

Topic 5

Updated Lagrangian Formulation for Incremental General Nonlinear Analysis

Contents:

- Principle of virtual work in terms of 2nd Piola-Kirchhoff stresses and Green-Lagrange strains referred to the configuration at time t
- Incremental stress and strain decompositions in the updated Lagrangian form of the principle of virtual work
- Linear and nonlinear strain increments
- Consistent linearization of terms in the principle of virtual work
- The “out-of-balance” virtual work term
- Iterative equations for modified Newton-Raphson solution
- Flow chart of complete solution
- Comparison to total Lagrangian formulation

Textbook:

Section 6.2.3

A SUMMARY OF THE T.L.F

• THE BASIC EQN. WE USE IS

$$\int_{t+\Delta t V} \tau_{ij} \delta_{t+\Delta t} \epsilon_{ij} dV = {}^{t+\Delta t} R$$

• WE INTRODUCE

$$\begin{matrix} {}^{t+\Delta t} \sigma_{ij} & {}^{t+\Delta t} \tau_{ij} \\ {}^{t+\Delta t} \epsilon_{ij} & {}^{t+\Delta t} \epsilon_{ij} \end{matrix}$$

• WE DECOMPOSE

$$\begin{matrix} {}^{t+\Delta t} \sigma_{ij} = {}^t \sigma_{ij} + {}^z \sigma_{ij} \\ {}^{t+\Delta t} \epsilon_{ij} = {}^t \epsilon_{ij} + {}^z \epsilon_{ij} \end{matrix}$$

• WE NOTE

$$\epsilon_{ij} = \epsilon_{ij} + \eta_{ij}$$

LINEAR / NONLINEAR

IN $u_i \leftarrow$ particle displ.

• WE OBTAIN

$$\int_{0V} {}^{t+\Delta t} \sigma_{ij} \delta {}^{t+\Delta t} \epsilon_{ij} dV = {}^{t+\Delta t} R$$

• SUBSTITUTION AND LINEARIZATION GIVES

$$\begin{aligned} & \int_{0V} c_{ijrs} \epsilon_{rs} \delta \epsilon_{ij} dV \\ & + \int_{0V} {}^t \sigma_{ij} \delta \eta_{ij} dV \\ & = {}^{t+\Delta t} R - \int_{0V} {}^t \sigma_{ij} \delta \epsilon_{ij} dV \end{aligned}$$

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• IN THE ITERATION WE HAVE

$$\dots = {}^{t+\Delta t} R - \int_{0V} {}^{t+\Delta t (k-1)} \sigma_{ij} \delta {}^{t+\Delta t (k-1)} \epsilon_{ij} dV$$

$k = 1, 2, 3, \dots$

• THE F.E. DISCRETIZATION GIVES

$${}^t \underline{K} \Delta \underline{u}^{(i)} = {}^{t+\Delta t} \underline{R} - {}^{t+\Delta t} \underline{F}^{(i-1)}$$

$i = 1, 2, 3, \dots$

• AT CONVERGENCE

$${}^{t+\Delta t} \underline{R} = {}^{t+\Delta t} \underline{F}$$

WE SATISFY :

- COMPATIBILITY
- STRESS-STRAIN LAW
- EQUILIBRIUM / NODAL POINT EQUILIBRIUM
- LOCAL EQUILIBRIUM IF MESH IS FINE ENOUGH

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UPDATED LAGRANGIAN FORMULATION

Because ${}^{t+\Delta t}S_{ij}$ and ${}^{t+\Delta t}\epsilon_{ij}$ are energetically conjugate,

the principle of virtual work

$$\int_{{}^{t+\Delta t}V} {}^{t+\Delta t}\tau_{ij} \delta {}^{t+\Delta t}e_{ij} {}^{t+\Delta t}dV = {}^{t+\Delta t}\mathcal{R}$$

can be written as

$$\int_{{}^tV} {}^{t+\Delta t}S_{ij} \delta {}^{t+\Delta t}\epsilon_{ij} {}^tdV = {}^{t+\Delta t}\mathcal{R}$$

