

6.730 Physics for Solid State Applications

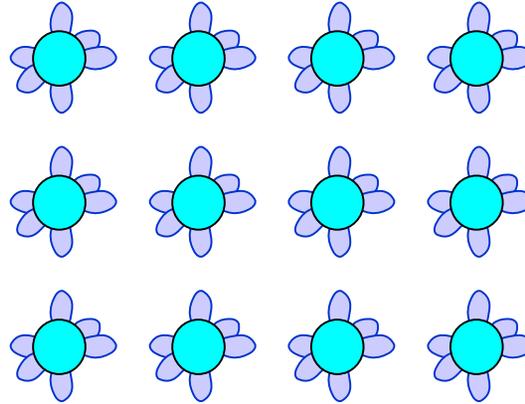
Lecture 15: Electrons in a Periodic Solid

Outline

- Review 2-D Tight-binding
- 3-D Tight-binding
- Semiconductor Fermi Energy
- Silicon Bandstructure

2D Monatomic Square Crystals

$$\mathbf{c}[\mathbf{R}_n] = \begin{pmatrix} c_s \\ c_{p_x} \\ c_{p_y} \\ c_{p_z} \end{pmatrix}$$



$$\mathbf{H}(\mathbf{k}) = \begin{pmatrix} \langle \phi_s | & \langle \phi_{p_x} | & \langle \phi_{p_y} | & \langle \phi_{p_z} | \\ \left(\begin{array}{cccc} E_s + V_{ss\sigma}g_0 & V_{sp\sigma}g_1 & V_{sp\sigma}g_2^* & 0 \\ V_{sp\sigma}g_1^* & E_p + V_{pp\pi}g_3 + V_{pp\sigma}g_4 & 0 & 0 \\ V_{sp\sigma}g_2 & 0 & E_p + V_{pp\pi}g_4 + V_{pp\sigma}g_3 & 0 \\ 0 & 0 & 0 & E_p + V_{pp\pi}g_0 \end{array} \right) \end{pmatrix}$$

$$g_0 = e^{-ik_x a} + e^{ik_x a} + e^{-ik_y a} + e^{ik_y a}$$

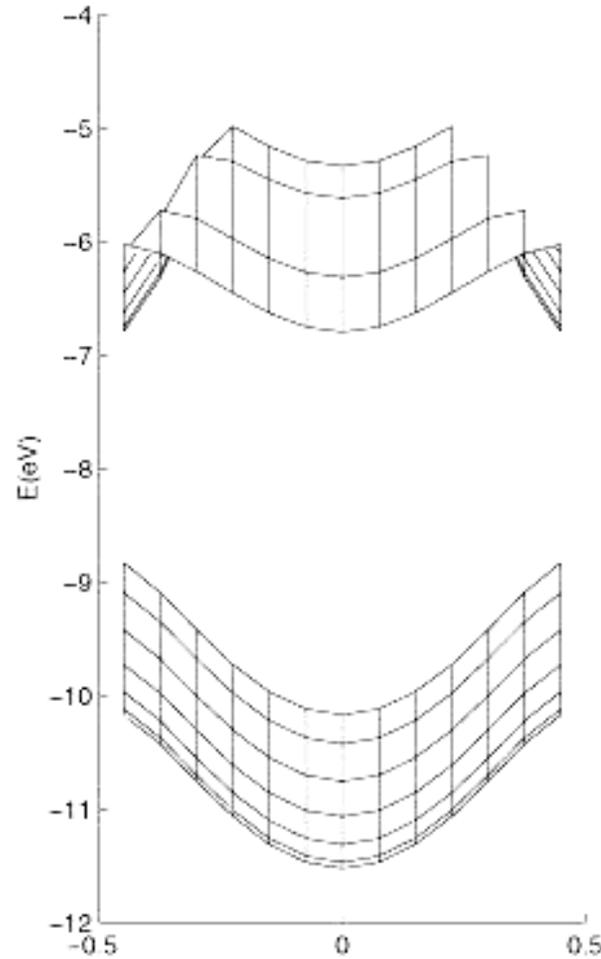
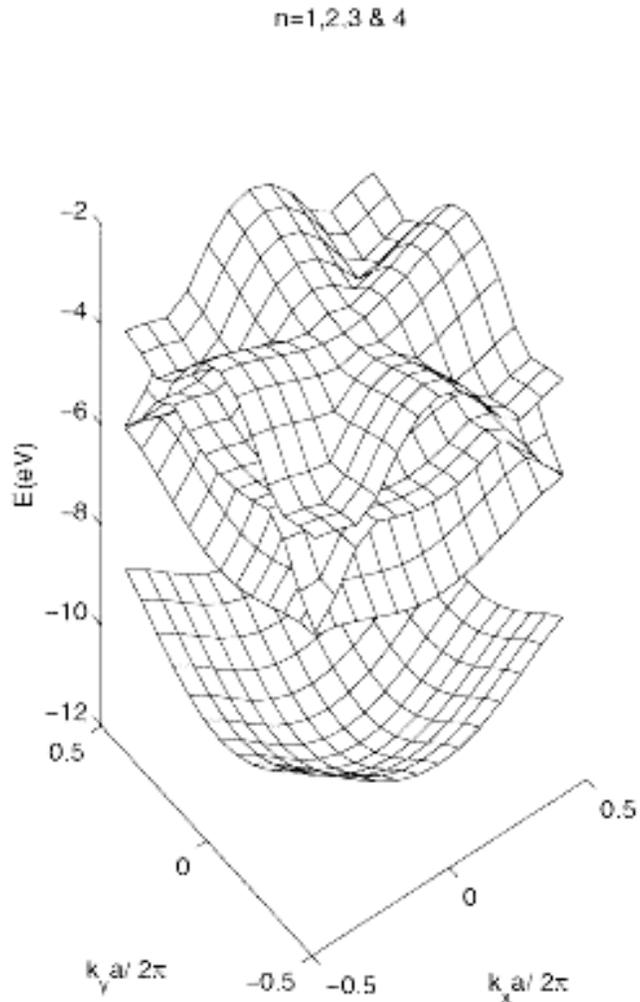
$$g_1 = e^{-ik_x a} - e^{ik_x a}$$

