

22.01 Fall 2015, Quiz 2 Study Sheet

November 14, 2015

Familiarize yourself with all these concepts from an *intuitive* and a *quantitative* point of view. Keep in mind that the material we covered in class, especially in Friday's recitation, will be particularly relevant.

1 Photon Interactions with Matter

- Three main mechanisms of interaction
 - Photoelectric Effect
 - Compton Scattering
 - Pair Production
- Mechanisms and Energetics
 - Photoelectric Effect
 - * Photon is absorbed by an electron, causes an ejection
 - * $T_{e^-} = \hbar\omega - E_{binding}$
 - * Usually an inner shell electron
 - Level 1 electrons about 80% of the time
 - * $\sigma_{\tau} \propto Z^5 \left(\frac{1}{\alpha}\right)^{\frac{7}{2}}$
 - Strongly more likely with increasing atomic number
 - Strongly less likely with increasing photon energy
 - Compton Scattering
 - * Photon gets scattered by an electron, leaves at a lower energy and at a different angle
 - * Key relationships: See Equations 10.6 - 10.9 in Yip
 - Equation 10.6 - Shows the *increase* in photon wavelength
 - Equation 10.7 - Relates the change in energy to photon energy & angle
 - Equation 10.8 - Gives the energy of the Compton electron
 - Equation 10.9 - Relates the exit angle of the scattered photon and the Compton electron
 - * σ_C : See Equation 10.17 (Klein-Nishina cross section)
 - **Know what this means in terms of forward scattering bias**
 - * Understand the *Compton Edge* - Backscattered photon is more likely (see Klein Nishina cross section), they impart the most energy to the Compton electron
 - * It is the Compton *electrons* that are counted in detectors, along with any other ionized electrons from any process
 - Pair Production
 - * Photon above 1.022MeV gets near an electron cloud
 - * Spontaneously changes into an electron/positron pair

- * Electron energy is counted by an ionization cascade
- * Positron moves some distance in a material, annihilates with another electron
- * Gives off two 511keV gammas in opposite directions
- * Single escape peak - One gets out of the detector, one is reabsorbed (medium-sized detectors)
- * Double escape peak - Both get out of the detector (small-sized detectors)
- * No escape peaks - All photons get absorbed (large detectors)
- * 511keV peak - Pair production happened somewhere *outside* the detector, then a 511keV photon enters the detector and undergoes photoelectric or Compton scattering
- See Equations 10.43 - 10.45 for cross section comparisons
- Mass attenuation coefficients - density-normalized photon interaction probabilities

$$I = I_0 e^{-\left(\frac{\mu}{\rho}\right)\rho x} \quad (1)$$

- **Know how to interpret photon spectra from detectors** - see Figure 10.18 for an example

- Ion/Electron Interactions

- Four methods: elastic & inelastic collisions, with other electrons and with nuclei
 - * Elastic collisions with electrons are the main mechanism of energy loss
- Derived using the hollow cylinder approach (see pp. 225-226) as a Coulomb force balance & integration
- $p_{e^-} = \int F_y(t) dt = m \int a(t) dt = mv$
- Stopping power formula: $-\frac{dT}{dx} = \frac{4\pi N Z_1^2 Z_2 e^4}{m_e v^2} \ln\left(\frac{2m_e v^2}{I}\right) = \frac{4\pi N Z_1^2 Z_2 e^4}{E_i} \frac{M}{m_e} \ln\left(\frac{\gamma_e E_i}{I}\right)$; $\gamma_s = \frac{4m_e M}{(M+m_e)^2}$
- Relativistic correction: $-\frac{dT}{dx} = \frac{4\pi N Z_1^2 Z_2 e^4}{m_e v^2} \left[\ln\left(\frac{2m_e v^2}{I}\right) - \ln(1 - \beta^2) - \beta^2 \right]$
- Stopping power curve shape: follows 1/E times ln(E), see & understand Figure 11.4
- Number of ion pairs produced: $i = \frac{1}{W} \left(-\frac{dT}{dx}\right)$
- Cross section comparison - see Equations 11.18 - 11.21 (p. 233)
- Know when ionization loss is high (low energy) and radiation loss is high (very high energy)
- Range: When the ions come to rest - see p. 235 for full explanation

- Neutron interactions

- No charge, almost all nuclear elastic collisions
 - * Never forget the Q-equation! See p. 143 to refresh your memory
 - * Assumptions:
 - Q=0 (elastic scattering)
 - M1=M3=1 amu, M2=M4=A amu
 - A neutron can lose at most $(1 - \alpha)$ of its energy in an elastic collision, $\alpha = \left(\frac{A-1}{A+1}\right)^2$
 - We assume that scattering is isotropic in angle, this is less true with increasing energy
- Inelastic Scattering
 - * One neutron goes in, *compound nucleus* formed, different neutron comes out
 - * See p. 244 for energy level diagram, p. 246 for diagram with resonance with nuclear energy levels
 - * Often $Q < 0$, cross section is zero until $E=Q$
- Fission

