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5.80 Small-Molecule Spectroscopy and Dynamics
Fall 2008

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Lecture #24: Pure Rotation Spectra of Polyatomic Molecules

Last time: Asymmetric tops

$$\mathbf{H}^{\text{ROT}} = \left[\bar{B}J(J+1) + \left(\frac{A-\bar{B}}{C-\bar{B}} \right) K^2 \right] |JK\rangle\langle JK| + \left(\frac{B-C}{A-B} \right) \frac{1}{4} [(J \mp K)(J \mp K - 1)(J \pm K + 2)(J \pm K + 1)]^{1/2} |JK\rangle\langle JK \pm 2|$$

asymmetry parameter $\kappa \equiv \frac{2B - A - C}{A - C}$

prolate	$B = C$	$\kappa = -1$
oblate	$A = B$	$\kappa = +1$
most asymmetric	$B = \frac{A+C}{2}$	$\kappa = 0$

Perturbation Theory: $\frac{H'_{JK,JK\pm 2}}{\Delta E^0} = \frac{1}{\kappa} f(J,K)$

Factor \mathbf{H}^{ROT} into four blocks — can label these blocks according to whether K_p and K_o (or K_a and K_c) are even or odd

e,e	e,o	o,e	o,o	correspond to rotational symmetry species — good quantum numbers [not related to geometric symmetry operations applied to rigid molecule].
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TODAY

pure rotation spectra

SYMMETRIC TOP $\rightarrow M_x = M_y = 0, \Delta K = 0$ like diatomic. Only get information about B from pure rotation spectrum.

ASYMMETRIC TOP $\rightarrow \bar{M}$ can have 1, 2, or 3 non-zero components

Transition	ΔK_p	ΔK_o
a-type	even	odd
b-type	odd	odd
c-type	odd	even

NOTATION/ASSIGNMENT

$$P_{if} \propto \left| \sum_{S,b} \epsilon_S \langle v_i | \mathbf{M}_{b,if}(\mathbf{Q}) | v_f \rangle \langle \Omega_i J_i M_i | \alpha_{Sb} | \Omega_f J_f M_f \rangle \right|^2$$

pure rotation
↓
 $M_b(Q_e)$

Hougen, p. 31
↓
 $f(J_i, J_f) g(J_i \Omega_i; J_f \Omega_f) h(J_i M_i; J_f M_f)$
body polarization lab polarization

BODY COMPONENT $b \rightarrow (a, b, \text{ or } c)$ $K \leftrightarrow \Omega$ Symmetric Tops - can't have nonzero dipole \perp to unique axis: z $z \rightarrow a$ prolate $z \rightarrow c$ oblate

"parallel" type transitions only

 $\Delta K = 0, \Delta J = \pm 1$ ($\Delta J = 0$ weak, except for $J = K$)if there were a $\perp \mu$ component, it would have to point in same direction after a $C_n(z)$ or $S_n(z)$ operation. $\mu_{\perp} = 0$ for $n \geq 2$.

$$E_{\text{ROT}} = BJ(J+1) + \begin{pmatrix} A-B \\ C-B \end{pmatrix} K^2$$

identical $BJ(J+1)$ stacks shifted by $\begin{pmatrix} A-B \\ C-B \end{pmatrix} K^2$

no inter-stack transitions

no information about other rotational constants

$$\Delta E_{JK \leftrightarrow J-1K} = 2BJ$$

(small splittings due to centrifugal distortion)

$$E_{JK} = E_{JK}^{\circ} - D_J J^2 (J+1)^2 - \underbrace{D_{JK} J(J+1)K^2}_{\text{centrifugal distortion}} - D_K K^4$$

but these give no information about A or C!

Get information about $\left\{ \begin{array}{l} \text{other rotational constants} \\ \text{equilibrium structure} \end{array} \right\}$ $\left\{ \begin{array}{l} \text{from IR and UV spectra} \\ \text{or from isotopomers} \end{array} \right\}$ $^{14}\text{NH}_3$ $^{15}\text{NH}_3$ $^{14}\text{ND}_3$ etc. and from non-symmetric species (non-symmetric tops!)Asymmetric Tops - more than 1 component of \overline{M} can be nonzero. $K \neq 0$ levels split by asymmetry splitting.Same $\Delta^K \Delta J$ type transitions from different K 's not superimposed.

No resemblance to diatomic molecule spectrum!

symmetry label - C_2 about specified axis

$$J_{K_p, K_o} \begin{pmatrix} \text{prolate axis} & \text{oblate axis} \\ \text{e or o} & \text{e or o} \\ K_p & K_o \end{pmatrix}$$

 M_a (e, o) M_b (o, o) M_c (o, e)

$$\left\langle J'_{K'_p K'_o} \left| M_{\begin{matrix} a \\ b \\ c \end{matrix}} \right| J''_{K''_p K''_o} \right\rangle \rightarrow \begin{pmatrix} e & e \\ o & o \end{pmatrix} \begin{pmatrix} M & \Delta K_p & \Delta K_c \\ a & e & o \\ b & o & o \\ c & o & e \end{pmatrix} \begin{pmatrix} K'_p & K''_o \\ e & e \\ o & o \end{pmatrix}$$

e × e × e = e o × o × e = e etc.

