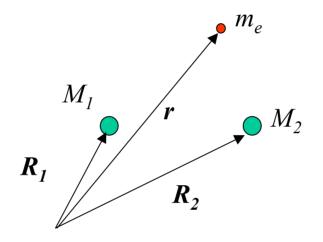
6.730 Physics for Solid State Applications

Lecture 2: Vibrational and Rotational States in Hydrogen

Rajeev J. Ram

Review Lecture 1: H₂

 $\psi(\mathbf{r},\mathbf{R_1},\mathbf{R_2})$, the wavefunction for the entire system of nuclei and electrons



$$\widehat{\mathcal{H}}\psi(\mathbf{r},\mathbf{R}_1,\mathbf{R}_2) = \mathbf{E}\psi(\mathbf{r},\mathbf{R}_1,\mathbf{R}_2)$$

$$\widehat{\mathcal{H}} = \frac{|\mathbf{P_1}|^2}{2M_1} + \frac{|\mathbf{P_2}|^2}{2M_2} + \frac{|\mathbf{p}|^2}{2m_e} - \frac{e^2}{4\pi\epsilon_o|\mathbf{r} - \mathbf{R_1}|} - \frac{e^2}{4\pi\epsilon_o|\mathbf{r} - \mathbf{R_2}|} + \frac{e^2}{4\pi\epsilon_o|\mathbf{R_1} - \mathbf{R_2}|}.$$

Approximate Models: Simplifying H₂

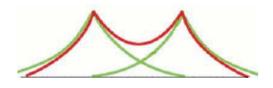
Born-Oppenheimer Approximation

- ◆ The electrons are much lighter than the nuclei (m_e/m_H≅1/1836), their motion is much faster than the vibrational and rotational motions of the nuclei within the molecule.
- → Works since vibrational and rotational energy of molecule is typically much less than the binding energy

$$\Psi(r, R_1, R_2) \approx \Phi(R_1, R_2) \psi(r; R_1, R_2)$$
.

Linear Combination of Atomic Orbitals (LCAO)

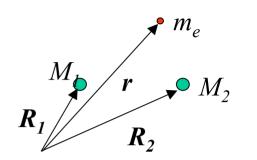
Even the electron part of the problem is too hard to solve exactly



$$\psi(r; R_1, R_2) = C \left[\phi_{1s}(r - R_1) + \phi_{1s}(r - R_2)\right]$$

Born-Oppenheimer Approximation

$$\Psi(\mathbf{r},\mathbf{R}_1,\mathbf{R}_2)\approx\Phi(\mathbf{R}_1,\mathbf{R}_2)\psi(\mathbf{r};\mathbf{R}_1,\mathbf{R}_2)$$



Electronic Part: $\psi(\mathbf{r}; \mathbf{R}_1, \mathbf{R}_2) \longrightarrow \psi(\mathbf{r})$

$$E_e(\mathbf{R_1}, \mathbf{R_2})\psi(\mathbf{r}) = \underbrace{\left[-\frac{\hbar^2 \nabla^2}{2m_e} - \frac{e^2}{4\pi\epsilon_o |\mathbf{r} - \mathbf{R_1}|} - \frac{e^2}{4\pi\epsilon_o |\mathbf{r} - \mathbf{R_2}|}\right]}_{\hat{\mathcal{H}}_e}\psi(\mathbf{r})$$

Nuclear Part:

$$E\Phi(\mathbf{R}_{1},\mathbf{R}_{2}) = \left[-\frac{\hbar^{2}\nabla_{1}^{2}}{2\mathbf{M}_{1}} - \frac{\hbar^{2}\nabla_{2}^{2}}{2\mathbf{M}_{2}} + \underbrace{E_{e}(\mathbf{R}_{1},\mathbf{R}_{2}) + \frac{\mathbf{e}^{2}}{4\pi\epsilon_{0}|\mathbf{R}_{1} - \mathbf{R}_{2}|}}_{\mathbf{V}_{eff}} \right] \Phi(\mathbf{R}_{1},\mathbf{R}_{2})$$

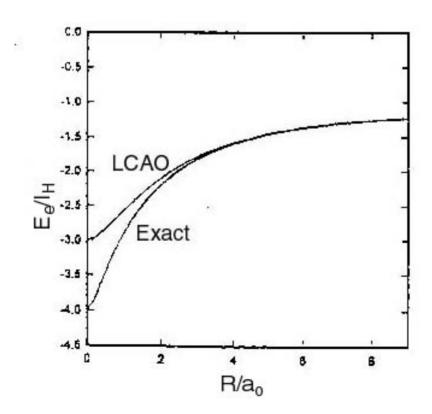
where E is the energy of the entire molecule