$$g_3 = e^{-ik_y a} + e^{ik_y a}$$

$$g_2 = -e^{-ik_y a} + e^{+ik_y a}$$

$$g_4 = e^{-ik_x a} + e^{ik_x a}$$

2D Monatomic Square Crystals Dispersion Relations



$$E_s = -10.11 \text{ eV}$$

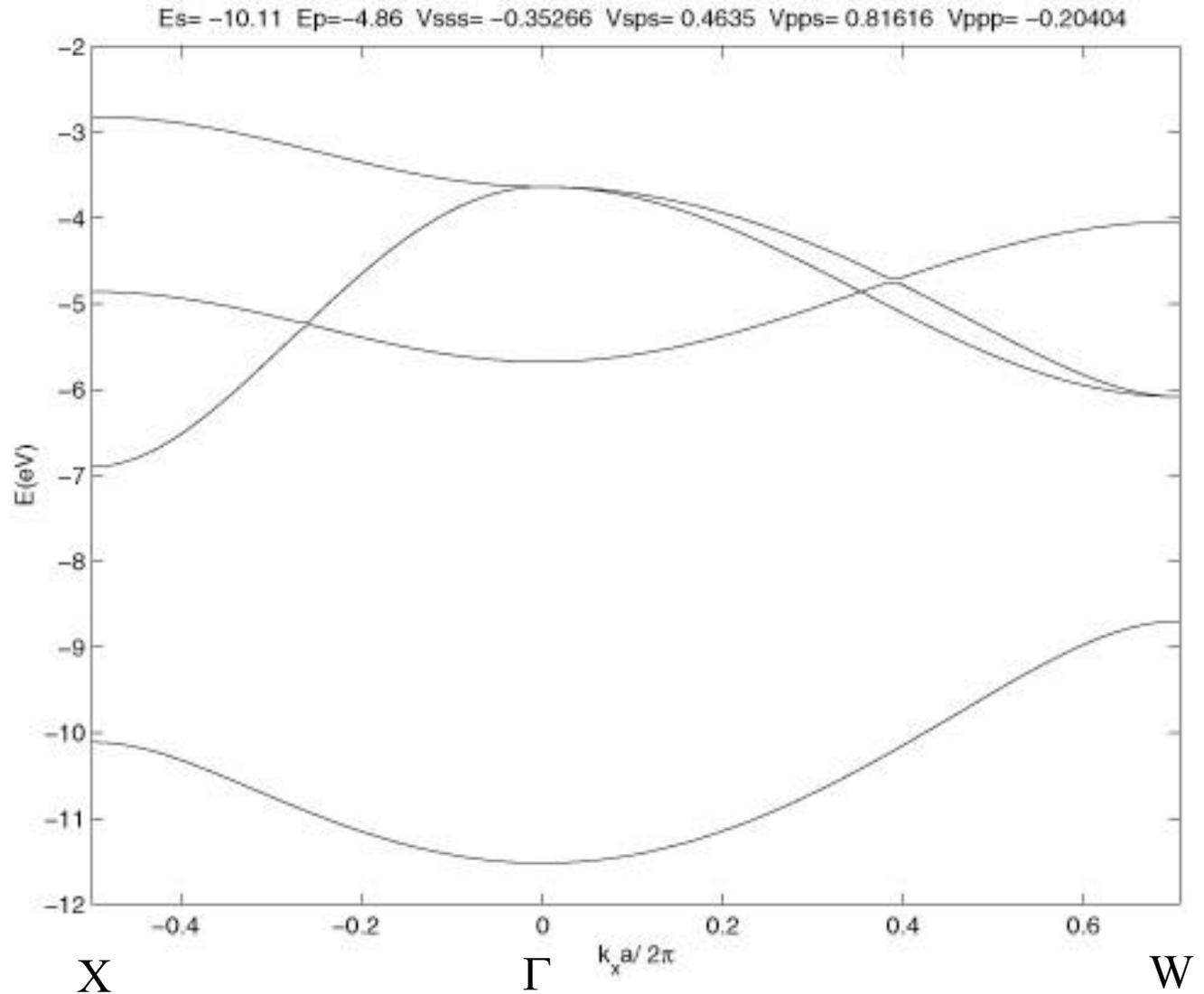
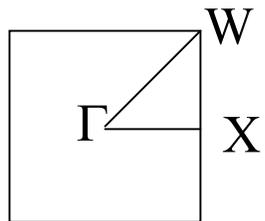
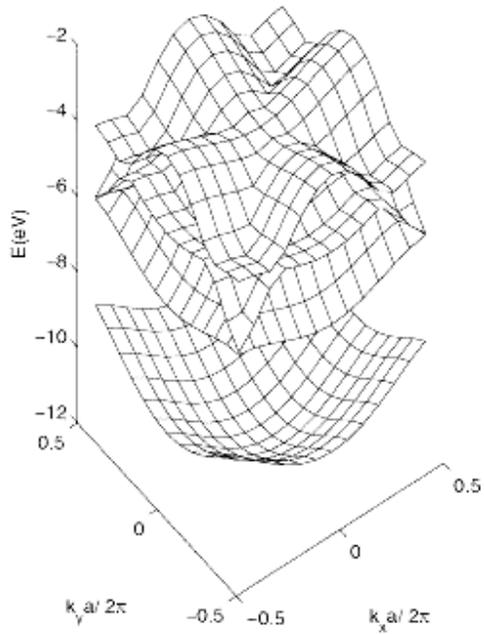
$$E_p = -4.86 \text{ eV}$$

$$a = 5.5 \text{ \AA}$$

2D Monatomic Square Crystals

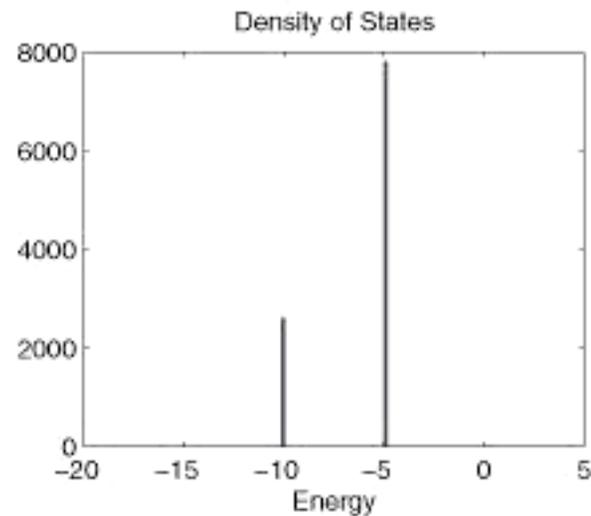
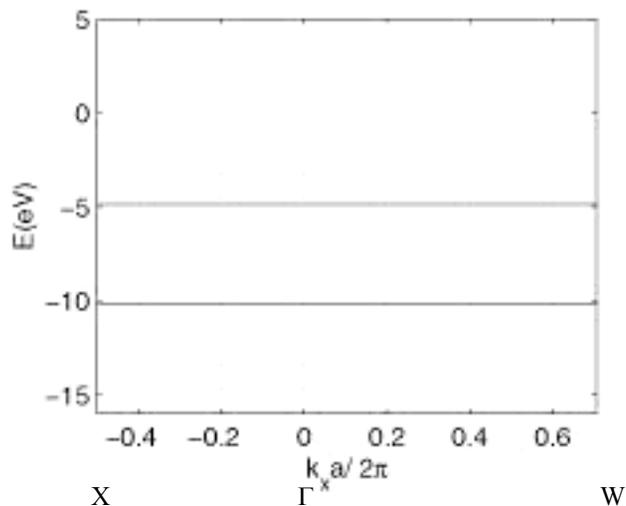
Dispersion Relations