This is not matrix multiplication. It is just a summary of possibilities.

integrand must be (e, e) with respect to both K_p and K_o for integral to be $\neq 0$

a-type selection rule	$\Delta K_a = \text{even}$	$\Delta K_c = \text{odd}$
b-type selection rule	$\Delta K_a = \text{odd}$	$\Delta K_c = \text{odd}$
c-type selection rule	$\Delta K_a = \text{odd}$	$\Delta K_c = \text{even}$

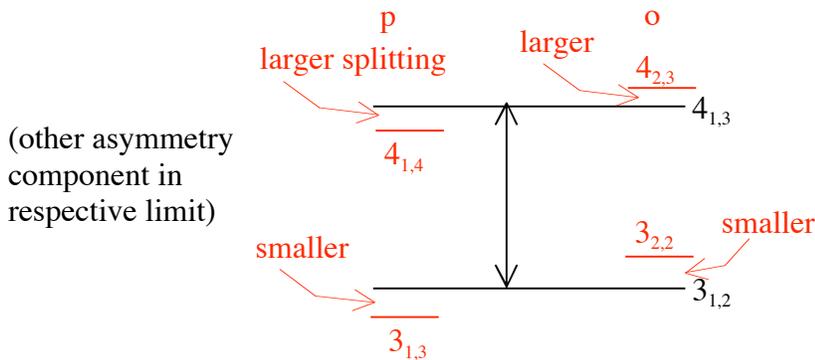
Want ΔK to be as small as possible. 0 or ± 1 for near symmetry axis, ± 2 or ± 3 possible for other axes.

K_p and K_o are (related to) symmetry labels. Only one of these two quantum numbers is an almost good projection quantum number.

$\Delta J = 0$ strong or weak depending on whether $\Delta K_{\text{near sym}} = \pm 1$ or 0.

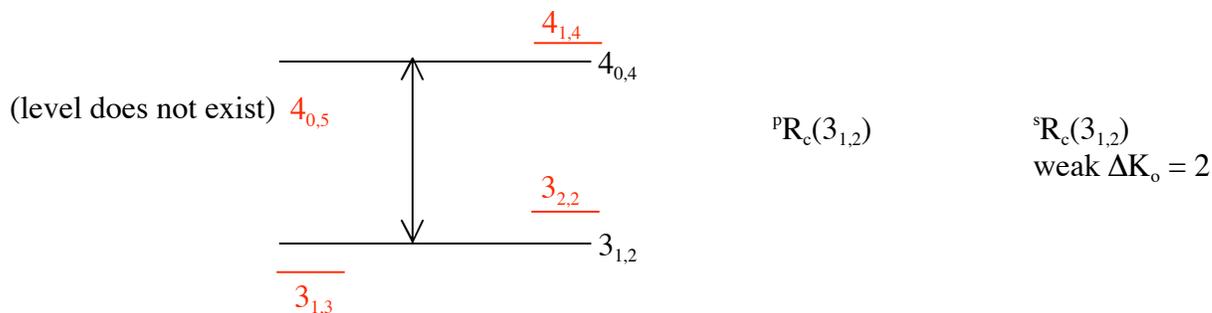
	near prolate		near oblate
a	$\Delta J = 0$ weak $\Delta K_a = 0$		\perp $\Delta J = 0$ strong $\Delta K_c = \pm 1$
b	\perp $\Delta J = 0$ strong $\Delta K_a = \pm 1$		\perp $\Delta J = 0$ strong $\Delta K_c = \pm 1$
c	\perp $\Delta J = 0$ strong $\Delta K_a = \pm 1$		$\Delta J = 0$ weak $\Delta K_c = \pm 0$

NOTATION examples



\Downarrow expressed as ...nopqrst...
 $\Delta K_p \text{ or } \Delta J_{a \text{ or } b \text{ or } c}$ (J_{K_p}, J_{K_o})
 near prolate near oblate

${}^q R_a(3_{1,2})$ ${}^r R_a(3_{1,2})$



Can decide whether the other asymmetry doublet transition is at higher or lower frequency. Asymmetry splitting is largest for largest J-K: $\nu(^pR_a(3_{1,2})) > \nu(^pR_a(3_{1,3}))$.

I am not sure that qualitative predictions can be made when $\Delta K = \Delta J$.

A lot of information about expected patterns can be deduced from qualitative prolate ↔ oblate correlation diagram and simple transition intensity ideas.

