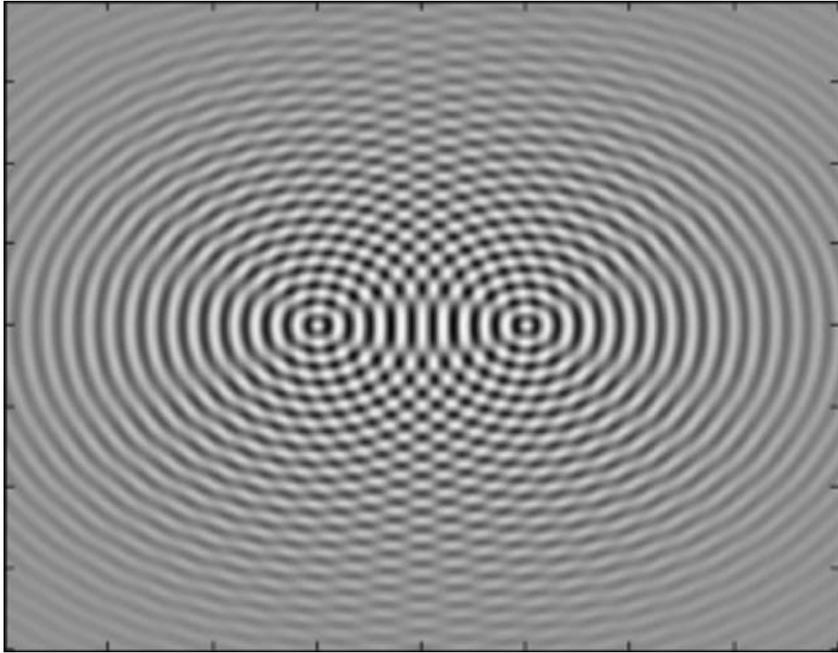
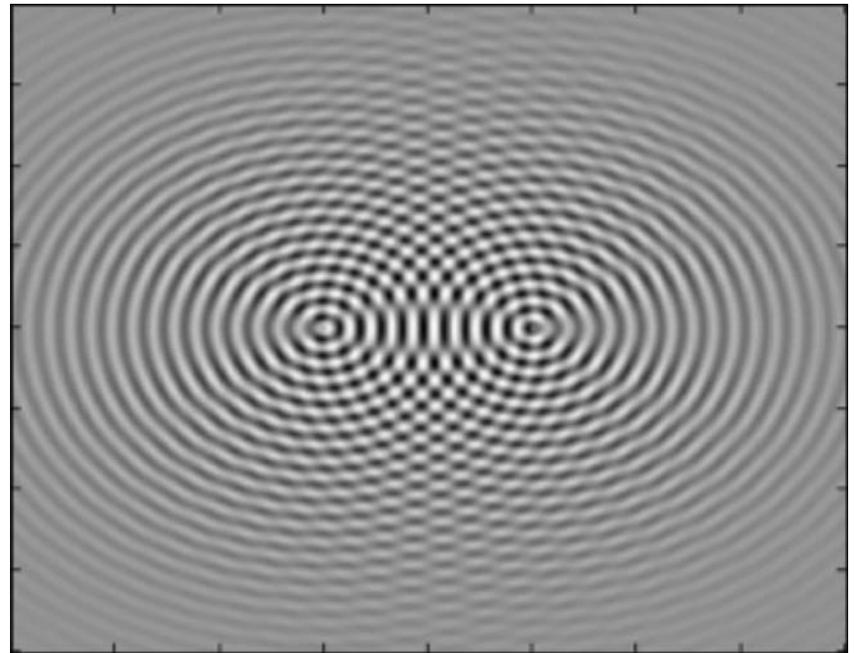


Phase delays and interference



in phase



out of phase

The 1D wave equation

Consider a moving disturbance with envelope ψ of the form

$$f(z, t) = \psi(z \pm ct).$$

The “-” sign denotes a wave moving to the right (forward), whereas “+” denotes a wave moving to the left (backwards.)