$n=1,2,3 \& 4$

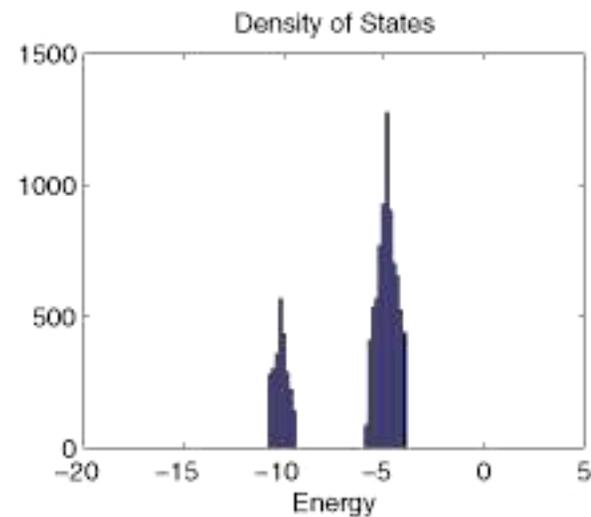
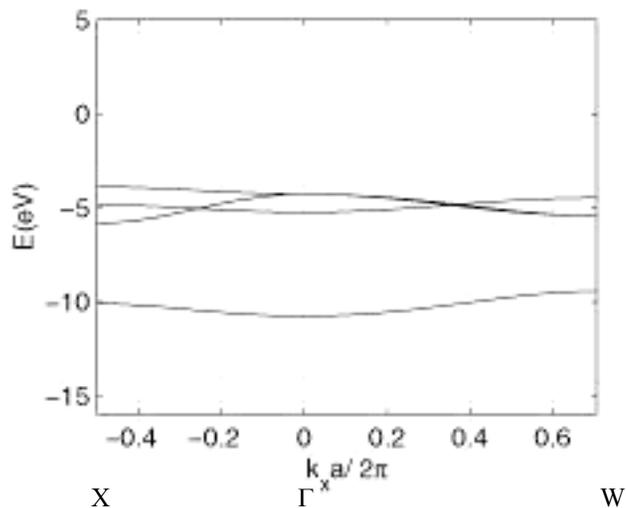


2D Monatomic Square Crystals

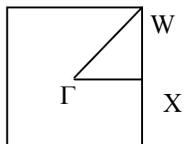
Variations with Lattice Constant



$a \rightarrow \infty$

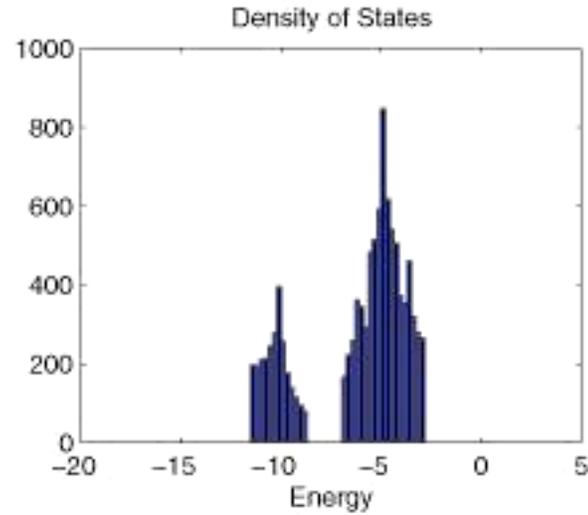
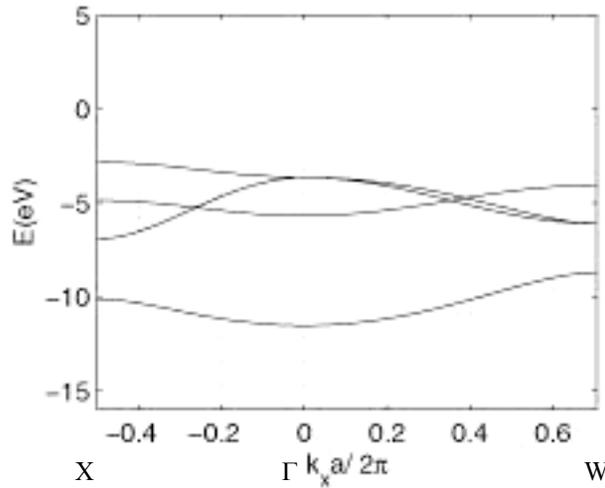


$a = 8.3 \text{ \AA}$

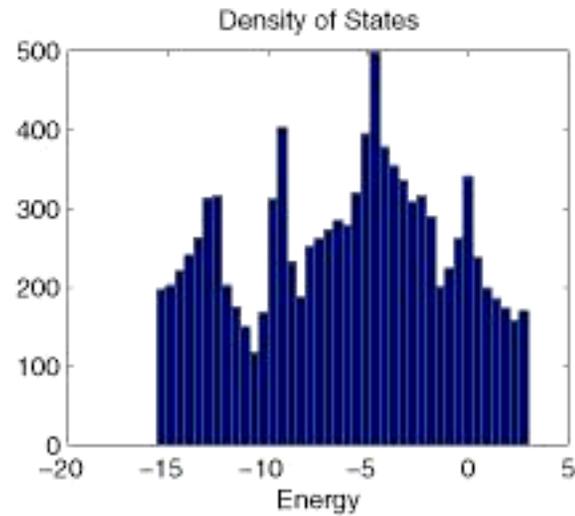
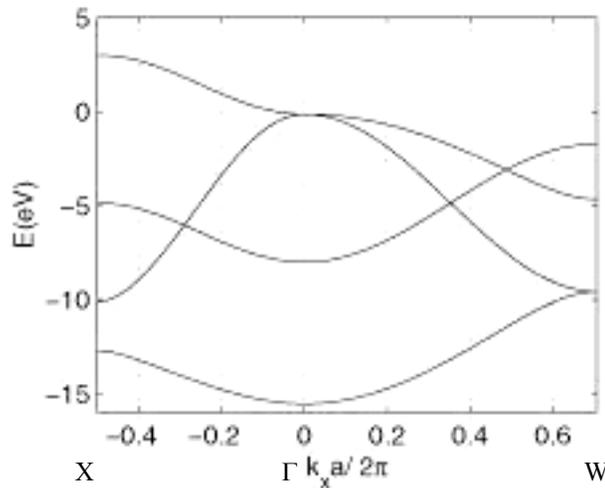


2D Monatomic Square Crystals

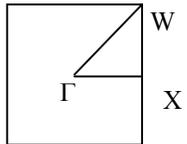
Dispersion Relations



$a = 5.5 \text{ \AA}$



$a = 2.8 \text{ \AA}$



2D Monatomic Square Crystals

Fermi Energy

How many states per band ?

