

## Lecture 9: Macroscopic Quantum Model

### Outline

1. Development of Quantum Mechanics
2. Schrödinger's Equation
  - Free particle
  - With forces
  - With Electromagnetic force
3. Physical meaning of Wavefunction
  - Probability density and Probability Current density
4. Macroscopic Quantum Model

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## Macroscopic Quantum Model

Superconductivity is a  
Quantum Phenomenon on a  
macroscopic length scale.

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## Development of Quantum Mechanics

### Free Particle (no forces or potentials)

*Wave-like properties*

*Particle-like properties*

$$\mathcal{E} = \hbar\omega \quad \mathbf{p} = \hbar\mathbf{k} \quad \mathcal{E} = \frac{1}{2} m (\mathbf{v} \cdot \mathbf{v}) = \frac{\mathbf{p} \cdot \mathbf{p}}{2m}$$

frequency                  wavenumber                  energy                  momentum  $\rightarrow \mathbf{p} = m\mathbf{v}$

Planck's constant  $\hbar \equiv \frac{h}{2\pi} = 1.0546 \times 10^{-34} \text{ J} \cdot \text{sec}$

Combine wave and particle properties, to find the dispersion relation:

$$\hbar\omega = \frac{\hbar^2}{2m} (\mathbf{k} \cdot \mathbf{k})$$

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## Schrödinger's Equation (free particle)

$$\hbar\omega = \frac{\hbar^2}{2m} (\mathbf{k} \cdot \mathbf{k})$$

Assume that this results from a uniform plane wave solution

$$\psi = \hat{\psi} e^{i(\mathbf{k} \cdot \mathbf{r} - \omega t)}$$

Then a good guess of the differential equation that gives the dispersion relation is

$$i\hbar \frac{\partial}{\partial t} \psi = -\frac{\hbar^2}{2m} \nabla^2 \psi$$

This guess is justified by experimental confirmation; this is *not* a derivation.

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## Schrödinger's Equation (with forces)

We present a plausibility argument, not a derivation, relating the classical formulation to the quantum formulation.

The energy for a particle in a force is, classically,

$$\mathcal{E} = \frac{1}{2} m (\mathbf{v} \cdot \mathbf{v}) + V(\mathbf{r})$$

Energy is conserved since the potential is independent of time.

$$\begin{aligned} 0 &= \frac{d\mathcal{E}}{dt} = m\mathbf{v} \cdot \frac{d\mathbf{v}}{dt} + \frac{d}{dt} V(\mathbf{r}) \\ &= m\mathbf{v} \cdot \frac{d\mathbf{v}}{dt} + \frac{\partial}{\partial t} V(\mathbf{r}) + (\mathbf{v} \cdot \nabla) V(\mathbf{r}) \\ &= \mathbf{v} \cdot \left( m \frac{d\mathbf{v}}{dt} + \nabla V \right) \implies m \frac{d\mathbf{v}}{dt} = -\nabla V \end{aligned}$$

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## Canonical Momentum & Schrödinger's Equation

$$\frac{d\mathbf{p}}{dt} = -\nabla V$$

$\frac{d}{dt}$  (canonical momentum) =  $-\nabla$  (generalized potential)

$$\mathbf{p} = m\mathbf{v} \quad \mathbf{V}$$

Here, the canonical momentum equals the kinematic momentum; and the generalized potential, the scalar potential.

$$\mathcal{E} = \frac{1}{2m} (\mathbf{p} \cdot \mathbf{p}) + V(\mathbf{r})$$

$$\hbar\omega = \frac{\hbar^2}{2m} (\mathbf{k} \cdot \mathbf{k}) + V(\mathbf{r}) \implies i\hbar \frac{\partial \psi}{\partial t} = -\frac{\hbar^2}{2m} \nabla^2 \psi + V(\mathbf{r})\psi$$

Schrödinger's Equation

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## Rules for Classical to Quantum

1. Write the classical equation of motion in terms of the canonical momentum,  $\mathbf{p}$ , and generalized potential,  $V$ :

$$\frac{d\mathbf{p}}{dt} = -\nabla V.$$

Indeed, this form identifies the precise expressions for  $\mathbf{p}$  and  $V$ .

2. Use these quantities to write the energy of the system:

$$\mathcal{E} = \frac{\mathbf{p} \cdot \mathbf{p}}{2m} + V.$$

3. Transform the classical expression into a quantum mechanical one by appealing to the Einstein-de Broglie relations. Since Schrödinger's equation is linear, these transformations are

$$\mathcal{E} = \hbar\omega \implies i\hbar \frac{\partial}{\partial t}$$

and

$$\mathbf{p} = \hbar\mathbf{k} \implies -i\hbar\nabla.$$

$$i\hbar \frac{\partial \psi}{\partial t} = -\frac{\hbar^2}{2m} \nabla^2 \psi + V(\mathbf{r})\psi$$

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## The wave function $\Psi$ : Real or Complex?

