbf{a}_2^\dagger)^2$$

$$\mathbf{h}_{12}/hc = k_{122}(\mathbf{a}_1^\dagger + \mathbf{a}_1)(\mathbf{a}_2^\dagger + \mathbf{a}_2)^2$$

where  $\omega_1$ ,  $x_{11}$ ,  $\omega_2$ ,  $x_{22}$ , and  $k_{122}$  are in  $\text{cm}^{-1}$  units.

The goal of this problem is to discover how (or whether) it is possible to determine experimentally the sign of  $k_{122}$ . This sign is physically significant because, if  $k_{122} > 0$ , then when bond 1 ( $\mathbf{Q}_1$ ) is stretched the oscillation frequency of bond 2 is *increased*, and, if  $k_{122} < 0$ , stretching bond 1 *decreases* the oscillation frequency of bond 2.

The energy levels of this two-mode oscillator are described by an  $\mathbf{H}^{\text{eff}}$  that has diagonal elements

$$E^{(0)}(v_1, v_2)/hc = \omega_1(v_1 + 1/2) + \omega_2(v_2 + 1/2) + x_{11}(v_1 + 1/2)^2 + x_{22}(v_2 + 1/2)^2 + x_{12}(v_1 + 1/2)(v_2 + 1/2)$$

and, because  $\omega_1 \approx 2\omega_2$ , is arranged in quasi-degenerate polyad blocks, each block with polyad quantum number  $P = 2v_1 + v_2$ . The off-diagonal matrix elements within each polyad block are

$$\mathbf{\Omega} + \mathbf{\Omega}^\dagger \equiv k_{122}[\mathbf{a}_1 \mathbf{a}_2^\dagger \mathbf{a}_2^\dagger + \mathbf{a}_1^\dagger \mathbf{a}_2 \mathbf{a}_2]$$

and the nonzero elements follow the selection rule  $2\Delta v_1 = -\Delta v_2$ .

- A. Set up the  $P = 10$  polyad block of  $\mathbf{H}$ . This is a  $6 \times 6$  matrix. Let  $\omega_1 = 1000 \text{ cm}^{-1}$ ,  $\omega_2 = 500 \text{ cm}^{-1}$ ,  $x_{11} = 10 \text{ cm}^{-1}$ ,  $x_{22} = 5 \text{ cm}^{-1}$ ,  $x_{12} = 0$ , and  $k_{122} = \pm 100 \text{ cm}^{-1}$ . Solve for the energy levels for  $k_{122} = 100 \text{ cm}^{-1}$  and  $k_{122} = -100 \text{ cm}^{-1}$ . Is the sign of  $k_{122}$  an observable quantity? If it is not,

you can skip part B, unless you have a clever idea for how to sample the sign of  $k_{122}$  in the time domain.

- B. Use the operator algebra from Lecture #8 (Resonance Operators: Equations of Motion) to propose a time domain observable quantity that is sensitive to the sign of  $k_{122}$ .
- C. Examine the time dependence of  $\langle \mathbf{a}_1^\dagger \mathbf{a}_1 \rangle_t$  for each of the six possible single basis-state plucks of the  $P = 10$  polyad. Which pluck is most sensitive to the value of  $k_{122}$ ?
- D. Examine the time dependence of the expectation value  $\langle \mathbf{\Omega} + \mathbf{\Omega}^\dagger \rangle_t$  for each of the six possible single basis-state plucks of the  $P = 10$  polyad. Which pluck is most sensitive to the value of  $k_{122}$ ?
- E. If you conclude that the sign of  $k_{122}$  is not an observable quantity, replace  $\mathbf{h}_{12}$  by

$$\mathbf{h}_{12} = k_{1122} (\mathbf{a}_1^\dagger + \mathbf{a}_1)^2 (\mathbf{a}_2^\dagger + \mathbf{a}_2)^2$$

and let  $\omega_1 = \omega_2 = 1000 \text{ cm}^{-1}$ ,  $x_{11} = x_{22} = 10 \text{ cm}^{-1}$ , and  $k_{1122} = \pm 100 \text{ cm}^{-1}$ .

Show that the sign of  $k_{1122}$  is an observable quantity.