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We already know the solution at time t (${}^tS_{ij}$, ${}^tu_{i,j}$, etc.). Therefore we decompose the unknown stresses and strains as

$${}^{t+\Delta t}S_{ij} = \underbrace{{}^tS_{ij}}_{\text{known}} + \underbrace{{}^tS_{ij}}_{\text{unknown increments}} = {}^t\tau_{ij} + {}^tS_{ij}$$

$${}^{t+\Delta t}\epsilon_{ij} = \underbrace{{}^t\epsilon_{ij}}_{\text{known}} + \underbrace{{}^t\epsilon_{ij}}_{\text{unknown increments}} = {}^t\epsilon_{ij}$$

↙
0

In terms of displacements, using

$${}^{t+\Delta t}{}_{t}\epsilon_{ij} = \frac{1}{2} \left({}^{t+\Delta t}{}_{t}u_{i,j} + {}^{t+\Delta t}{}_{t}u_{j,i} + {}^{t+\Delta t}{}_{t}u_{k,i} {}^{t+\Delta t}{}_{t}u_{k,j} \right)$$

we find

$${}_{t}\epsilon_{ij} = \underbrace{\frac{1}{2} ({}_{t}u_{i,j} + {}_{t}u_{j,i})}_{\text{linear in } u_i} + \underbrace{\frac{1}{2} {}_{t}u_{k,i} {}_{t}u_{k,j}}_{\text{nonlinear in } u_i}$$

(No initial displacement effect)

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We define

$${}_{t}e_{ij} = \frac{1}{2} ({}_{t}u_{i,j} + {}_{t}u_{j,i}) \quad \text{linear strain increment}$$

$${}_{t}\eta_{ij} = \frac{1}{2} {}_{t}u_{k,i} {}_{t}u_{k,j} \quad \text{nonlinear strain increment}$$

Hence

$$\begin{aligned} {}_{t}\epsilon_{ij} &= {}_{t}e_{ij} + {}_{t}\eta_{ij} \\ \delta {}_{t}\epsilon_{ij} &= \delta {}_{t}e_{ij} + \delta {}_{t}\eta_{ij} \end{aligned}$$

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The equation of the principle of virtual work becomes

$$\int_{t_V} {}^t\mathbf{S}_{ij} \delta_t \epsilon_{ij} {}^t dV + \int_{t_V} {}^t\boldsymbol{\tau}_{ij} \delta_t \eta_{ij} {}^t dV \\ = {}^{t+\Delta t} \mathcal{R} - \int_{t_V} {}^t\boldsymbol{\tau}_{ij} \delta_t \mathbf{e}_{ij} {}^t dV$$

Given a variation δu_i , the right-hand-side is known. The left-hand-side contains unknown displacement increments.

Important: So far, no approximations have been made.

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Just as in the total Lagrangian formulation,

- The equation of the principle of virtual work is in general a complicated nonlinear function in the unknown displacement increment.
- Therefore we linearize this equation to obtain the approximate equation

$$\underline{{}^t\mathbf{K}} \Delta \underline{\mathbf{U}} = {}^{t+\Delta t} \underline{\mathbf{R}} - \underline{{}^t\mathbf{F}}$$

We begin to linearize the terms containing the unknown displacement increments.

1) The term $\int_{tV} {}^t\tau_{ij} \delta_t \eta_{ij} {}^t dV$

is linear in u_i .

- ${}^t\tau_{ij}$ does not contain u_i .
- $\delta_t \eta_{ij} = \frac{1}{2} {}^t u_{k,i} \delta_t u_{k,j} + \frac{1}{2} \delta_t u_{k,i} {}^t u_{k,j}$
is linear in u_i .

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2) The term $\int_{tV} {}^t S_{ij} \delta_t \epsilon_{ij} {}^t dV$ contains

linear and higher-order terms in u_i .

- ${}^t S_{ij}$ is a nonlinear function (in general) of ${}^t \epsilon_{ij}$.
- $\delta_t \epsilon_{ij} = \delta_t e_{ij} + \delta_t \eta_{ij}$ is a linear function of u_i .

We need to neglect all higher-order terms in u_i .

