

13.022 - FALL 2001

QUIZ #3

1. $\phi_{tt} + g\phi_z + \mu\phi_t = 0, z=0$

LET $\phi = \operatorname{Re} \left\{ \frac{i\omega}{\omega} e^{kz - ikx + iwt} \right\}$

BY DEFINITION IT SATISFIES THE LAPLACE EQUATION
SUBSTITUTING IN THE FREE SURFACE CONDITION
WE OBTAIN:

$$-\omega^2 + gk + i\omega\mu = 0$$

FOR KNOWN REAL ω , THE WAVENUMBER k IS
COMPLEX AND GIVEN BY

$$k = \frac{\omega^2 - i\omega\mu}{g} = k_R - ik_I$$

FROM THE DEFINITION OF THE WAVE POTENTIAL
THE COMPLEX FACTOR BECOMES

$$e^{-ikx} = e^{-ik_R x} e^{-ik_I x} \text{ WHERE } k_I = \frac{\omega\mu}{g}$$

SO THE NEW EFFECTIVE WAVE AMPLITUDE BECOMES:

$$A_{BR} = A e^{-k_I x}$$

THUS THE AMPLITUDE DUE TO BREAKING DECAYS
AT AN EXPONENTIAL RATE WITH THE DECAY
RATE $k_I = \omega\mu/g$. —

b) IN WATER OF FINITE DEPTH:

$$\phi = \operatorname{Re} \left\{ \frac{iSA}{\omega} \frac{\cosh k(z+H)}{\cosh kh} e^{-ikx+i\omega t} \right\}$$

AGAIN ϕ SATISFIES THE LAPLACE EQUATION FOR
REAL OR COMPLEX k . UPON SUBSTITUTION IN THE
FREE SURFACE CONDITION:

$$-\omega^2 + gk \tanh kh - i\omega\mu = 0$$

$$\Rightarrow gk \tanh kh = \omega^2 - i\omega\mu$$

IN SHALLOW WATER $\tanh kh \approx kh$, SO

$$gkh^2 = \omega^2 - i\omega\mu$$

$$k^2 = \frac{\omega^2 - i\omega\mu}{gh} \Rightarrow k = \left(\frac{\omega^2 - i\omega\mu}{gh} \right)^{1/2}$$

FOR SMALL VALUES OF μ , WE MAY WRITE

$$\begin{aligned} k &= \left(\frac{\omega^2}{SH} - i \frac{\omega\mu}{gH} \right)^{1/2} \\ &= \frac{\omega}{\sqrt{SH}} \left(1 - i \frac{\omega\mu}{gH} \frac{gH}{\omega^2} \right)^{1/2} \\ &= \frac{\omega}{\sqrt{gH}} \left(1 - i \frac{\mu}{\omega} \right)^{1/2} = \frac{\omega}{\sqrt{gH}} - \frac{i}{2} \frac{\mu}{\sqrt{gH}} \end{aligned}$$

UPON SUBSTITUTION IN e^{-ikx} WE OBTAIN:

$$e^{-ikx} = e^{-i \frac{\omega}{\sqrt{SH}} x} e^{-k_I x}$$

WHERE $k_I = \frac{1}{2} \frac{\mu}{\sqrt{gH}}$.

THE RATE OF DECAY IN DEEP WATER IS: $\frac{\omega\mu}{g}$
AND IS LARGER THAN THAT IN SHALLOW
WATER WHEN:

$$\frac{\omega\mu}{g} > \frac{1}{2} \frac{\mu}{\sqrt{gH}} \Rightarrow \omega > \frac{1}{2} \sqrt{\frac{g}{H}}$$

$$P = -\rho \frac{\partial \phi}{\partial t} \Rightarrow -\rho \frac{\partial \phi}{\partial t} + \lambda \frac{\partial \phi}{\partial z} = 0$$

$$-i\omega\rho \phi + \lambda \phi_z = 0, \quad z = -H$$

$$\Rightarrow \phi_z - \frac{i\omega\rho}{\lambda} \phi = 0, \quad z = -H$$

ON $z=0$ $\phi_z - (\omega^2/g - i\omega\mu) \phi = 0$

IN MORE GENERAL NOTATION:

$$\phi_z - \alpha^2 \phi = 0, \quad z = -H$$

$$\phi_z - \beta^2 \phi = 0, \quad z = 0$$

WHERE $\alpha^2 = \frac{i\omega\rho}{\lambda}$, $\beta^2 = \frac{\omega^2}{g} - i\omega\mu$.