Approximate Electronic Energy

$$E_e(R) = \frac{E_s + V_{ss\sigma}}{1 + S(1)}$$

$$E_s(R) = E_s^o + \langle \phi_1(\mathbf{r}) | \hat{V}_2 | \phi_1(\mathbf{r}) \rangle$$

$$E_s^o = -I_H = -m_e e^4/(32\pi^2 \epsilon_o^2 \hbar^2) = -13.6 \text{ eV}$$



Nuclear and Electronic Energy Together

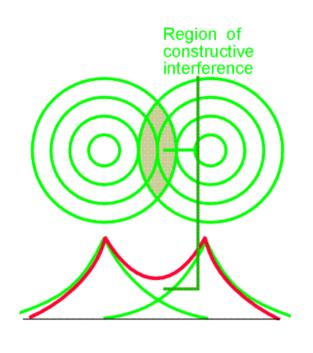
$$V_{\rm eff} = E_e(R) + \frac{e^2}{4\pi\epsilon_o R}$$

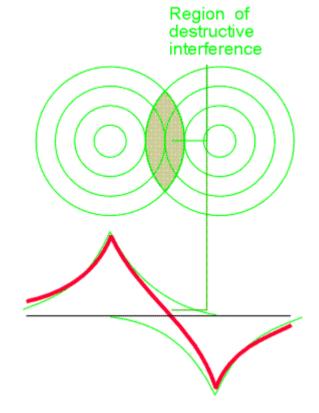
$$E_{\Phi(R_1,R_2)} = \begin{bmatrix} \frac{\hbar^2 \nabla_1^2}{2M_1} - \frac{\hbar^2 \nabla_2^2}{2M_2} + \underbrace{E_e(R_1,R_2) + \frac{e^2}{4\pi\epsilon_o(R_1-R_2)}} \\ \frac{e^2}{4\pi\epsilon_o(R_1-R_2)} \end{bmatrix} \Phi(R_1,R_2)$$

$$E(R_0) = V_{\rm eff}[LCAO](R_0) = -1.13 \, I_H$$

$$E(R_0) - E(R \to \infty) = -0.1 \, I_H = -1.4 \, {\rm eV} \, {\rm per \ bond}$$

First Excited State Energy: Antibonding





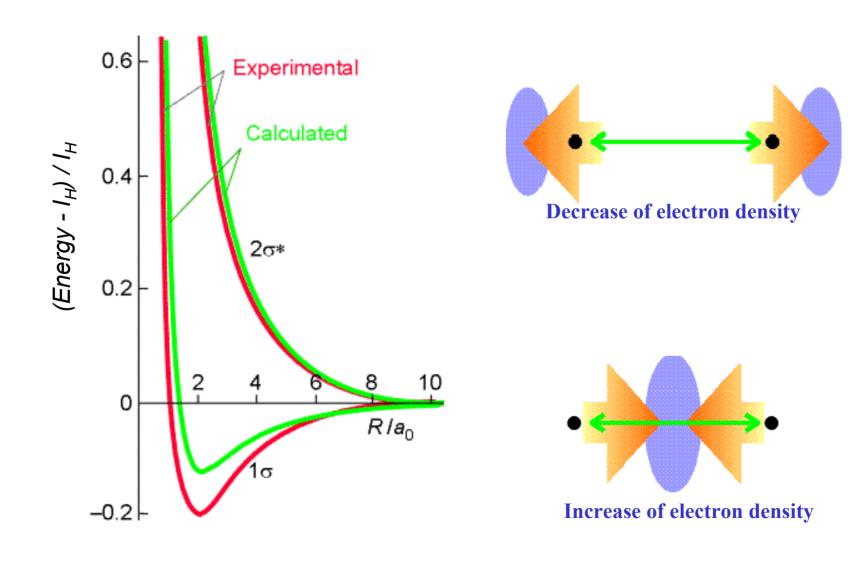
$$\psi_{1\sigma q}(\mathbf{r}) = \mathbf{C} \left[\phi(\mathbf{r} - \mathbf{R}_1) + \phi(\mathbf{r} - \mathbf{R}_2) \right]$$

$$E_{e,1\sigma g}(R) = \frac{E_s + V_{ss\sigma}}{1 + S(1)}$$

$$\psi_{1\sigma u}(\mathbf{r}) = C \left[\phi(\mathbf{r} - \mathbf{R}_1) - \phi(\mathbf{r} - \mathbf{R}_2) \right]$$