Denoting by ψ' the envelope’s derivative with respect to its argument, and taking the first partial derivatives with respect to z and t ,

$$\frac{\partial f}{\partial z} = \psi' \quad \frac{\partial f}{\partial t} = \pm c\psi' \Rightarrow \frac{\partial f}{\partial z} \mp \frac{1}{c} \frac{\partial f}{\partial t} = 0.$$

Taking second derivatives,

$$\frac{\partial^2 f}{\partial z^2} = \psi'' \quad \frac{\partial^2 f}{\partial t^2} = c^2\psi'' \Rightarrow \frac{\partial^2 f}{\partial z^2} - \frac{1}{c^2} \frac{\partial^2 f}{\partial t^2} = 0.$$

This is referred to as the **Helmholtz Wave Equation**, or simply the **Wave Equation**.

We can solve the wave equation for specific envelopes with proper initial and boundary conditions, e.g.

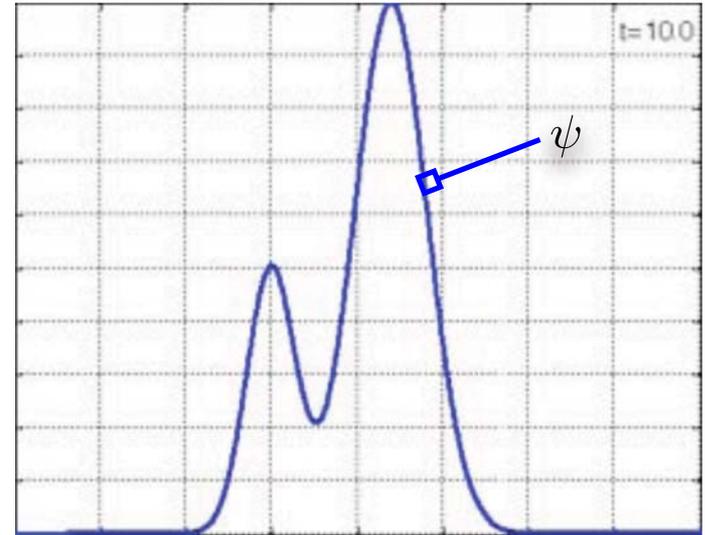
$$f(z, t = 0) = \psi(z) = a_0 \cos(kz),$$

where k is a shorthand for $2\pi/\lambda$ and λ is the spatial period (wavelength.)

We can easily see that

$$f(z, t) = a_0 \cos(k(z \pm ct))$$

satisfy the Wave Equation and the initial condition, so they are solutions propagating forward and backward, respectively.



The solution is often written in equivalent forms, as

$$f(z, t) = a_0 \cos(k(z \pm ct)) = a_0 \cos(kz \pm \omega t) = a_0 \cos\left(\frac{2\pi}{\lambda}(z \pm ct)\right).$$

These are consistent given the definitions

$$\omega = 2\pi\nu, \quad k = \frac{2\pi}{\lambda}, \quad \text{and the dispersion relation } c = \lambda\nu.$$

Linear superposition of waves

$$\frac{\partial^2 f}{\partial z^2} = \frac{1}{c^2} \frac{\partial^2 f}{\partial t^2}.$$

This is a *linear* partial differential equation. We can see easily that if f_1 is a solution and f_2 is another solution, then $a_1 f_1 + a_2 f_2$ is also a solution for arbitrary constants a_1, a_2 .

For example, $a_- \psi(z - ct) + a_+ \psi(z + ct)$,

a superposition of a forward and backward propagating waves, is also a solution to the Wave Equation.

It is known as a **standing wave**.

Let us derive the standing wave for the case of forward and backward waves of equal amplitudes.