Compare the plane wave solutions for QM and E&M

$$\psi = \hat{\psi} e^{i(\mathbf{k} \cdot \mathbf{r} - \omega t)}$$

$$\mathbf{H} = \text{Re} \left\{ \widehat{\mathbf{H}} e^{i(\mathbf{k} \cdot \mathbf{r} - \omega t)} \right\}$$

$$i\hbar \frac{\partial \psi}{\partial t} = -\frac{\hbar^2}{2m} \nabla^2 \psi + V(\mathbf{r})\psi$$

$$\nabla^2 \mathbf{H} - \frac{1}{c^2} \frac{\partial}{\partial t} \mathbf{H} = 0$$

$\Psi$  must be a complex function; it is not a mathematical convenience.

$\mathbf{H}$  is a real function; it is only a mathematical convenience to consider it a complex function. The real part must be taken.

*What does it mean to have a complex wave function?*

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## The physical meaning of the wave function $\Psi$

The absolute phase of a plane wave should *not* influence the overall physics of a system.

So Max Born hypothesized in ~1927 that the square of the magnitude of the wave function  $\Psi$  was equal to the *probability* of a quantum mechanical particle to be at the location  $\mathbf{r}$  at time  $t$ .

$$\wp(\mathbf{r}, t) \equiv |\psi(\mathbf{r}, t)|^2 = \psi^*(\mathbf{r}, t)\psi(\mathbf{r}, t)$$

With the normalization condition (particle must be somewhere)

$$\int d\mathbf{r} \psi^*(\mathbf{r}, t)\psi(\mathbf{r}, t) = 1$$



## Evolution of Probability

Multiply the S-Eqn by  $\Psi^*$  and (S-Eqn)\* by  $\Psi$

$$\begin{aligned} i\hbar\psi^* \frac{\partial\psi}{\partial t} &= -\frac{\hbar^2}{2m}\psi^*\nabla^2\psi + V\psi^*\psi \\ -i\hbar\psi \frac{\partial\psi^*}{\partial t} &= -\frac{\hbar^2}{2m}\psi\nabla^2\psi^* + V\psi\psi^* \end{aligned}$$

$$\begin{aligned} i\hbar \frac{\partial}{\partial t} (\psi\psi^*) &= -\frac{\hbar^2}{2m} (\psi^*\nabla^2\psi - \psi\nabla^2\psi^*) \\ i\hbar \frac{\partial}{\partial t} (\psi\psi^*) &= \underbrace{-\frac{\hbar^2}{2m} (\psi^*\nabla^2\psi - \psi\nabla^2\psi^*)}_{\mathcal{P}} \underbrace{+ (\nabla \cdot (\psi^*\nabla\psi - \psi\nabla\psi^*))}_{\nabla \cdot \mathbf{J}_{\mathcal{P}}} \end{aligned}$$



## Probability Current

Therefore we find that the probability

$$\wp(\mathbf{r}, t) \equiv |\psi(\mathbf{r}, t)|^2 = \psi^*(\mathbf{r}, t)\psi(\mathbf{r}, t)$$

and the probability current

$$\mathbf{J}_\wp \equiv \frac{\hbar}{2im}(\psi^*\nabla\psi - \psi\nabla\psi^*) = \text{Re} \left\{ \psi^* \frac{\hbar}{im} \nabla\psi \right\}$$

satisfy a continuity relation

$$\frac{\partial \wp}{\partial t} = -\nabla \cdot \mathbf{J}_\wp$$



## Schrödinger's Equation with E&M Fields

For a charged particle, we want the classical equations such that

$$\frac{d}{dt}(\text{canonical momentum}) = -\nabla(\text{generalized potential})$$

Start with the Lorentz Force Law

$$m \frac{d\mathbf{v}}{dt} = q(\mathbf{E} + (\mathbf{v} \times \mathbf{B}))$$

and use the vector and scalar potentials

$$\mathbf{B} = \nabla \times \mathbf{A} \quad \text{and} \quad \mathbf{E} = -\frac{\partial \mathbf{A}}{\partial t} - \nabla\phi \quad \text{to find}$$