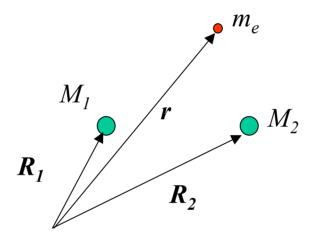
$$E_{e,1\sigma u}(R) = \frac{E_s - V_{ss\sigma}}{1 - S(1)}$$

First Excited State Energy: LCAO



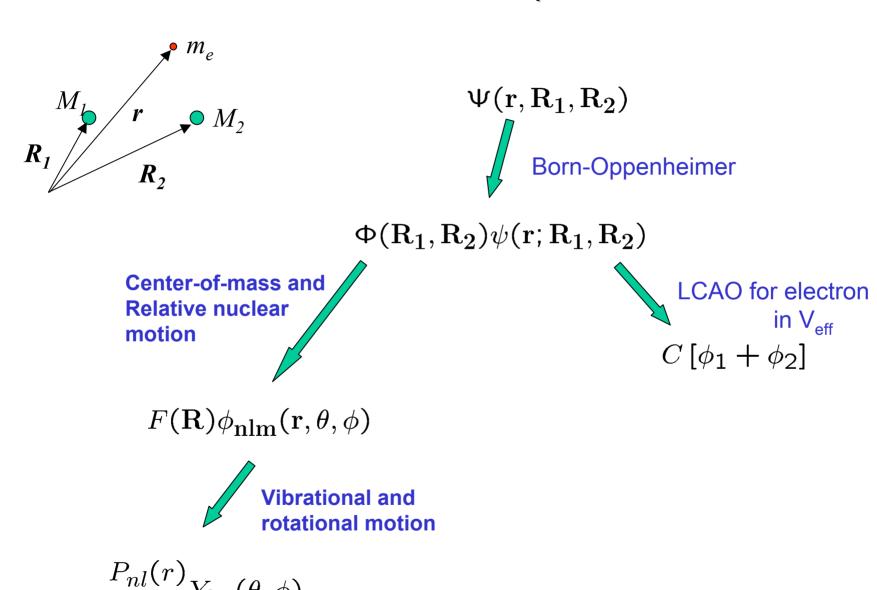
A Closer Look at Nuclear Motion

Molecular Vibration and Rotation



$$\begin{split} \mathit{E}\Phi(R_1,R_2) = \left[\frac{|P|_1^2}{2M_1} + \frac{|P|_2^2}{2M_2} + V_{\text{eff}}(|R_1-R_2|)\right]\Phi(R_1,R_2) \\ \text{LCAO for electron} \\ &\text{in V_{eff}} \end{split}$$

Divide and Conquer



Center-of-Mass and Relative Nuclear Motion

$$\mathit{E}\Phi(R_1,R_2) = \left[\frac{|P|_1^2}{2M_1} + \frac{|P|_2^2}{2M_2} + V_{\text{eff}}(|R_1 - R_2|)\right]\Phi(R_1,R_2)$$

$$E\Phi(\mathbf{r}, \mathbf{R}) = \left[\frac{|\mathbf{P}|^2}{2\mathbf{M}} + \frac{|\mathbf{p}|^2}{2\mu} + \mathbf{V}_{\mathsf{eff}}(|\mathbf{r}|)\right]\Phi(\mathbf{r}, \mathbf{R})$$

Center-of-mass

$$MR = M_1R_1 + M_2R_2$$

Note that R is the C-of-M coordinate now

$$P = p_1 + p_2$$

$$M = M_1 + M_2$$

Relative

$$r = R_2 - R_1$$

$$\frac{\mathbf{p}}{\mu} = \frac{\mathbf{P_2}}{M_2} - \frac{\mathbf{P_1}}{M_1}$$

$$\frac{1}{\mu} = \frac{1}{M_1} + \frac{1}{M_2}$$

Center-of-Mass and Relative Nuclear Motion

$$E\Phi(\mathbf{r}, \mathbf{R}) = \left[\frac{|\mathbf{P}|^2}{2\mathbf{M}} + \frac{|\mathbf{p}|^2}{2\mu} + \mathbf{V}_{\mathsf{eff}}(|\mathbf{r}|)\right]\Phi(\mathbf{r}, \mathbf{R})$$

$$\Phi(\mathbf{r}, \mathbf{R}) = \mathbf{F}(\mathbf{R})\phi(\mathbf{r})$$

$$E_{\mathbf{R}}F(\mathbf{R}) = \frac{|\mathbf{P}|^2}{2\mathbf{M}}F(\mathbf{R})$$
 $E_{\mathbf{r}}\phi(\mathbf{r}) = \left[\frac{|\mathbf{p}|^2}{2\mu} + V_{\mathsf{eff}}(|\mathbf{r}|)\right]\phi(\mathbf{r})$

Total energy is the sum of CM motion and relative:

$$E = E_{\mathbf{R}} + E_{\mathbf{r}}$$

If this is the total energy, where is the electron energy?