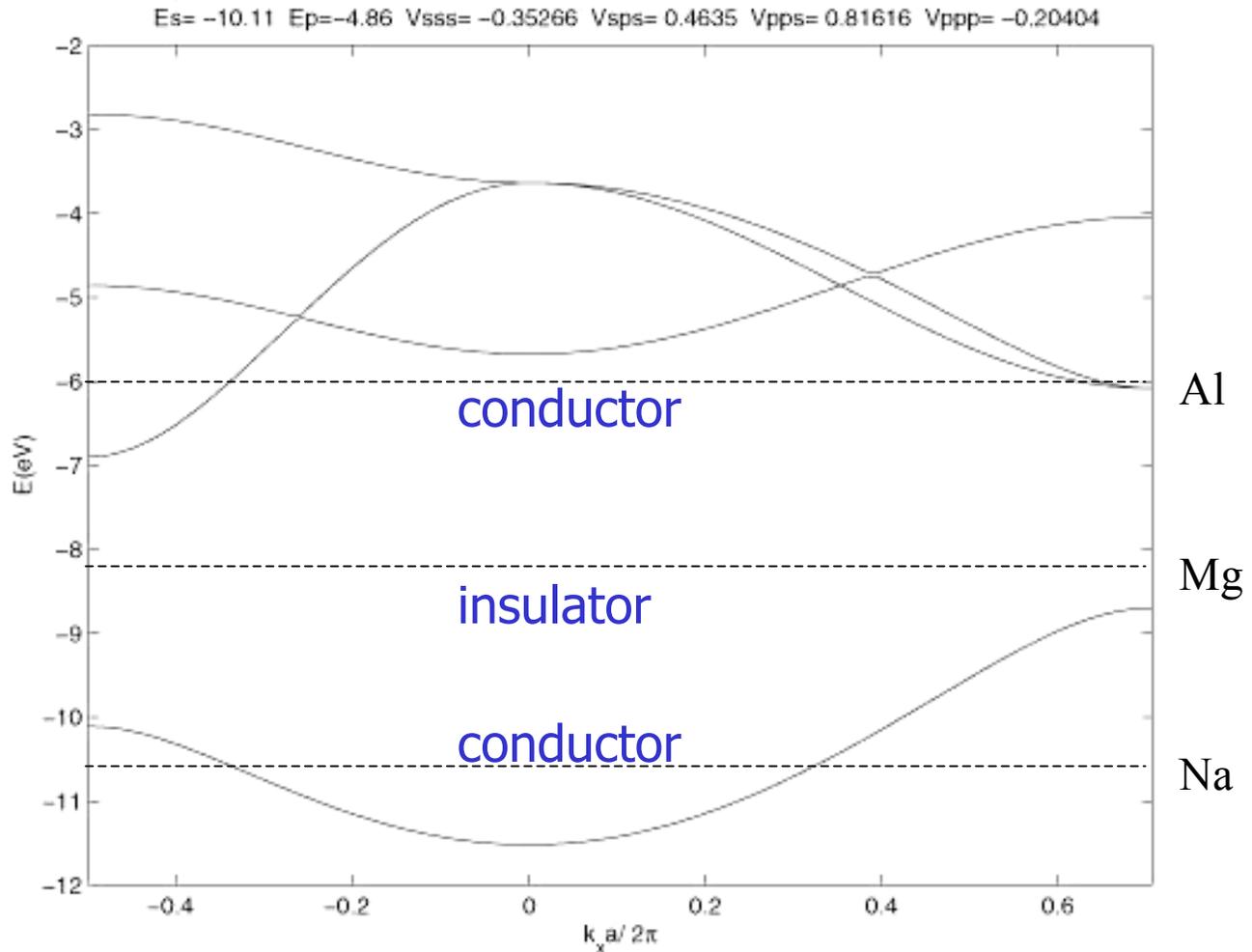
$$2 \cdot \frac{\text{Area of BZ}}{\text{Area per state}} = \frac{2 \cdot \left(\frac{2\pi}{a}\right)^2}{\left(\frac{2\pi}{L}\right)^2} = 2n$$

where n is the areal density of atoms

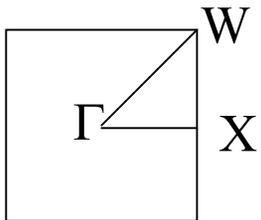
To estimate Fermi energy we need to know the number of outermost valence electrons each atom has...

I	II	III	IV	V	VI	VII	VIII
11	12	13	14	15	16	17	18
Na	Mg	Al	Si	P	S	Cl	Ar

2D Monatomic Square Crystals Dispersion Relations



$a = 5.5 \text{ \AA}$



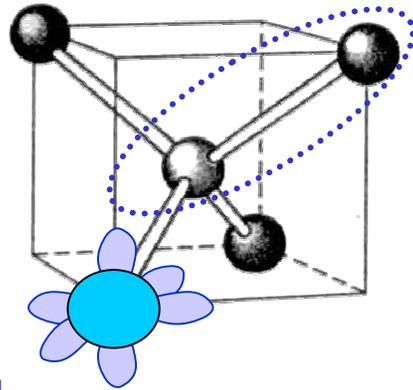
Reducing a , makes Mg a conductor (semimetal) !