LET ϕ BE OF THE FORM

$$\phi = \frac{isA}{\omega} f(z) e^{-ikx+i\omega t}$$

WHERE THE RELATION BETWEEN k & ω IS TO
BE DETERMINED. BY VIRTUE OF THE LAPLACE
EQUATION:

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$$f''(z) - k^2 f = 0 ; \quad -H < z < 0$$

$$f'(z) - \alpha^2 f = 0, \quad z = -H$$

$$f'(z) - \beta^2 f = 0, \quad z = 0$$

$$F(z) = A e^{kz} + B e^{-kz}$$

$$\begin{aligned} z = -H : \quad & A k e^{-kH} - B k e^{+kH} - \alpha^2 (A e^{-kH} + B e^{+kH}) \\ & = 0 \\ \Rightarrow \quad & A e^{-kH} (k - \alpha^2) = B e^{+kH} (k + \alpha^2) = 0 \end{aligned}$$

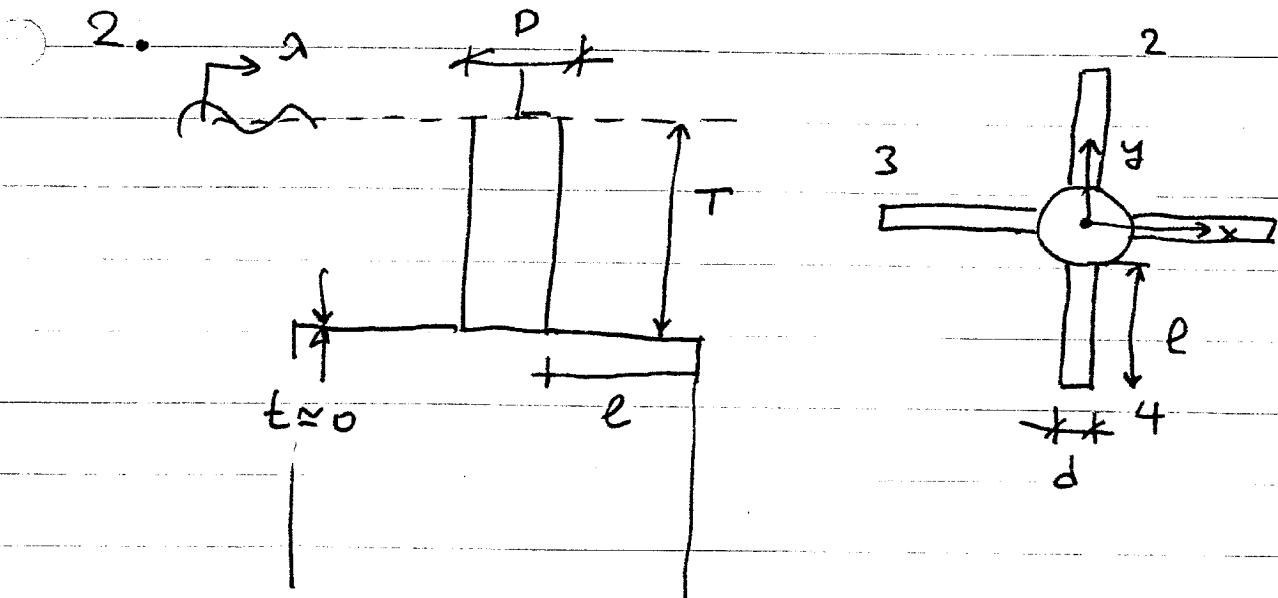
$$z = 0 : \quad A (k - \beta^2) - B (k + \beta^2) = 0$$

IT FOLLOWS THAT:

$$\frac{A}{B} = \frac{e^{kH} (k + \alpha^2)}{e^{-kH} (k - \alpha^2)} = \frac{k + \beta^2}{k - \beta^2} \Rightarrow$$

$$\frac{k + \beta^2}{k - \beta^2} = e^{2kH} \frac{k + \alpha^2}{k - \alpha^2}$$

$$\text{WITH : } \alpha^2 = \frac{i\omega e}{\lambda} \text{ AND } \beta^2 = \frac{\omega^2}{g} - i\omega\mu.$$



$$\lambda \gg D, d, \quad \lambda \sim T, l.$$

$$\phi = \operatorname{Re} \left\{ \frac{i\omega}{\omega} e^{kz - ikxt + i\omega t} \right\}$$