- * Also a compound nucleus formation, but this time the “liquid drop” splits into two uneven sized droplets
- * First, the neutron rich fission products shed 1-2 neutrons each
- * Then, beta decay reduces their assymetry further
- * See p. 255 (Figure 12.11) for a time diagram of what happens

- Neutron Transport

- Balance of neutron gains & losses from a location (dV), in a certain energy range (dE), going in a certain solid angle direction (dΩ)
- Gains
 - * Fission
 - * External sources
 - * Scattering down into our energy group
- Losses
 - * Scattering down out of our energy group
 - * Absorption
 - * Leakage
- Terminology
 - * $\phi(E)$: Angular flux, $\Phi(E)$: Total flux
 - * $\Sigma_t = \Sigma_a + \Sigma_s$ (total = absorption + fission)
 - * $\Sigma_a = \Sigma_f + \Sigma_\gamma$ (absorption = fission + capture)
 - * $\Sigma_s = \Sigma_{s,el} + \Sigma_{s,non-el}$ (scattering = elastic + inelastic)
 - * E is our energy, E' is some other energy
 - * Ω is our angle, Ω' is some other angle
- Equational form:

$$\begin{aligned} \frac{dn(\mathbf{r}, E, \mathbf{\Omega}, t)}{dt} &= \frac{\nu\chi(E)}{4\pi} \int_V \int_E \int_{\Omega'} d^3r dE' d\Omega' \Sigma_f(E') \phi(\mathbf{r}, E', \mathbf{\Omega}, t) \\ &+ S_0(\mathbf{r}, E, \mathbf{\Omega}, t) + \int_V \int_E \int_{\Omega'} d^3r dE d\Omega \Sigma_s(E') \phi(\mathbf{r}, E, \mathbf{\Omega}', t) F(E' \rightarrow E, \Omega' \rightarrow \Omega) \\ &- \int_V d^3r dE d\Omega \Sigma_t(E') \phi(\mathbf{r}, E, \mathbf{\Omega}, t) - \int_V d^3r dE d\Omega \cdot \nabla \phi(\mathbf{r}, E, \mathbf{\Omega}, t) \end{aligned} \quad (2)$$

- Terms: change = fission + external + scattering in - all collisions - flow outwards
- Simplifications: make everything isotropic (eliminate angular dependence), any volume element is the same (eliminate r-dependence)
- Big simplification: Assume all neutrons are in one energy group, homogeneous (perfectly mixed) reactor with one average material
- New equation: gains = losses

$$\Sigma_a \Phi - \nabla D \cdot \nabla \Phi = \frac{\nu \Sigma_f \Phi}{k_{eff}} \quad (3)$$

- Solution takes form of $\frac{-\nabla^2 \Phi}{\Phi} = constants$; *sin & cos*
- Eliminate sin due to symmetry concerns
- Solution is $\frac{-\nabla^2 \Phi}{\Phi} = B^2$; B is the “buckling”

* Material buckling & criticality:

$$\Sigma_a \Phi - \nabla D \cdot \nabla \Phi = \frac{\nu \Sigma_f \Phi}{k_{eff}} \implies \Sigma_a \Phi - DB^2 \Phi = \frac{\nu \Sigma_f \Phi}{k_{eff}} \implies k_{eff} = \frac{\nu \Sigma_f}{\Sigma_a + DB^2} \quad (4)$$

* k_{eff} is a balance between neutron production & destruction

* If $k_{eff} = 1$, the reactor is critical (perfect balance between production & loss)

* Then, using Equation 4, we can give a criticality condition based on the materials in a reactor:

$$B_{material}^2 = \frac{\nu \Sigma_f - \Sigma_a}{D} \quad (5)$$

* and its geometry:

$$\frac{-\nabla^2 \Phi}{\Phi} = B_{geometry}^2; \quad \Phi = A \cos(B_g x) \quad (6)$$

* The reactor is perfectly critical when

$$B_g^2 = B_m^2 \quad (7)$$

– Another criticality condition: the six-factor formula

$$k_{eff} = \eta f p \epsilon P_{FNL} P_{TNL} \quad (8)$$

* Assumes two groups (fast & thermal) of neutron energies

* η - average number of neutrons per fission

* f - thermal utilization factor - probability that an absorption takes place in the fuel and not elsewhere

* p - resonance escape probability, the probability that a neutron slows down through and out of the resonance region of cross sections into the thermal energies

* ϵ - fast fission factor (not all fissions are thermal)

* P_{FNL} - probability that a given fast neutron stays in the reactor (not that high)

* P_{TNL} - probability that a given thermal neutron stays in the reactor (pretty high)

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22.01 Introduction to Nuclear Engineering and Ionizing Radiation
Fall 2015

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