I	II
11 Na	12 Mg

III	IV	V	VI	VII	VIII
13 Al	14 Si	15 P	16 S	17 Cl	18 Ar

LCAO Basis for FCC Crystals

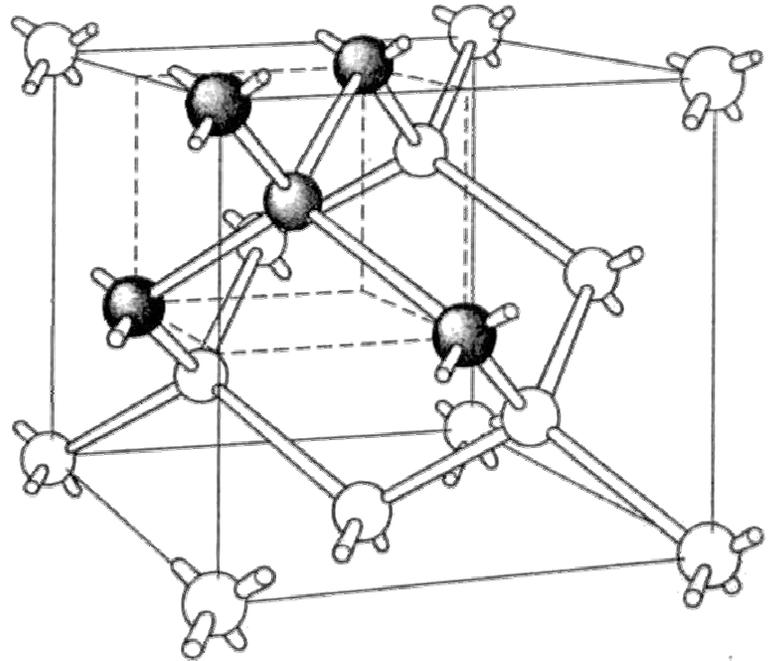
$$\psi(\mathbf{r}) = \sum_{\alpha} \sum_{\mathbf{n}=\mathbf{R}_{\ell}} \mathbf{c}_{\alpha}[\mathbf{R}_{\ell}] \phi_{\alpha}(\mathbf{r} - \mathbf{R}_{\ell})$$



Ga: [Ar]3d¹⁰ 4s² 4p¹

As: [Ar]3d¹⁰ 4s² 4p³

$$\mathbf{c}[\mathbf{R}_{\mathbf{n}}] = \begin{pmatrix} c_{s1} \\ c_{p_{x1}} \\ c_{p_{y1}} \\ c_{p_{z1}} \\ c_{s2} \\ c_{p_{x2}} \\ c_{p_{y2}} \\ c_{p_{z2}} \end{pmatrix}$$



Tight-binding for 3-D Crystals

$$\psi(\mathbf{r}) = \sum_{\alpha} \sum_{\mathbf{n}=\mathbf{R}_{\ell}} c_{\alpha}[\mathbf{R}_{\ell}] \phi_{\alpha}(\mathbf{r} - \mathbf{R}_{\ell})$$

Best estimate for energy with LCAO basis....

$$\sum_{\alpha} \sum_{\mathbf{R}_m} \widetilde{\mathbf{H}}_{\beta,\alpha}(\mathbf{R}_n, \mathbf{R}_m) c_{\alpha}[\mathbf{R}_m] = E \sum_{\alpha} \sum_{\mathbf{R}_p} \widetilde{\mathbf{S}}_{\beta,\alpha}(\mathbf{R}_n, \mathbf{R}_p) c_{\alpha}[\mathbf{R}_p]$$

Hamiltonian matrix....

$$\widetilde{\mathbf{H}}_{\beta,\alpha}(\mathbf{R}_n, \mathbf{R}_m) = \langle \phi_{\beta}(\mathbf{r} - \mathbf{R}_n) | \hat{\mathcal{H}} | \phi_{\alpha}(\mathbf{r} - \mathbf{R}_m) \rangle$$

Overlap matrix....

$$\widetilde{\mathbf{S}}_{\beta,\alpha}(\mathbf{R}_n, \mathbf{R}_p) = \langle \phi_{\beta}(\mathbf{r} - \mathbf{R}_n) | \phi_{\alpha}(\mathbf{r} - \mathbf{R}_p) \rangle$$

Tight-binding for 3-D Crystals

$$\psi(\mathbf{r}) = \sum_{\alpha} \sum_{\mathbf{n}=\mathbf{R}_{\ell}} \mathbf{c}_{\alpha}[\mathbf{R}_{\ell}] \phi_{\alpha}(\mathbf{r} - \mathbf{R}_{\ell})$$

Since the probability of finding electrons at each lattice site is equal...

$$\mathbf{c}[\mathbf{R}_{\mathbf{n}} + \mathbf{R}_{\ell}] = e^{i\mathbf{k}\cdot\mathbf{R}_{\ell}} \mathbf{c}[\mathbf{R}_{\mathbf{n}}]$$

$$\mathbf{c}[\mathbf{R}_{\mathbf{n}}] = e^{i\mathbf{k}\cdot\mathbf{R}_{\mathbf{n}}} \tilde{\epsilon}$$

Consequently...

$$\mathbf{H}(\mathbf{k}) \tilde{\epsilon} = \mathbf{E} \mathbf{S}(\mathbf{k}) \tilde{\epsilon}$$

$$\mathbf{H}(\mathbf{k}) = \sum_{\mathbf{R}_{\mathbf{p}}} \tilde{\mathbf{H}}(\mathbf{R}_{\mathbf{p}}) e^{-i\mathbf{k}\cdot\mathbf{R}_{\mathbf{p}}}$$

$$\mathbf{S}(\mathbf{k}) = \sum_{\mathbf{R}_{\mathbf{p}}} \tilde{\mathbf{S}}(\mathbf{R}_{\mathbf{p}}) e^{-i\mathbf{k}\cdot\mathbf{R}_{\mathbf{p}}}$$

Energy Band for 1-D Lattice

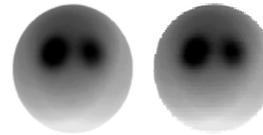
Two orbital, single atom basis

Hamiltonian Matrix

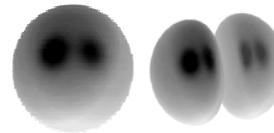
$$E_s = \langle \phi_s(\mathbf{r}) | \hat{\mathcal{H}} | \phi_s(\mathbf{r}) \rangle$$

$$E_p = \langle \phi_p(\mathbf{r}) | \hat{\mathcal{H}} | \phi_p(\mathbf{r}) \rangle$$

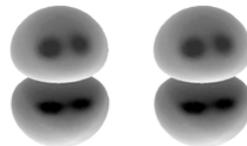
$$V_{ss\sigma} = \langle \phi_s(\mathbf{r}) | \hat{\mathcal{H}} | \phi_s(\mathbf{r} - \mathbf{a}_x) \rangle$$



$$V_{sp\sigma} = \langle \phi_s(\mathbf{r}) | \hat{\mathcal{H}} | \phi_{p_x}(\mathbf{r} - \mathbf{a}_x) \rangle$$