NON-LECTURE

Table of transition intensities in Townes and Schawlow, pages 557-612 (similar to diatomic M-averaged rotational linestrength factors).

$$^x S_{J_{kl} J'_{mn}}(\kappa) \equiv 10^4 (2J + 1) \left| \left\langle J_{kl} \left| \overset{\text{M averaged}}{M_x \alpha_x \bar{S}} \right| J'_{mn} \right\rangle \right|^2$$

listed for
 $\kappa = \pm 1, \pm 0.5, 0$

$x = a, b, c$
 $K_p: k, m$
 $K_o: l, n$

Transitions are listed by sub-branch and labeled by lower state quantum numbers. Notation for rotational level and rotational transition.

$$\begin{array}{c}
 \mathbf{J} \\
 \begin{array}{cc}
 \uparrow & \uparrow \\
 K_p & K_o \\
 \uparrow & \uparrow \\
 K_{-1} & K_{+1}
 \end{array}
 \end{array}
 \quad
 \mathbf{R}_{\Delta K_p \Delta K_o} \left(\mathbf{J}_{K_p, K_o} \right)$$

2 left columns in T and S table are for upper sign of κ
 2 right columns in T and S table are for lower sign of κ

But we are modern and have computers. How to do it?

* specify A, B, C

* diagonalize \mathbf{H}^{ROT} $\begin{cases} E(J_{K_p K_o}) \\ \Psi_{J_{K_p K_o}} \end{cases}$

* evaluate $M_b \alpha_{bS}$ matrix elements following Hougen, pp. 32-41 (especially 37-41)

* sum over M_J components

What does actual microwave spectrum look like?

Complicated $J \rightarrow 2J + 1$ sublevels for each J

Many rotational branches

No obvious pattern for Q branches

Do H_2CO instead of NO_2 .

Which is a, b, c

oblate or prolate? (possibility of switch to oblate?)

where is μ ? what is type? selection rules?

Example: NO_2

$$\kappa \equiv \frac{2B - A - C}{A - C}$$

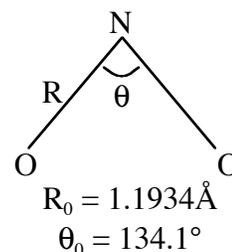
$$A_0 = 8.0012 \text{ cm}^{-1}$$

$$B_0 = 0.43364$$

$$C_0 = 0.41040$$

$$\kappa = -0.8857 \text{ (near prolate)}$$

μ is in-plane \perp to O—O direction



light axis: a, \parallel to O—O direction (through center of mass)

heavy axis: c, \perp to plane

$\therefore \mu$ is \parallel b axis

Special relationship for planar molecules

$$\frac{1}{A} + \frac{1}{B} = \frac{1}{C}$$

$$\text{inertial defect: } -\left(\frac{1}{A_e} + \frac{1}{B_e} - \frac{1}{C_e}\right) = +5.6 \times 10^{-3}$$

(not quite 0 because $v = 0$ not e constants used)

b-type transitions: $\Delta K_p = \text{odd}$, $\Delta K_o = \text{odd}$, $\Delta J = 0$ strong, $\Delta J = \pm 1$

types of
branches

${}^p\text{P}$
downward

${}^p\text{Q}$
downward

${}^p\text{R}$
downward
at low J

${}^r\text{P}$
upward at
low J

${}^r\text{Q}$
upward

${}^r\text{R}$
upward

~ 3 upward transitions for each $J_{K_a K_c}$

$\sim 3(2J + 1)$ upward transitions for each J

How to assign?

* Look for patterns

series of lines spaced by $\sim 2\bar{B}$ (R-branches)

also repeated K-stack spacings

pairs of lines corresponding to asymmetry doublets (R-branches)

size of asymmetry splitting - large for large J-K

absence of asymmetry splitting : $K = 0 \leftrightarrow 0, 1 \leftrightarrow 0, 0 \leftrightarrow 1$

* T-dependence

high E^{ROT} lines increase in intensity rapidly as $T \uparrow$

low E^{ROT} lines decrease in intensity slightly as $T \uparrow$

* guess structure and compute spectrum — frequencies and relative intensities

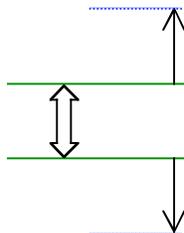
* tricks:

Stark (ϵ -field) modulation

picks out low-J lines where asymmetry splitting is small $J \approx K$

Double Resonance — bootstrap — transfer one assignment to neighboring linked transitions

saturate modulate label



Recall, for planar molecule there are only 2 independent bits of information. 2 correct assignments lead to accurate prediction of entire spectrum.

* isotope shifts. Change mass of atom that lies on one a, b, or c axis - corresponding rotational constant unchanged, others affected!