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${}^tS_{ij}$ can be written as a Taylor series in ${}^t\varepsilon_{ij}$:

$${}^tS_{ij} = \underbrace{\frac{\partial {}^tS_{ij}}{\partial {}^t\varepsilon_{rs}} \Big|_t}_{\text{known}} \underbrace{{}^t\varepsilon_{rs}}_{\substack{\text{linear and} \\ \text{quadratic in } u_i}} + \text{higher-order terms}$$

$$\doteq \frac{\partial {}^tS_{ij}}{\partial {}^t\varepsilon_{rs}} \Big|_t \left(\underbrace{{}^te_{rs}}_{\substack{\text{linear} \\ \text{in } u_i}} + \underbrace{{}^t\eta_{rs}}_{\substack{\text{quadratic} \\ \text{in } u_i}} \right) \doteq \underbrace{{}^tC_{ijrs}}_{\substack{\text{linearized term} \\ \text{in } u_i}} {}^te_{rs}$$

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Hence we obtain

$$\underbrace{{}^tS_{ij}}_{\text{known}} \delta {}^t\varepsilon_{ij} \doteq \underbrace{{}^tC_{ijrs} {}^te_{rs}}_{\text{linear in } u_i} (\delta {}^te_{ij} + \delta {}^t\eta_{ij})$$

$$= \underbrace{{}^tC_{ijrs} {}^te_{rs}}_{\substack{\text{linear in } u_i \\ \text{does not} \\ \text{contain } u_i}} \underbrace{\delta {}^te_{ij}}_{\text{linear in } u_i} + \underbrace{{}^tC_{ijrs} {}^te_{rs}}_{\text{quadratic in } u_i} \underbrace{\delta {}^t\eta_{ij}}_{\text{linear in } u_i}$$

$$\doteq \underbrace{{}^tC_{ijrs} {}^te_{rs} \delta {}^te_{ij}}_{\text{linearized result}}$$

The final linearized equation is

$$\int_{tV} {}_t\mathbf{C}_{ijrs} {}_t\mathbf{e}_{rs} \delta {}_t\mathbf{e}_{ij} {}^t dV + \int_{tV} {}^t\boldsymbol{\tau}_{ij} \delta {}_t\eta_{ij} {}^t dV$$

$$\underline{\delta \mathbf{U}}^T \underline{{}^t\mathbf{K}} \underline{\Delta \mathbf{U}}$$

$$= \underline{{}^{t+\Delta t}\mathcal{R}} - \int_{tV} {}^t\boldsymbol{\tau}_{ij} \delta {}_t\mathbf{e}_{ij} {}^t dV$$

$$\underline{\delta \mathbf{U}}^T (\underline{{}^{t+\Delta t}\mathbf{R}} - \underline{{}^t\mathbf{F}})$$

when discretized using the finite element method

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An important point is that

$$\int_{tV} {}^t\boldsymbol{\tau}_{ij} \delta {}_t\mathbf{e}_{ij} {}^t dV$$

is the virtual work due to element internal stresses at time t . We interpret

$$\underline{{}^{t+\Delta t}\mathcal{R}} - \int_{tV} {}^t\boldsymbol{\tau}_{ij} \delta {}_t\mathbf{e}_{ij} {}^t dV$$

as an “out-of-balance” virtual work term.

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Solution using updated Lagrangian formulation

Displacement iteration:

$${}^{t+\Delta t}\underline{u}_i^{(k)} = {}^{t+\Delta t}\underline{u}_i^{(k-1)} + \Delta \underline{u}_i^{(k)}, \quad {}^{t+\Delta t}\underline{u}_i^{(0)} = {}^t\underline{u}_i$$