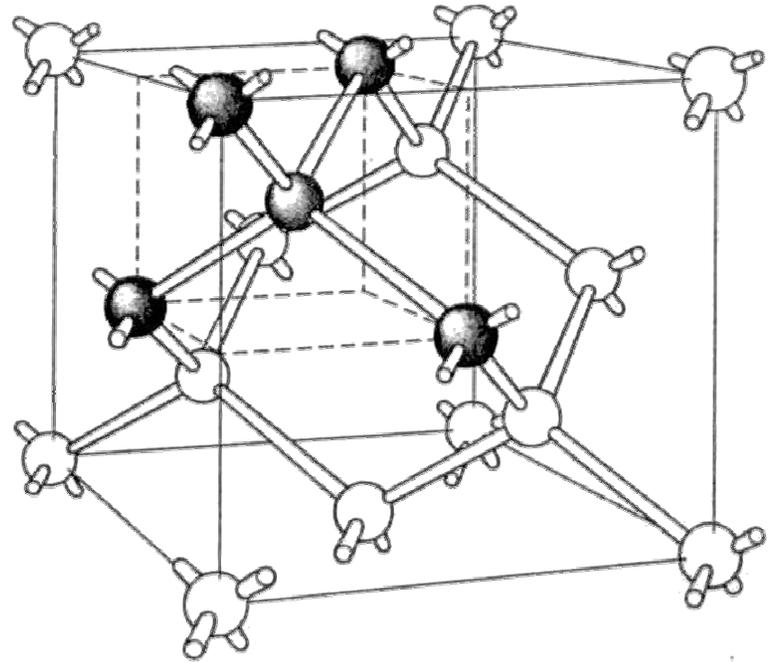
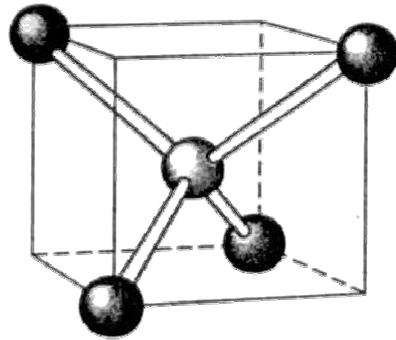
$$V_{pp\sigma} = \langle \phi_{p_x}(\mathbf{r}) | \hat{\mathcal{H}} | \phi_{p_x}(\mathbf{r} - \mathbf{a}_x) \rangle$$



$$V_{pp\pi} = \langle \phi_{p_y}(\mathbf{r}) | \hat{\mathcal{H}} | \phi_{p_y}(\mathbf{r} - \mathbf{a}_x) \rangle$$

Orbital Overlaps for 3-D Crystals

Diamond and Zincblende



$$V_{\ell\ell'm} = \eta_{\ell\ell'm} \frac{\hbar^2}{2md^2}$$

$$\eta_{ss\sigma} = -1.40 \approx -9\pi^2/64$$

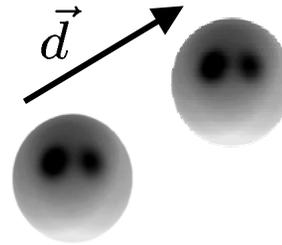
$$\eta_{pp\sigma} = 3.24$$

$$\eta_{sp\sigma} = 1.84$$

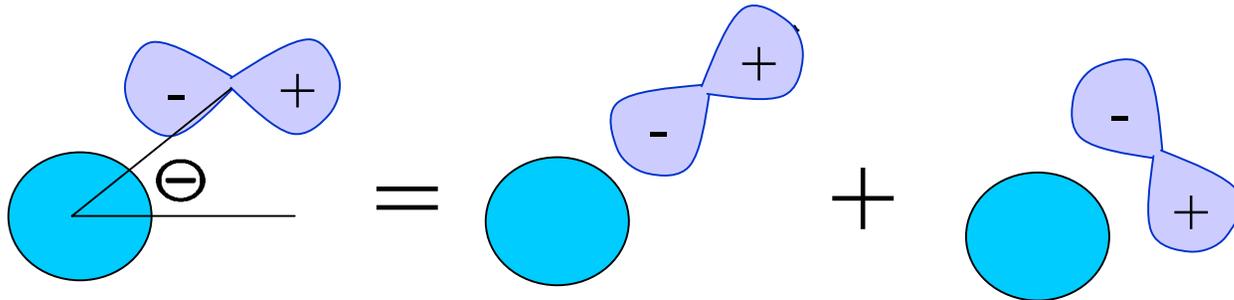
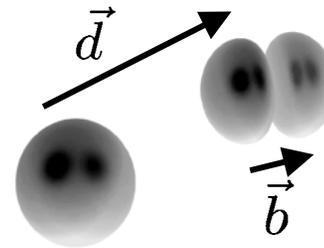
$$\eta_{pp\pi} = -0.81$$

Orbital Overlaps for 3-D Crystals

$$\langle \phi_{s1} | \hat{\mathcal{H}} | \phi_{s2} \rangle = V_{ss\sigma}$$

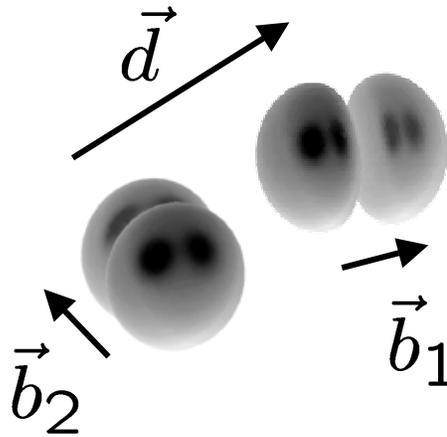


$$\langle \phi_{s1} | \hat{\mathcal{H}} | \phi_{p2} \rangle = V_{sp\sigma} \hat{\mathbf{d}} \cdot \hat{\mathbf{b}}$$