## Canonical Momentum and Energy

$$\frac{d\mathbf{p}}{dt} = -\nabla \left( \underbrace{q\phi - \frac{q}{m} \mathbf{p} \cdot \mathbf{A} + \frac{q^2}{2m} \mathbf{A} \cdot \mathbf{A}}_V \right)$$

where the canonical momentum is

$$\mathbf{p} = m\mathbf{v} + q\mathbf{A}$$

Kinematic momentum

Field momentum

The energy follows from the above to be

$$\mathcal{E} = \frac{\mathbf{p} \cdot \mathbf{p}}{2m} + \left( q\phi - \frac{q}{m} \mathbf{p} \cdot \mathbf{A} + \frac{q^2}{2m} \mathbf{A} \cdot \mathbf{A} \right)$$

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## Classical to Quantum

The energy can be written as

$$\mathcal{E} = \frac{1}{2m} (\mathbf{p} - q\mathbf{A}) \cdot (\mathbf{p} - q\mathbf{A}) + q\phi$$

The transition to quantum mechanics is done as before

$$\mathcal{E} \implies i\hbar \frac{\partial}{\partial t} \quad \text{and} \quad \mathbf{p} \implies -i\hbar \nabla$$

The S-Eqn becomes

$$i\hbar \frac{\partial \psi}{\partial t} = \frac{1}{2m} \left( \frac{\hbar}{i} \nabla - q\mathbf{A} \right)^2 \psi + q\phi \psi$$

With probability current

$$\mathbf{J}_\varphi = \text{Re} \left\{ \psi^* \left( \frac{\hbar}{im} \nabla - \frac{q}{m} \mathbf{A} \right) \psi \right\}$$

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# Macroscopic Quantum Model

1. The wave function describes the whole ensemble of superelectrons such that

$$\Psi^*(\mathbf{r}, t)\Psi(\mathbf{r}, t) = n^*(\mathbf{r}, t) \rightarrow \text{density}$$

and

$$\int d\mathbf{r} \Psi^*(\mathbf{r}, t)\Psi(\mathbf{r}, t) = N^* \rightarrow \text{Total number}$$

2. The flow of probability becomes the flow of particles, with the physical current density given by

$$\mathbf{J}_S = q^* \text{Re} \left\{ \Psi^* \left( \frac{\hbar}{im^*} \nabla - \frac{q^*}{m^*} \mathbf{A} \right) \Psi \right\}$$



## MQM cont.

3. This macroscopic quantum wavefunction follows

$$i\hbar \frac{\partial}{\partial t} \Psi(\mathbf{r}, t) = \frac{1}{2m^*} \left( \frac{\hbar}{i} \nabla - q^* \mathbf{A}(\mathbf{r}, t) \right)^2 \Psi(\mathbf{r}, t) + q^* \phi(\mathbf{r}, t) \Psi(\mathbf{r}, t)$$

Writing  $\Psi(\mathbf{r}, t) = \sqrt{n^*(\mathbf{r}, t)} e^{i\theta(\mathbf{r}, t)}$ , we find

$$\mathbf{J}_S = q^* n^*(\mathbf{r}, t) \left( \frac{\hbar}{m^*} \nabla \theta(\mathbf{r}, t) - \frac{q^*}{m^*} \mathbf{A}(\mathbf{r}, t) \right)$$

## The Supercurrent Equation



## The Supercurrent Equation

$$\mathbf{J}_s = q^* n^*(\mathbf{r}, t) \left( \frac{\hbar}{m^*} \nabla \theta(\mathbf{r}, t) - \frac{q^*}{m^*} \mathbf{A}(\mathbf{r}, t) \right)$$

$\mathbf{J}_s$  is a unique physical quantity, but  $\mathbf{A}$  and  $\theta$  are not.

If a new vector and scalar potential are found in another gauge such that

$$\mathbf{A}' \equiv \mathbf{A} + \nabla \chi \qquad \phi' \equiv \phi - \frac{\partial \chi}{\partial t}$$

Then  $\mathbf{B} = \nabla \times \mathbf{A}'$   $\mathbf{E} = -\frac{\partial \mathbf{A}'}{\partial t} - \nabla \phi'$

and  $\theta' = \theta + \frac{q^*}{\hbar} \chi$   $\mathbf{J}_s = q^* n^*(\mathbf{r}, t) \left( \frac{\hbar}{m^*} \nabla \theta'(\mathbf{r}, t) - \frac{q^*}{m^*} \mathbf{A}'(\mathbf{r}, t) \right)$

$\mathbf{B}$ ,  $\mathbf{E}$ , and  $\mathbf{J}$  are gauge invariant,  $\mathbf{A}$ ,  $\phi$ , and  $\theta$  are not.

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