Modified Newton iteration:

$$\begin{aligned} & \int_{V} {}^t C_{ijrs} \Delta {}^t e_{rs}^{(k)} \delta {}^t e_{ij} {}^t dV + \int_{V} {}^t T_{ij} \delta \Delta {}^t \eta_{ij}^{(k)} {}^t dV \\ & = {}^{t+\Delta t} \mathcal{R} - \int_{V} {}^{t+\Delta t} T_{ij}^{(k-1)} \delta {}^{t+\Delta t} e_{ij}^{(k-1)} {}^{t+\Delta t} dV \end{aligned}$$

$$k = 1, 2, \dots$$

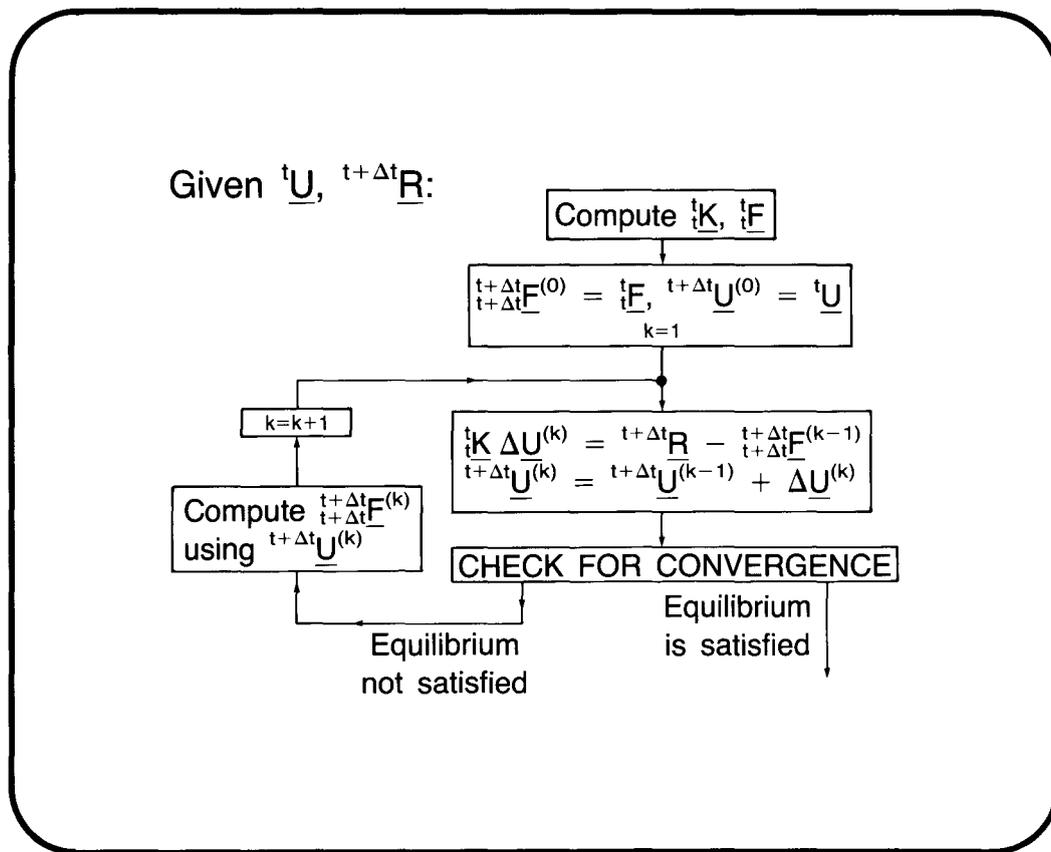
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which, when discretized, gives

$$\underline{K} \Delta \underline{U}^{(k)} = {}^{t+\Delta t} \underline{R} - \underbrace{{}^{t+\Delta t} \underline{F}^{(k-1)}}_{\substack{\text{computed} \\ \text{from } {}^{t+\Delta t} \underline{u}_i^{(k-1)}}}$$

(for $k = 1, 2, 3, \dots$)

Note that ${}^{t+\Delta t} \underline{U}^{(k)} = {}^t \underline{U} + \sum_{j=1}^k \Delta \underline{U}^{(j)}$.



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Comparison of T.L. and U.L. formulations

- In the T.L. formulation, all derivatives are with respect to the initial coordinates whereas in the U.L. formulation, all derivatives are with respect to the current coordinates.
- In the U.L. formulation we work with the actual physical stresses (Cauchy stress).

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The same assumptions are made in the linearization and indeed the same finite element stiffness and force vectors are calculated (when certain transformation rules are followed).

MIT OpenCourseWare
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Resource: Finite Element Procedures for Solids and Structures
Klaus-Jürgen Bathe

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