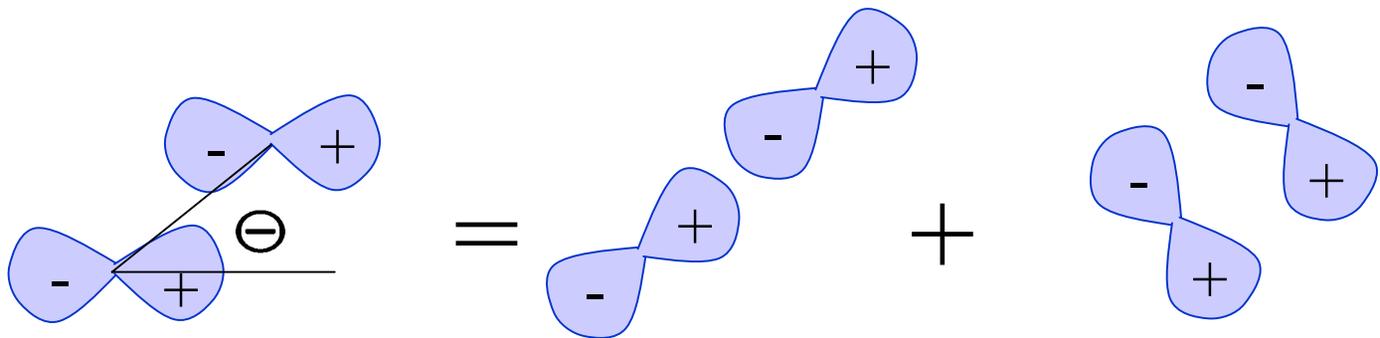


$$\langle \phi_{s1} | \hat{\mathcal{H}} | \phi_{p_y2} \rangle = V_{sp\sigma} \cos\Theta + 0 \sin\Theta$$

Orbital Overlaps for 3-D Crystals



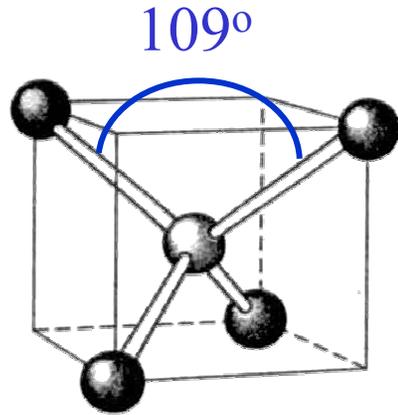
$$\langle \phi_{p1} | \hat{\mathcal{H}} | \phi_{p2} \rangle = (\hat{\mathbf{d}} \cdot \hat{\mathbf{b}}_1)(\hat{\mathbf{d}} \cdot \hat{\mathbf{b}}_2) V_{pp\sigma} + [\hat{\mathbf{b}}_1 - \hat{\mathbf{d}}(\hat{\mathbf{b}}_1 \cdot \hat{\mathbf{d}})] \cdot [\hat{\mathbf{b}}_2 - \hat{\mathbf{d}}(\hat{\mathbf{b}}_2 \cdot \hat{\mathbf{d}})] V_{pp\pi}$$



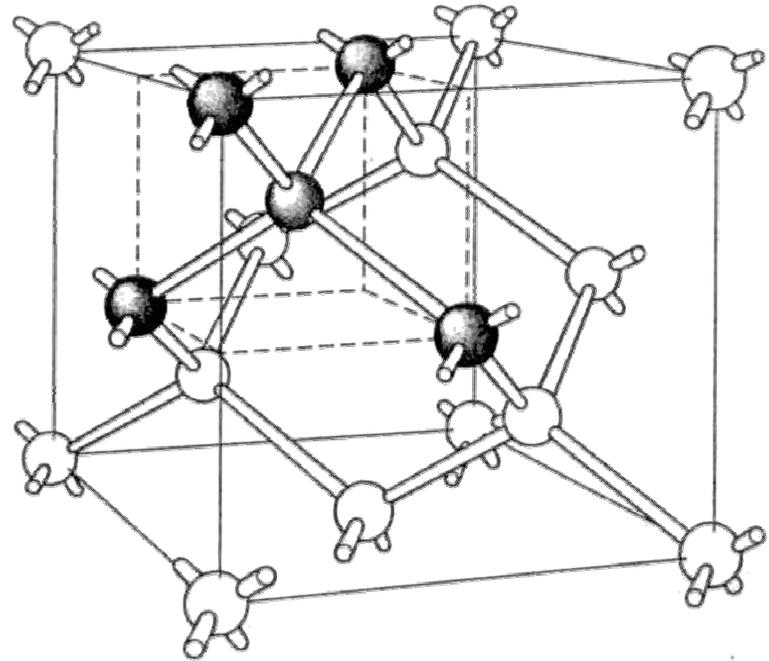
$$\langle \phi_{p_{y1}} | \hat{\mathcal{H}} | \phi_{p_{y2}} \rangle = V_{pp\sigma} \cos^2 \Theta + V_{pp\pi} \sin^2 \Theta$$

Orbital Overlaps for 3-D Crystals

Diamond and Zincblende



$$\Theta = \frac{109^\circ}{2}$$



$$\langle \phi_{s1} | \hat{\mathcal{H}} | \phi_{p_{y2}} \rangle = V_{sp\sigma} \cos \Theta = V_{sp\sigma} \frac{1}{\sqrt{3}}$$

$$\langle \phi_{p_{y1}} | \hat{\mathcal{H}} | \phi_{p_{y2}} \rangle = V_{pp\sigma} \cos^2 \Theta + V_{pp\pi} \sin^2 \Theta = V_{pp\sigma} \frac{1}{3} + V_{pp\pi} \frac{2}{3}$$

Zincblende LCAO Bands

Reduced Hamiltonian Matrix

$$\begin{array}{l}
 \langle \phi_{s1} | \\
 \langle \phi_{p_x1} | \\
 \langle \phi_{p_y1} | \\
 \langle \phi_{p_z1} | \\
 \langle \phi_{s2} | \\
 \langle \phi_{p_x2} | \\
 \langle \phi_{p_y2} | \\
 \langle \phi_{p_z2} |
 \end{array}
 \begin{pmatrix}
 |\phi_{s1}\rangle & |\phi_{p_x1}\rangle & |\phi_{p_y1}\rangle & |\phi_{p_z1}\rangle & |\phi_{s2}\rangle & |\phi_{p_x2}\rangle & |\phi_{p_y2}\rangle & |\phi_{p_z2}\rangle \\
 E_s & 0 & 0 & 0 & E_{ss}g_0 & E_{sp}g_1 & E_{sp}g_2 & E_{sp}g_3 \\
 0 & E_p & 0 & 0 & -E_{sp}g_1 & E_{xx}g_0 & E_{xy}g_3 & E_{xy}g_2 \\
 0 & 0 & E_p & 0 & -E_{sp}g_2 & E_{xy}g_3 & E_{xx}g_0 & E_{xy}g_1 \\
 0 & 0 & 0 & E_p & -E_{sp}g_3 & E_{xy}g_2 & E_{xy}g_1 & E_{xx}g_0 \\
 E_{ss}g_0^* & -E_{sp}g_1^* & -E_{sp}g_2^* & -E_{sp}g_3^* & E_s & 0 & 0 & 0 \\
 E_{sp}g_1^* & E_{xx}g_0^* & E_{xy}g_3^* & E_{xy}g_2^* & 0 & E_p & 0 & 0 \\
 E_{sp}g_2^* & E_{xy}g_3^* & E_{xx}g_0^* & E_{xy}g_1^* & 0 & 0 & E_p & 0 \\
 E_{sp}g_3^* & E_{xy}g_2^* & E_{xy}g_1^* & E_{xx}g_0^* & 0 & 0 & 0 & E_p
 \end{pmatrix}$$