The superposition is

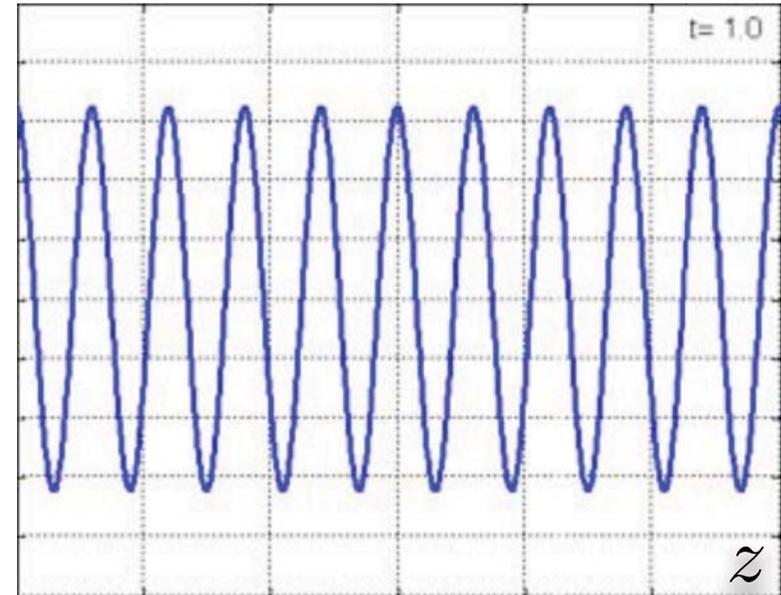
$$f(z, t) = a \cos(kz - \omega t + \phi) + a \cos(kz + \omega t + \phi).$$

At this point we must recall the trigonometric identity

$$\cos A + \cos B = 2 \cos \frac{A+B}{2} \cos \frac{A-B}{2}.$$

Using the identity, we rewrite the superposition as

$$\begin{aligned} f(z, t) &= 2a \cos \frac{(kz - \omega t + \phi) + (kz + \omega t + \phi)}{2} \times \\ &\quad \times \cos \frac{(kz - \omega t + \phi) - (kz + \omega t + \phi)}{2} \\ &= 2a \cos(kz + \phi) \cos(\omega t). \end{aligned}$$



$$f(z, t) = 2a \cos(kz + \phi) \cos(\omega t)$$

is a spatial sinusoid of period $\lambda = 2\pi/k$, phase delay ϕ and is also oscillating with temporal period $T = 2\pi/\omega$.

It is *stationary*, that is non-propagating.

Hence the term **standing wave**.

Complex (phasor) representation of waves

Consider two simple wave forms with amplitude a , wave number k , angular frequency ω , and phase delay ϕ :

$$f_c(z, t) = a \cos(kz - \omega t + \phi), \quad f_s(z, t) = a \sin(kz - \omega t + \phi).$$

These are both solutions to the Wave Equation (but with different initial conditions.) Hence, their superposition is also a solution to the Wave Equation. We form the following special superposition

$$\begin{aligned} \hat{f}(z, t) &= f_c(z, t) + i f_s(z, t) \\ &= a \exp \{ i(kz - \omega t + \phi) \}. \end{aligned}$$

This is the **complex representation** of the wave. The complex exponential provides immense mathematical convenience, as we will see, but it is important to remember that *only its real part has physical significance*. Further acknowledging that in linear media the temporal frequency ω of the wave does not change, it is common to drop the $e^{-i\omega t}$ term, resulting in

$$a \exp \{ i(kz + \phi) \}.$$

This reduced complex representation is sometimes referred to as **phasor**.

