

## Lecture 8

**Music and wave equations:** Spent a little time relating the 18.303 theory of the vibrating string to what you hear when you listen to a stringed instrument; scales, harmonics, transposition, timbre and the Fourier series, etcetera. (See [notes](#).) Performed a little demo on my Yamaha [guitalele](#). To obtain a chromatic scale, each fret on the guitalele (or guitar) shortens the strings by a factor of  $2^{1/12}$  (and this is why the frets get closer together as you go up the neck: they are equally spaced on a log scale).

**New topic: Separation of variables:** (See [notes](#).) This is a technique to *reduce the dimensionality* of a PDE by representing the solution as a product of lower-dimensional functions. It *only works in a handful of cases*, usually as a consequence of *symmetry*, but those cases are common enough that it is important to know them. It also gives us our only analytically solvable PDE examples in more than 1d; otherwise we will have to use the computer.

**Separation of Time:** The most important case is the one we've already done, under another name. We solved  $Au = \partial u / \partial t$  by looking for eigenfunctions  $Au = \lambda u$ , and then multiplying by  $\exp(\lambda t)$  to get the time dependence. Similarly for  $Au = \partial^2 u / \partial t^2$  except with sines and cosines. In both cases, we wrote the solution as a sum of products of purely spatial functions (the eigenfunctions) and purely temporal functions like  $\exp(\lambda t)$ . The key point here is that we aren't assuming that the *solution* is separable, only that it can be decomposed into *linear combination* of separable functions.

**Separation of Space:** Here, we try to solve problems in more than one *spatial* dimension by factoring out 1d problems in one or more dimension. In particular, we will try to find *eigenfunctions* in separable form, and then write any solution as a linear combination of eigenfunctions as usual. In practice, this mainly works only in a few important cases, especially when one direction is *translationally invariant* or when the problem is *rotationally invariant*. In the former case, translational invariance in one direction (say  $z$ ) allows us to write the eigenfunctions in separable form as  $X(x,y)Z(z)$ , where it turns out that  $Z(z) = \exp(ikz)$  for some  $k$  (and  $X$  and  $\lambda$  will then depend on  $k$ ). In the latter case, we get separable eigenfunctions  $R(r)\exp(im\theta)$  where  $m$  is an integer, in 2d, and  $R(r)Y_{l,m}(\theta,\phi)$  in 3d, where  $Y_{l,m}(\theta,\phi)$  is a [spherical harmonic](#). Also, we can *sometimes* get separable solutions for finite "box-like" domains, i.e. translationally invariant problems that have been truncated to a finite length in  $z$ .

To start with, we looked at  $\nabla^2 u = \lambda u$  in a 2d  $L_x \times L_y$  box with Dirichlet boundary conditions, and looked for separable solutions of the form  $X(x)Y(y)$ . Plugging this in and dividing by  $XY$  (the standard techniques), we get 1d eigenproblems for  $X$  and  $Y$ , and these eigenproblems ( $X'' = X \times \text{constant}$  and  $Y'' = Y \times \text{constant}$ ) just give us our familiar sine and cosine solutions. Adding in the boundary condition, we get  $\sin(n_x \pi x / L_x) \sin(n_y \pi y / L_y)$  eigenfunctions with eigenvalues  $\lambda = -(n_x \pi / L_x)^2 - (n_y \pi / L_y)^2$ . As expected, these are real and negative, and the eigenfunctions are orthogonal...giving us a 2d Fourier sine series. For example, this gives us the "normal modes" of a square drum surface.

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