$$E_{ss} = V_{ss\sigma}$$

$$E_{sp} = -\frac{1}{\sqrt{3}}V_{sp\sigma}$$

$$E_{xy} = \frac{1}{3}V_{pp\sigma} - \frac{1}{3}V_{pp\pi}$$

$$E_{xx} = \frac{1}{3}V_{pp\sigma} + \frac{2}{3}V_{pp\pi}$$

Zincblende LCAO Bands

Nearest Neighbors

$$\mathbf{R}_2 = -\frac{a}{2}(\hat{y} + \hat{z})$$

$$\mathbf{R}_3 = -\frac{a}{2}(\hat{z} + \hat{x})$$

$$\mathbf{R}_4 = -\frac{a}{2}(\hat{x} + \hat{y})$$

Zincblende LCAO Bands

Reduced Hamiltonian Matrix

$$\begin{array}{c}
 \langle \phi_{s1} | \\
 \langle \phi_{p_x1} | \\
 \langle \phi_{p_y1} | \\
 \langle \phi_{p_z1} | \\
 \langle \phi_{s2} | \\
 \langle \phi_{p_x2} | \\
 \langle \phi_{p_y2} | \\
 \langle \phi_{p_z2} |
 \end{array}
 \begin{pmatrix}
 |\phi_{s1}\rangle & |\phi_{p_x1}\rangle & |\phi_{p_y1}\rangle & |\phi_{p_z1}\rangle & |\phi_{s2}\rangle & |\phi_{p_x2}\rangle & |\phi_{p_y2}\rangle & |\phi_{p_z2}\rangle \\
 E_s & 0 & 0 & 0 & E_{ss}g_0 & E_{sp}g_1 & E_{sp}g_2 & E_{sp}g_3 \\
 0 & E_p & 0 & 0 & -E_{sp}g_1 & E_{xx}g_0 & E_{xy}g_3 & E_{xy}g_2 \\
 0 & 0 & E_p & 0 & -E_{sp}g_2 & E_{xy}g_3 & E_{xx}g_0 & E_{xy}g_1 \\
 0 & 0 & 0 & E_p & -E_{sp}g_3 & E_{xy}g_2 & E_{xy}g_1 & E_{xx}g_0 \\
 E_{ss}g_0^* & -E_{sp}g_1^* & -E_{sp}g_2^* & -E_{sp}g_3^* & E_s & 0 & 0 & 0 \\
 E_{sp}g_1^* & E_{xx}g_0^* & E_{xy}g_3^* & E_{xy}g_2^* & 0 & E_p & 0 & 0 \\
 E_{sp}g_2^* & E_{xy}g_3^* & E_{xx}g_0^* & E_{xy}g_1^* & 0 & 0 & E_p & 0 \\
 E_{sp}g_3^* & E_{xy}g_2^* & E_{xy}g_1^* & E_{xx}g_0^* & 0 & 0 & 0 & E_p
 \end{pmatrix}$$

$$g_0 = 1 + e^{-i\mathbf{k}\cdot\mathbf{R}_2} + e^{-i\mathbf{k}\cdot\mathbf{R}_3} + e^{-i\mathbf{k}\cdot\mathbf{R}_4}$$

$$g_1 = 1 + e^{-i\mathbf{k}\cdot\mathbf{R}_2} - e^{-i\mathbf{k}\cdot\mathbf{R}_3} - e^{-i\mathbf{k}\cdot\mathbf{R}_4}$$

$$g_3 = 1 - e^{-i\mathbf{k}\cdot\mathbf{R}_2} + e^{-i\mathbf{k}\cdot\mathbf{R}_3} - e^{-i\mathbf{k}\cdot\mathbf{R}_4}$$

$$g_4 = 1 - e^{-i\mathbf{k}\cdot\mathbf{R}_2} - e^{-i\mathbf{k}\cdot\mathbf{R}_3} + e^{-i\mathbf{k}\cdot\mathbf{R}_4}$$

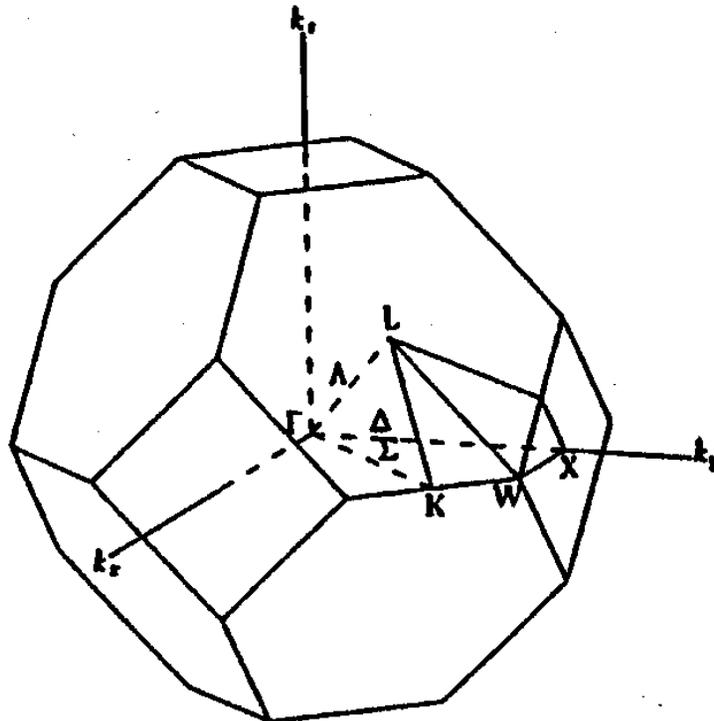
Silicon Bandstructure

Si: [Ne] $3s^2 3p^2$

4 e- per silicon atom

2 silicon atoms per lattice site

total: 8 electrons at each site



Silicon and Germanium Bandstructure



LCAO and Nearly Free Electron Bandstructure

$$\psi_i(r) = \sum_{\alpha} \sum_{\mathbf{R}_n} c_{i,\alpha[\mathbf{R}_n]} \phi_{\alpha}(r - \mathbf{R}_n) \quad \psi(r) = \sum_{\mathbf{R}} c_{\mathbf{k}} e^{i\mathbf{k}r}$$