Solution to the standing wave using phasors

Recall the superposition of forward and backward waves resulting in the standing wave,

$$f(z, t) = a \cos(kz - \omega t + \phi) + a \cos(kz + \omega t + \phi).$$

Some care needs to be taken in deriving the phasor of the backward wave. We need to get rid of a term of the form $e^{-i\omega t}$ term, so we write

$$a \cos(kz + \omega t + \phi) = a \cos(-kz - \omega t - \phi) = \operatorname{Re} \left\{ e^{-i(kz + \omega t + \phi)} \right\}$$

Therefore, the proper phasor for the backward propagating wave is

$$a \exp \{ -ikz - \phi \}.$$

The superposition is now written as

$$ae^{i(kz + \phi)} + ae^{-i(kz + \phi)} = 2a \cos(kz + \phi).$$

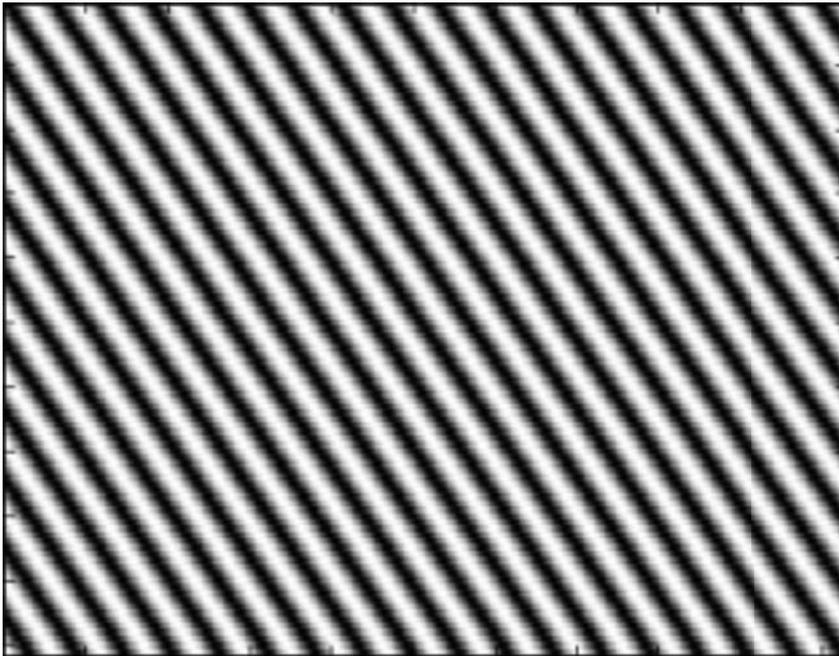
This becomes identical to our earlier result on the standing wave after we put back the $\cos(\omega t)$ term that is implicit in the phasor notation; however, the derivation was relatively painless, without need for the trigonometric formulae.

The 3D wave equation

In three-dimensions, the Wave Equation is generalized as

$$\frac{\partial^2 f}{\partial x^2} + \frac{\partial^2 f}{\partial y^2} + \frac{\partial^2 f}{\partial z^2} - \frac{1}{c^2} \frac{\partial^2 f}{\partial t^2} = 0.$$

Our familiar plane and spherical waves are special solutions.