Center-of-Mass Nuclear Motion

Schrodinger equation for center-of-mass is same as free particle:

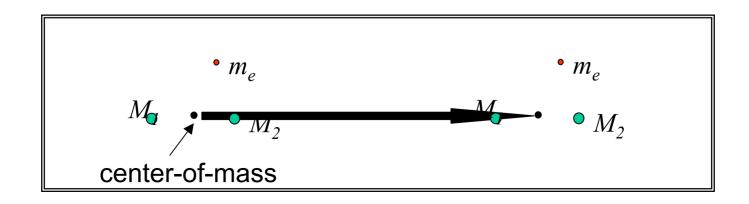
$$E_{\mathbf{R}}F(\mathbf{R}) = \frac{|\mathbf{P}|^2}{2\mathbf{M}}F(\mathbf{R})$$

Eigenstate:

$$F(\mathbf{R}) = e^{i\mathbf{K}\cdot\mathbf{R}}$$

Eigenenergy:

$$E_{\mathbf{R}} = \frac{\hbar^2 |\mathbf{K}|^2}{2M}$$



Relative Nuclear Motion

Schrodinger equation for relative motion is a central potential problem just like the hydrogen atom:

$$E_{\mathbf{r}}\phi(\mathbf{r}) = \left[\frac{|\mathbf{p}|^2}{2\mu} + V_{\text{eff}}(|\mathbf{r}|)\right]\phi(\mathbf{r})$$

$$E_{\mathbf{r}}\phi(r,\theta,\phi) = \begin{bmatrix} \frac{p_r^2}{2\mu} + \frac{|\mathbf{L}(\theta,\phi)|^2}{2\mu r^2} + V_{\text{eff}}(r) \end{bmatrix} \phi(r,\theta,\phi)$$
Radial kinetic energy

Angular kinetic energy

 $\mathbf{L}(heta,\phi)$ is the angular momentum operator

Relative Nuclear Motion Separation of Radial and Angular Components

$$E_{\mathbf{r}}\phi(r,\theta,\phi) = \left[\frac{p_r^2}{2\mu} + \frac{|\mathbf{L}(\theta,\phi)|^2}{2\mu r^2} + V_{\mathsf{eff}}(r)\right]\phi(r,\theta,\phi)$$

$$\phi_{nlm}(r,\theta,\phi) = \frac{P_{nl}(r)}{r} Y_{lm}(\theta,\phi)$$

 $Y_{lm}(\theta,\phi)$ is the spherical Bessel function

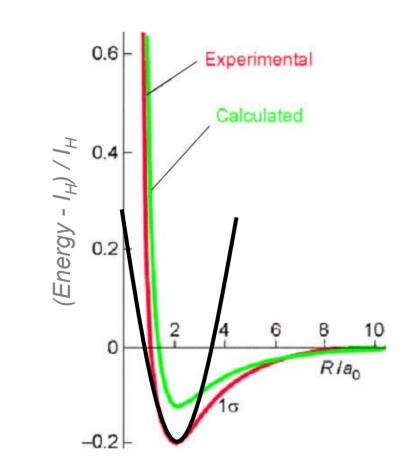
$$|L(\theta,\phi)|^2 Y_{lm}(\theta,\phi) = \hbar^2 l(l+1) Y_{lm}(\theta,\phi)$$

$$E_{\mathbf{r}}P_{nl}(r) = \left[-\frac{\hbar^2}{2\mu} \frac{d^2}{dr^2} + \frac{\hbar^2 l(l+1)}{2\mu r^2} + V_{\text{eff}}(r) \right] P_{nl}(r)$$

Vibrational Motion of Nuclei Harmonic Oscillator

For no rotation, this simplifies to...

$$E_{\mathbf{r}}^{n0}P_{n0}(r) = \left[-\frac{\hbar^2}{2\mu} \frac{d^2}{dr^2} + V_{\text{eff}}(r) \right] P_{n0}(r)$$



$$V_{\text{eff}}(r) = V_o + \frac{1}{2}(r - R_o)^2 \left(\frac{d^2V}{dr^2}\right)_{R_o}$$