THE VERTICAL EXCITING FORCES ON THE TETHERS
 1, 2, 3 & 4 DEPEND UPON THE HEAVE EXCITING
 FORCE AND THE PITCH EXCITING MOMENT ON THE
 PLATFORM ABOUT THE REFERENCE COORDINATE
 SYSTEM. LET

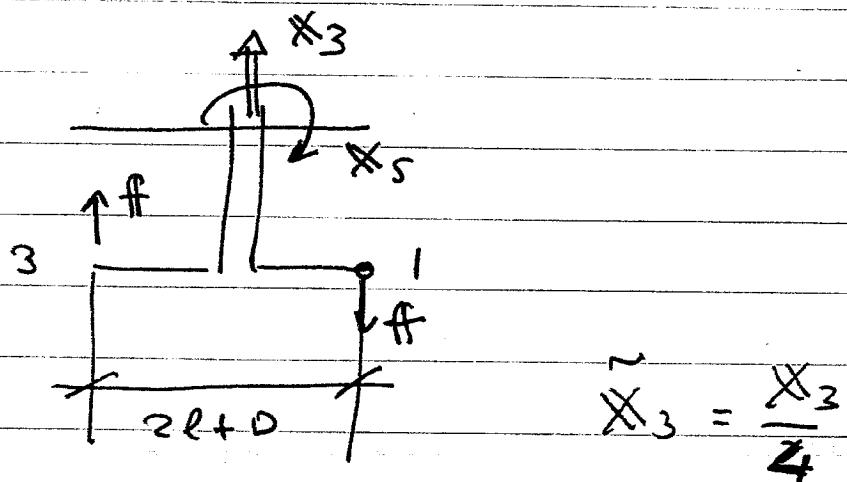
$$J(x, t) = \operatorname{Re} \left\{ A e^{-ikx + i\omega t} \right\}$$

$$F_3(t) = \operatorname{Re} \left\{ X_3 e^{i\omega t} \right\}$$

$$F_5(t) = \operatorname{Re} \left\{ X_5 e^{i\omega t} \right\}.$$

WHERE \hat{X}_3 AND \hat{X}_5 ARE THE COMPLEX HEAVE EXCITING FORCE AND PITCH EXCITING MOMENT AMPLITUDES OR RAOS. THEY WILL BE DETERMINED IN A MOMENT.

IN TERMS OF \hat{X}_3 & \hat{X}_5 THE VERTICAL EXCITING FORCES ON EACH TETHER ARE OBTAINED AS FOLLOWS:



$$\tilde{X}_3 = \frac{\hat{X}_3}{4}$$

- TETHERS 2 & 4 : $F_2(t) = \text{Re} \{ \tilde{X}_3 e^{i\omega t} \}$

THE PITCH MOMENT DOES NOT INDUCE A VERTICAL FORCE UPON TETHERS 2 & 4 TO LEADING ORDER IN THE WAVE AMPLITUDE

TETHERS 1 & 3 !

STATIC EQUILIBRIUM REQUIRES THAT THE PITCH MOMENT \hat{X}_5 IS BALANCED BY TWO FORCES

ff OF EQUAL AND OPPOSITE MAGNITUDE. WITH
THE PITCH MOMENT \ddot{X}_S DRAWN IN ^{THE} FIGURE
THE FORCES ff AS DRAWN ACT ON THE TETHERS AND

$$\ddot{X}_S = ff(2\ell + D)$$

$$\Rightarrow ff = \frac{\ddot{X}_S}{2\ell + D}$$

SO THE VERTICAL FORCE ON EACH TETHER IS:

TETHER 1: $F_1 = \tilde{X}_3 - \frac{\ddot{X}_S}{2\ell + D}, \tilde{X}_3 = \frac{\dot{X}_3}{2}$

TETHER 3: $F_2 = \tilde{X}_3 + \frac{\ddot{X}_S}{2\ell + D} .$

THERE REMAINS TO DETERMINE \dot{X}_3 & \ddot{X}_S . WE
WILL USE GI TAYLOR'S THEOREM SINCE λ IS
LARGE COMPARED TO THE CROSS-SECTIONAL
DIMENSIONS OF THE PLATFORM.