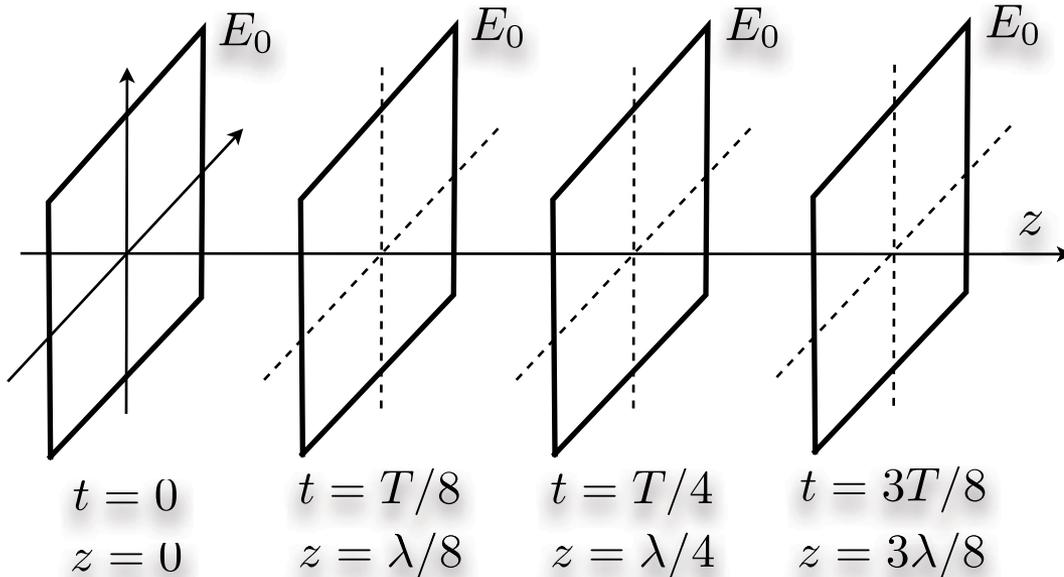


Plane wave



Spherical wave

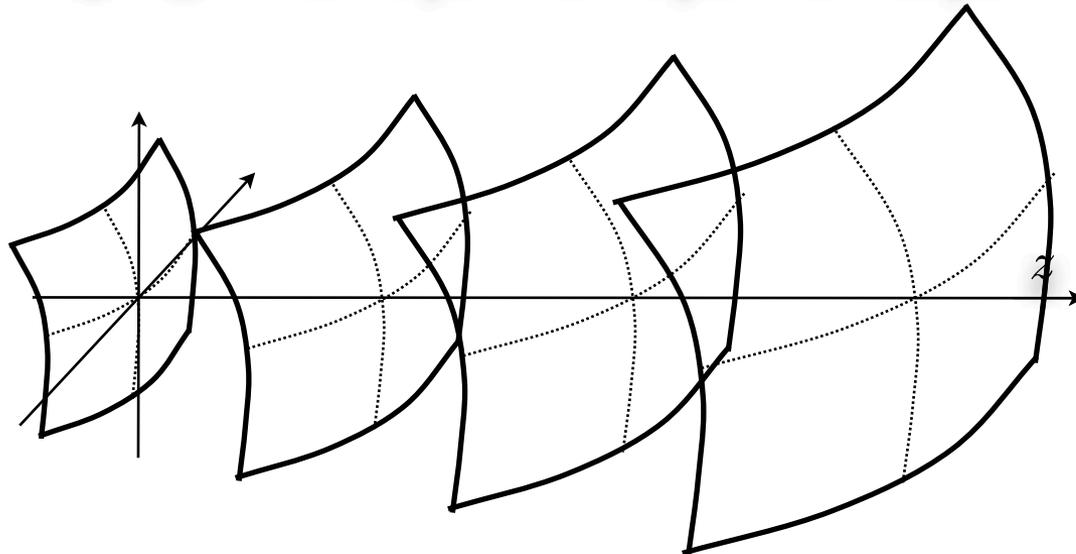
Planar and Spherical Wavefronts



Planar wavefront (plane wave):

The wave phase is constant along a **planar** surface (the wavefront).

As time evolves, the wavefronts propagate at the wave speed without changing; we say that the wavefronts are *invariant to propagation* in this case.