Vibrational Motion of Nuclei Harmonic Oscillator

$$E_{\mathbf{r}}^{n0}P_{n0}(r) = \left[-\frac{\hbar^2}{2\mu} \frac{d^2}{dr^2} + V_o + \frac{1}{2}(r - R_o)^2 \left(\frac{d^2V}{dr^2} \right)_{R_o} \right] P_{n0}(r)$$

$$x = r - R_o \qquad \frac{1}{2}(r - R_o)^2 \left(\frac{d^2V}{dr^2} \right)_{R_o} = \frac{1}{2}\mu\omega_o^2 x^2$$

$$\left[E_{\mathbf{r}}^{n0} - V_o \right] P_{n0}(x) = \left[-\frac{\hbar^2}{2\mu} \frac{d^2}{dx^2} + \frac{1}{2} \mu \omega_o^2 x^2 \right] P_{n0}(x)$$

$$E_{\mathbf{r}}^{n0} = V_o + \hbar\omega_o(n + \frac{1}{2})$$

Approximation: Born-Oppenheimer, parabolic effective potential

Vibrational Motion of Nuclei Rigid Rotor

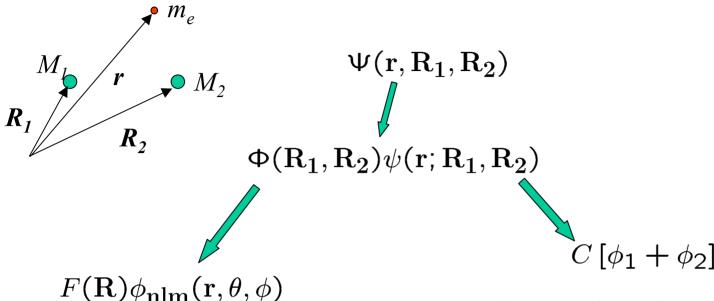
$$E_{\mathbf{r}}^{nl}P_{nl}(r) = \left[-\frac{\hbar^2}{2\mu} \frac{d^2}{dr^2} + \frac{\hbar^2 l(l+1)}{2\mu r^2} + V_o + \frac{1}{2}(r - R_o)^2 \left(\frac{d^2 V}{dr^2} \right)_{R_o} \right] P_{nl}(r)$$

Assuming that the vibrational motion produces only small displacements... $r \approx R_o$

$$\frac{\hbar^2 l(l+1)}{2\mu r^2} \approx \frac{\hbar^2 l(l+1)}{2\mu R_o^2}$$

$$E_{\mathbf{r}}^{nl} = V_o + \hbar \omega_o (n + \frac{1}{2}) + \frac{\hbar^2 l(l+1)}{2\mu R_o^2}$$

Divide and Conquer



Approximations

Born-Oppenheimer

Nuclei inside electron cloud act as if they are embedded in an elastic medium (V_{eff})

- Effective potential (V_{eff}) is parabolic Vibrations of simple harmonic oscillator
- Rigid rotor

Vibrations only displace nuclei slightly from equilibrium bond length

Total Energy of the H₂ Molecule

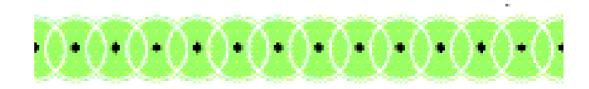
$$E = E_{n\ell mK} = E_{\mathbf{r}} + E_{\mathbf{R}}$$

$$E_{n\ell mK} = \underbrace{V_o}_{\text{Binding}} + \underbrace{\hbar \omega_o(n+\frac{1}{2})}_{\text{Vibrational}} + \underbrace{\frac{\hbar^2 l(l+1)}{2\mu R_o^2}}_{\text{Rotational}} + \underbrace{\frac{\hbar^2 |\mathbf{K}|^2}{2M}}_{\text{Translational}}$$

	Binding	Vibrational	Rotational	Translational
H ₂	1.4 eV	$0.5 (n + \frac{1}{2}) \text{ eV}$	7.5 / (/+1) meV	
O ₂	1.0 eV	0.1 $(n + \frac{1}{2})$ eV	30 / (/ +1) meV	
	30 THz	3-30 THz	0.3-3 THz	

Generalizations from Molecules to Solids

- The source of the binding energy is primarily the electrostatic potential between the nuclei and the electrons. The localization energy can also play a role (metal).
- Nuclear motions of the ions contribute a very small part to the binding energy.
- Sharing electrons between nuclei lowers the energy of the solid.
- The potential between the nuclei is of the same form as the molecule.
- Exicted states exists.



Assumptions for Electronic States

- One electron energy levels
- No spin or exchange energies
- LCAO a good approximation
- Ignore motion of the nuclei to first order