THE HEAVE EXCITING FORCE IS ARISING FROM
THE HEAVE FORCE ON EACH HORIZONTAL PLATE
AND THE HEAVE FORCE ON THE BOTTOM OF THE
PONTON MODELED AS A $\frac{1}{2}$ SPHERE

ACCORDING TO GI TAYLOR:

$$F_z = -\rho \left(\frac{a}{e} + \rho A \right) \frac{\partial P}{\partial z}$$

ON $z = -H$:

$$\begin{aligned} \frac{\partial P}{\partial z} &= \frac{\partial}{\partial z} \left[-\rho \frac{\partial \phi}{\partial t} \right] = -\rho \frac{\partial^2}{\partial z \partial t} \left\{ \frac{ig_A}{\omega} e^{kz - ikx + i\omega t} \right\} \\ &= -\rho \left\{ \frac{ig_A}{\omega} k(i\omega) e^{-ikx} e^{-kH} e^{i\omega t} \right\} \end{aligned}$$

Re

THE ADDED MASS OF A PLATE OF WIDTH d IS:

$$a_{33} \Big|_{\text{PLATE}} = \pi \rho \frac{d^2}{4}$$

ITS DISPLACEMENT IS ASSUMED TO BE ZERO FOR A SMALL THICKNESS. THE ADDED MASS OF A SPHERE IS $\frac{1}{2}$ ITS DISPLACED VOLUME

$$a_{33} \Big|_{\text{SPHERE}} = \frac{1}{2} \rho A = \frac{1}{2} \rho \frac{4}{3} \pi \left(\frac{D}{2}\right)^3$$

SO THE HEAVE EXCITING FORCE ON THE TLP BECOMES

$$F_3(t) \cancel{x_3} = -\rho \left(\frac{a_{33}}{e} + \frac{\gamma}{2} \right) \frac{\partial P}{\partial z} \Big|_{z=-H}$$

SPHERE

$$-\rho \int_{C_1 + C_2} dl \left(\frac{a_{33}}{e} \right)_{\text{PLATE}} \frac{\partial P}{\partial z} \Big|_{z=-H} = \operatorname{Re} \{ X_3 e^{i\omega t} \}$$

$+ C_3 + C_4$

WHERE THE CONTOUR INTEGRALS ARE OVER THE AXES (RECTILINEAR) OF THE FOUR HORIZONTAL PLATES.

THE PITCH MOMENT $F_S(t) = \operatorname{Re} \{ X_5 e^{i\omega t} \}$ WILL BE CONTRIBUTED FROM THE TWO HORIZONTAL PLATES 1 & 3 IN LINE WITH THE DIRECTION OF PROPAGATION OF THE WAVES. HENCE

$$F_S(t) = -\rho \int_{C_1 + C_2} dx \left(\frac{a_{33}}{e} \right)_{\text{PLATE}} \frac{\partial P}{\partial z} \Big|_{z=-H} (-x)$$

WHERE $(-x)$ IS THE MOMENT ARM OF THE DIFFERENTIAL ELEMENT dx ALONG PLATES 1 & 3. —

C) HEAVE IS RESTORED EQUILIBRIUM BY ALL FOUR TETHERS. SO THE EFFECTIVE SPRING CONSTANT IS :

$$C_{33} = \rho g A_w + 4k$$

WHERE $A_w = \pi D^2/4$ WITH D BEING THE DIAMETER OF THE VERTICAL CYLINDER

THE EFFECTIVE MASS IS GIVEN BY THE MASS OF THE PLATFORM, THE HEAVE ADDED MASS OF THE 4 HORIZONTAL PLATES + $\frac{1}{2}$ SPHERES, AND THE EFFECTIVE INERTIA OF EACH TETHER:

$$M_{TOT} = A_{33} + M + 4m$$

WHERE $A_{33} = 4A_{PLATE} + \frac{1}{2} A_{SPHERE}$

$$A_{PLATE} = l \frac{\pi d^2}{4}, A_{SPHERE} = \frac{1}{2} \nabla = \frac{1}{2} \frac{4}{3} \pi \left(\frac{D}{2}\right)^3$$

SO THE HEAVE NATURAL PERIOD OF THE TLP IS:

$$\omega^2 = \frac{\rho g A_w + 4k}{A_{33} + M + 4m}.$$