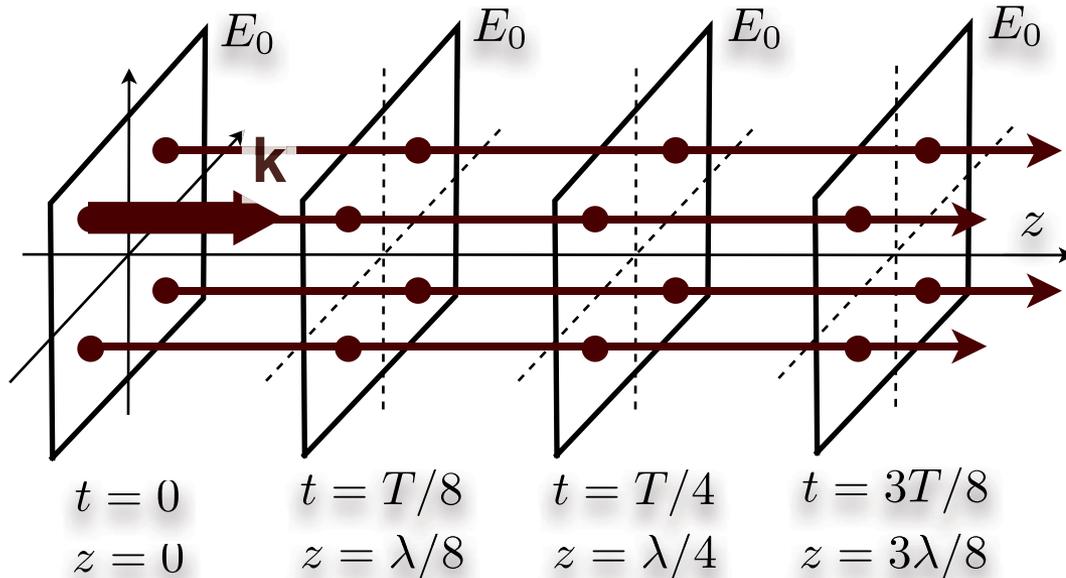


Spherical wavefront (spherical wave):

The wave phase is constant along a **spherical** surface (the wavefront).

As time evolves, the wavefronts propagate at the wave speed and expand outwards while preserving the wave's energy.

Wavefronts, rays, and wave vectors



Rays are:

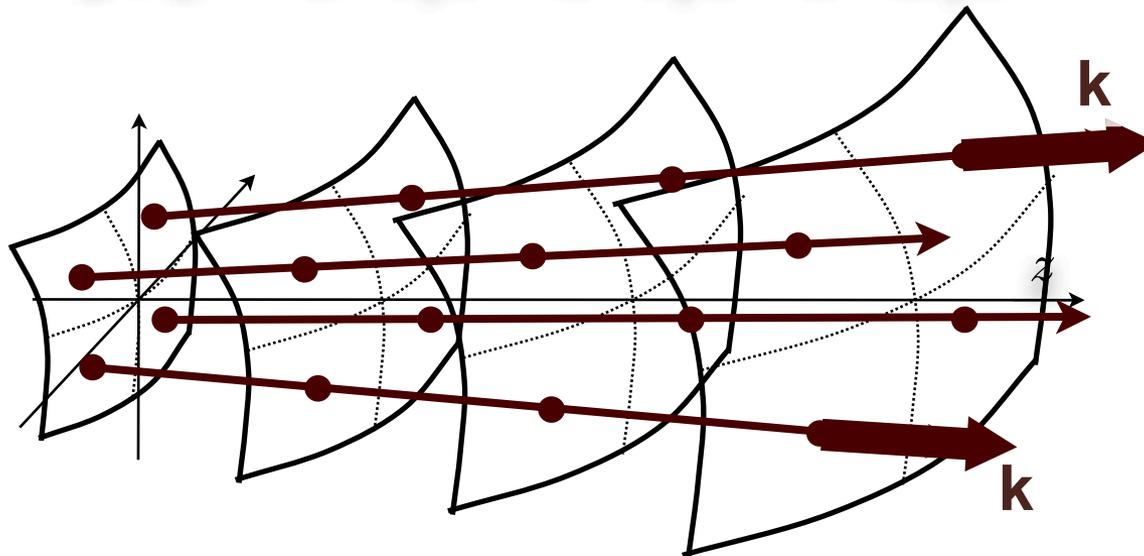
- 1) normals to the wavefront surfaces
- 2) trajectories of "particles of light"

Wave vectors:

At each point on the wavefront, we may assign a normal vector \mathbf{k}

This is known as the wave vector; its magnitude k is the wave number and it is defined as

$$k \equiv |\mathbf{k}| = \frac{2\pi}{\lambda}$$



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